

Gravitational collapse and integrable singularities

Roberto Casadio

D.I.F.A. “A. Righi”

Bologna University

INFN

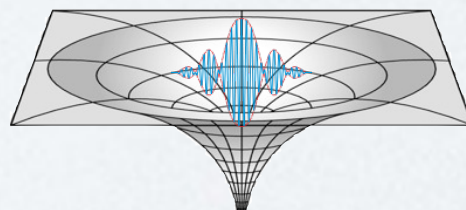
From Black Holes to the Cosmos

26/08/2026



Theory and Phenomenology
of Fundamental Interactions

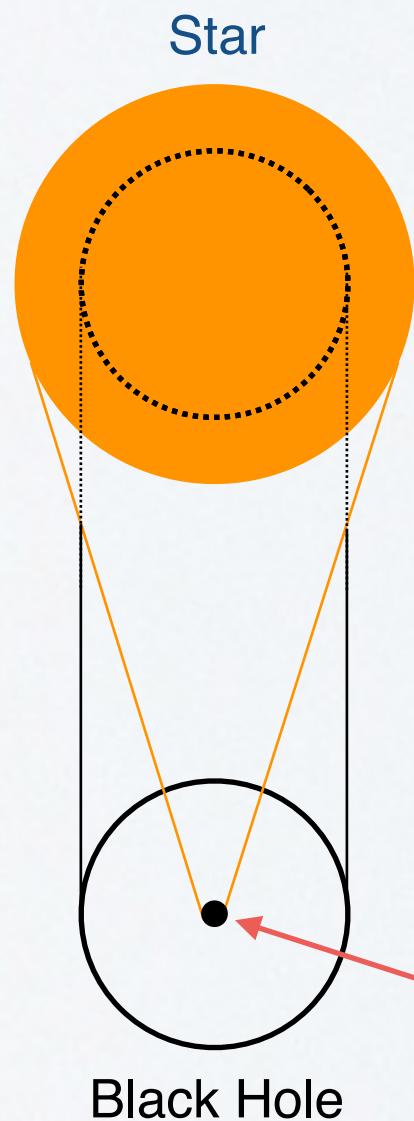
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- **Preamble:** gravitational collapse
- Regular and integrable densities
- **Toy kinematic model:** collapsing Schwarzschild mimickers
- **Quantum Gravity:** gravitational collare and dust core

Preamble - QG and gravitational collapse

BH form from collapse - the classical view:



- $|\text{matter}\rangle \sim$ very large number of SM particles ($M_{\odot} \sim 10^{57}$ neutrons)
- $|\text{gravity}\rangle \sim$ very large number of “gravitons” [1] ($N_G \sim M_{\odot}^2/m_p^2 \sim 10^{76}$)



Reduce and (not over-)simplify
(how...?)

Singularity ~ CED

QED

... but ...

GR ~ QCD

“Atomic collapse”

Atoms

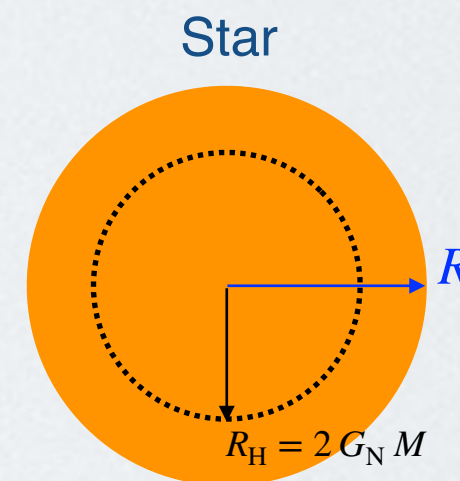
Regular and Integrable densities

MisnerSharpHernandez mass: $ds^2 = g_{tt} dt^2 + g_{rr} dr^2 + r^2 d\Omega^2$

↓ Einstein equations

$$g^{rr} = 1 - \frac{2 G_N m(r)}{r}$$

$$m(r) = 4 \pi \int_0^r \rho(x) x^2 dx \quad 0 \leq r \leq R$$



Regular (central) density: $\rho(r \sim 0) \sim \rho_0 \implies m(r \sim 0) \sim r^3$



Finite geometric invariants $\sim$ Finite tidal forces



Classical configuration

Regular and Integrable densities

MSH mass function [1]:

$$ds^2 = g_{tt} dt^2 + g_{rr} dr^2 + r^2 d\Omega^2$$



Einstein equations

$$g^{rr} = 1 - \frac{2 G_N m(r)}{r}$$

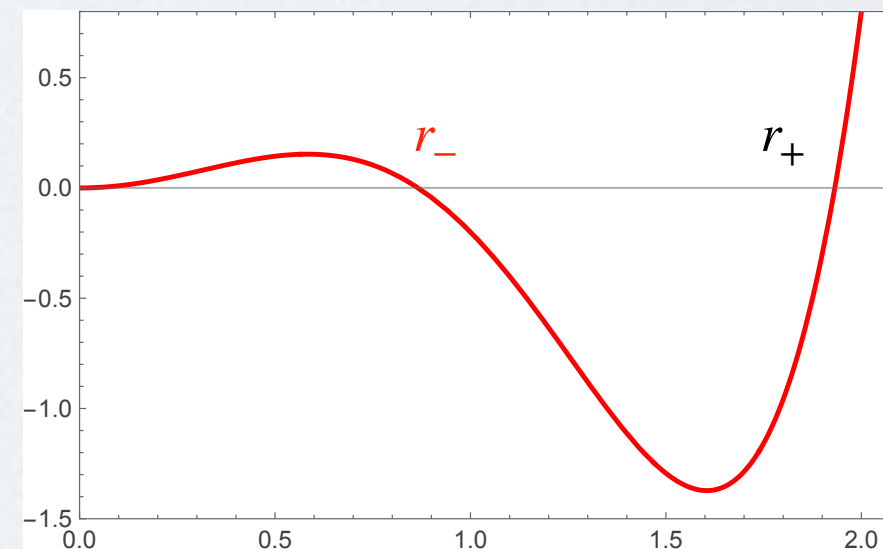
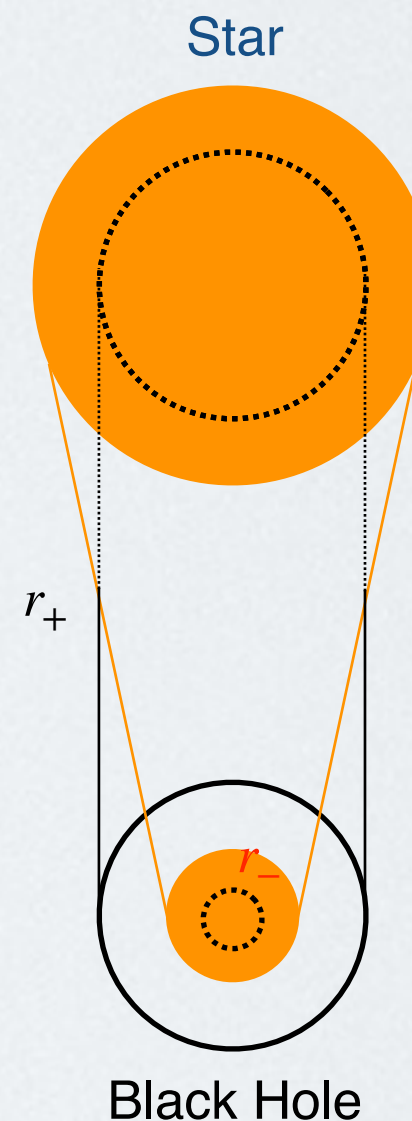
$$m(r) = 4\pi \int_0^r \rho(x) x^2 dx \quad 0 \leq r \leq R < r_+ = R_H$$

Regular (central) density: $\rho(r \sim 0) \sim \rho_0 \implies m(r \sim 0) \sim r^3$

Horizon(s): $g^{rr} = 0 \implies \Delta = r_{\pm}^2 - 2 r_{\pm} m(r_{\pm}) = 0$



Inner Cauchy horizon



Regular and Integrable densities

MSH mass function [1]:

$$ds^2 = g_{tt} dt^2 + g_{rr} dr^2 + r^2 d\Omega^2$$



Einstein equations

$$g^{rr} = 1 - \frac{2 G_N m(r)}{r}$$

$$m(r) = 4 \pi \int_0^r \rho(x) x^2 dx \simeq 4 \pi \int_0^r \mu |\psi(x)|^2 x^2 dx < \infty$$

Regular (central) density: $\rho(r) \sim \rho_0 \implies m(r) \sim r^3$

Integrable density: $m(r) \sim r^{1+n} < \infty \implies \rho(r) \sim r^{n-2} \quad (-1 < n < 2)$

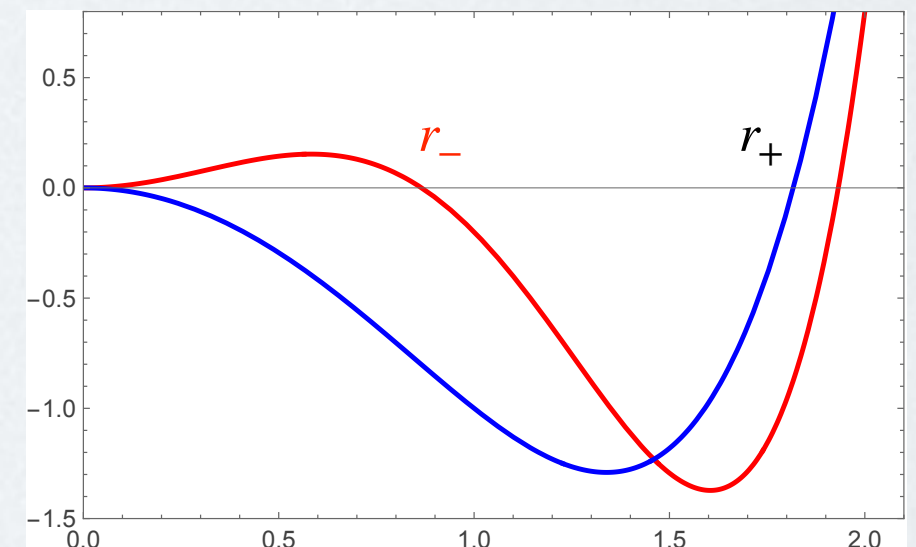
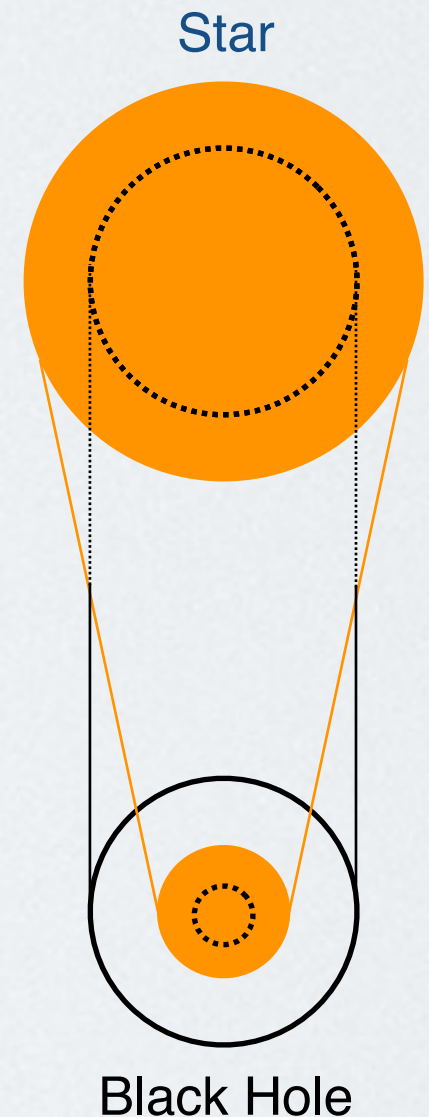
Horizon(s): $g^{rr} = 0 \implies \Delta = r_{\pm}^2 - 2 r_{\pm} m(r_{\pm}) = 0$

$$r = r_{\pm} \quad (0 < n < 2)$$

$$r = r_{+} \quad (-1 < n < 0)$$



“Quantum” configuration ~ no Cauchy horizon



Regular and Integrable densities

MSH mass function [1,2]:

$$ds^2 = g_{tt} dt^2 + g_{rr} dr^2 + r^2 d\Omega^2$$



Einstein equations

$$g^{rr} = 1 - \frac{2 G_N m(r)}{r}$$

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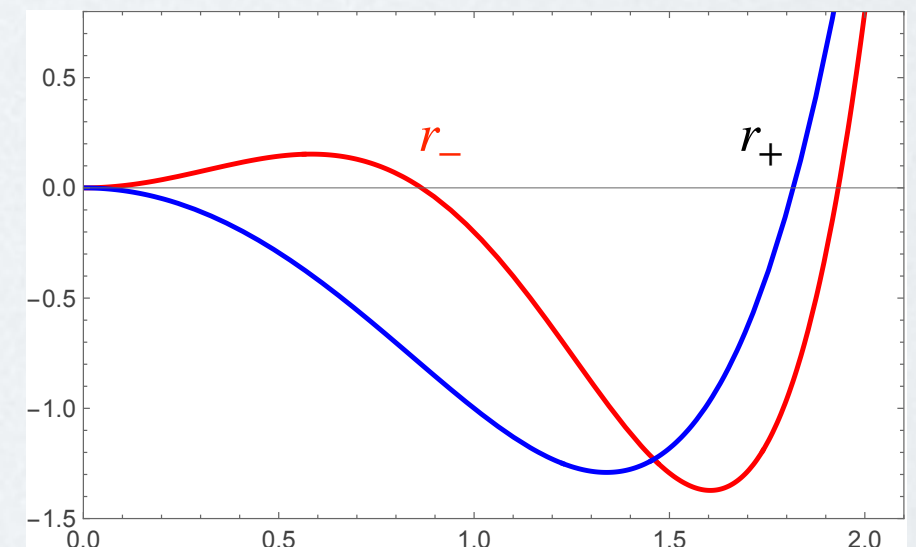
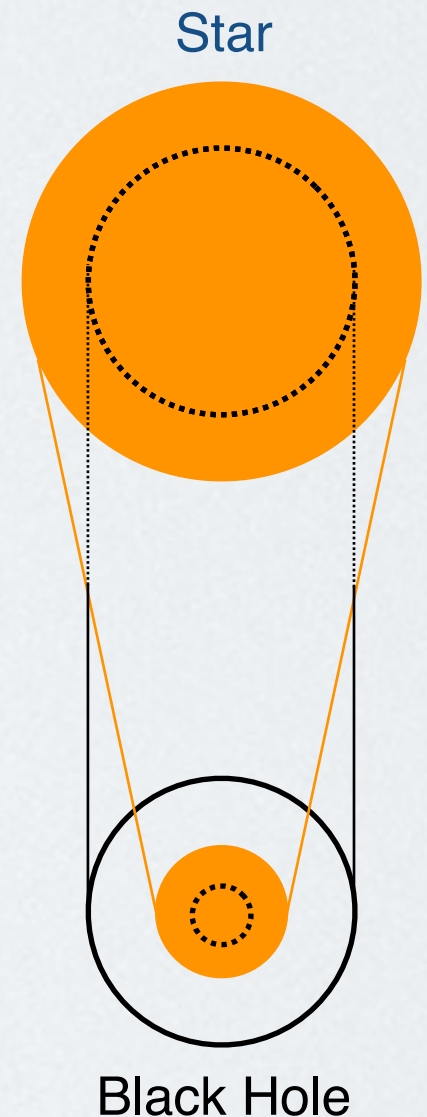
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Schwarzschild: $m(r) = M \implies \Delta = r_H^2 - 2 r_H G_N M \quad (n \rightarrow -1)$



[1] R.C., A. Giusti, J. Ovalle, JHEP 05 (2023) 118 [arXiv:2303.02713]

[2] R.C., A. Giusti, A. Kamenshchik, J. Ovalle, to appear in EPJC [arXiv:2605.01808]

Collapsing Schwarzschild mimickers

Mimicker geometry [1,2,3]:

$$ds^2 = -f(r) dv^2 + 2 dv dr + r^2 d\Omega^2$$

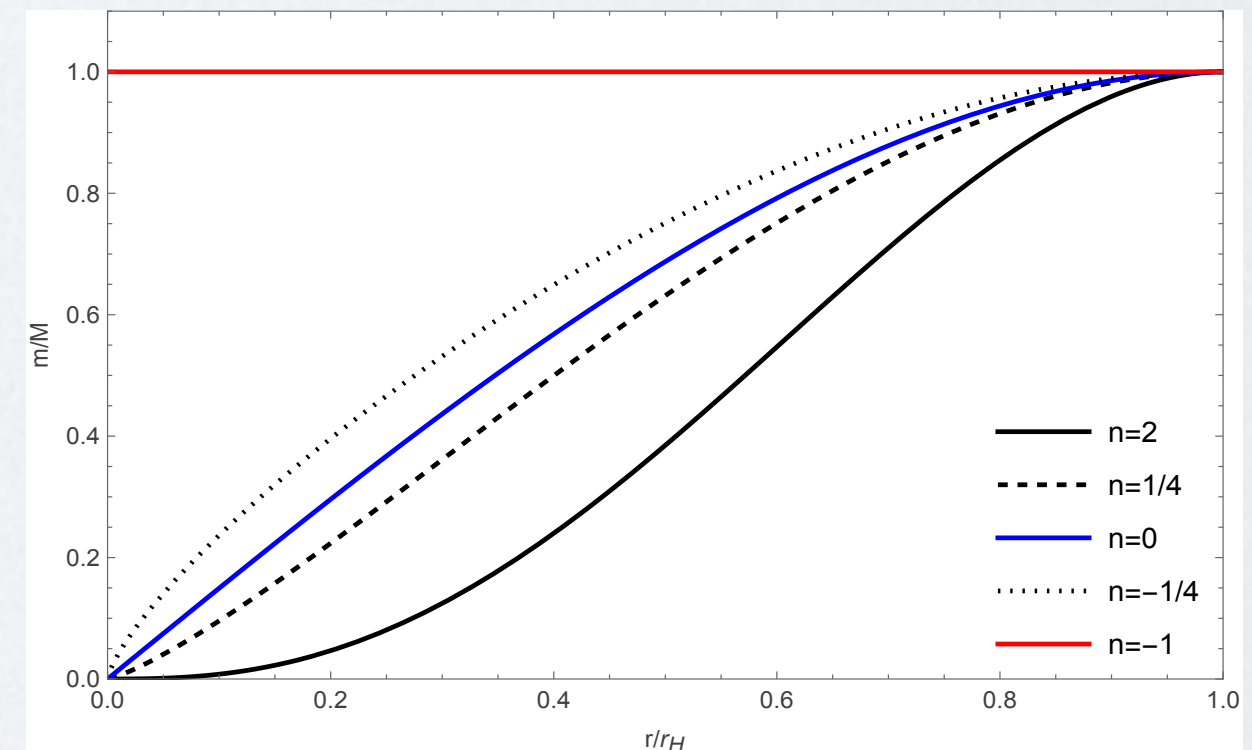
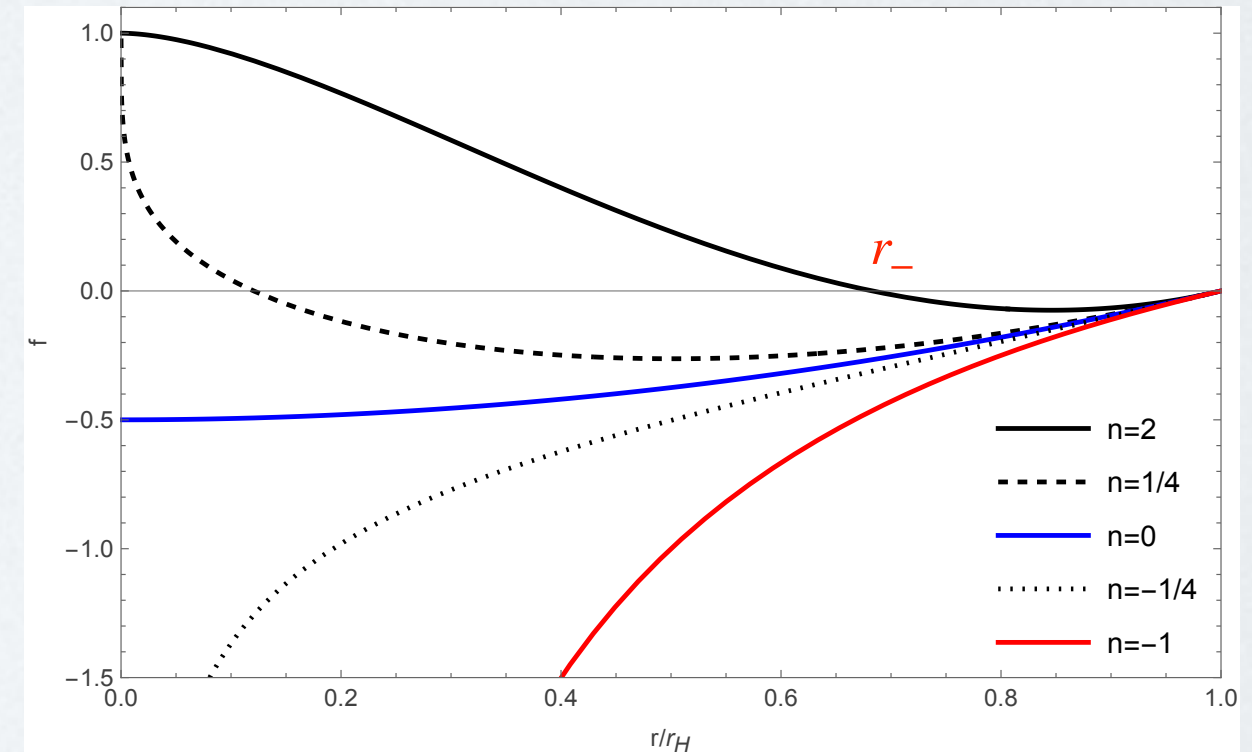
(ingoing light rays: $v = \text{constant}$)

$$f = \begin{cases} 1 - \frac{2 G_N m(r)}{r} & 0 \leq r \leq r_H = 2 G_N M \\ 1 - \frac{2 G_N M}{r} & r \geq r_H \end{cases}$$

Minkowski breaking [2,3]

MSH mass function [1,2,3]: $f(r), m(r) \in C^1(\mathbb{R}^+)$

$$m = \frac{r}{2 G_N (n-2)} \left[(n+1) \frac{r^2}{r_H^2} - 3 \left(\frac{r}{r_H} \right)^n \right]$$



[1] J. Ovalle, PRD 109 (2024) 104032 [arXiv:2405.06731]; J. Ovalle, arXiv:2509.00816

[2] J. Ovalle, R.C. A. Kamenshchik, PRD113 (2026) 064042 [arXiv:2603.06451]

[3] R.C., A. Giusti, A. Kamenshchik, J. Ovalle, to appear in EPJC [arXiv:2605.01808]

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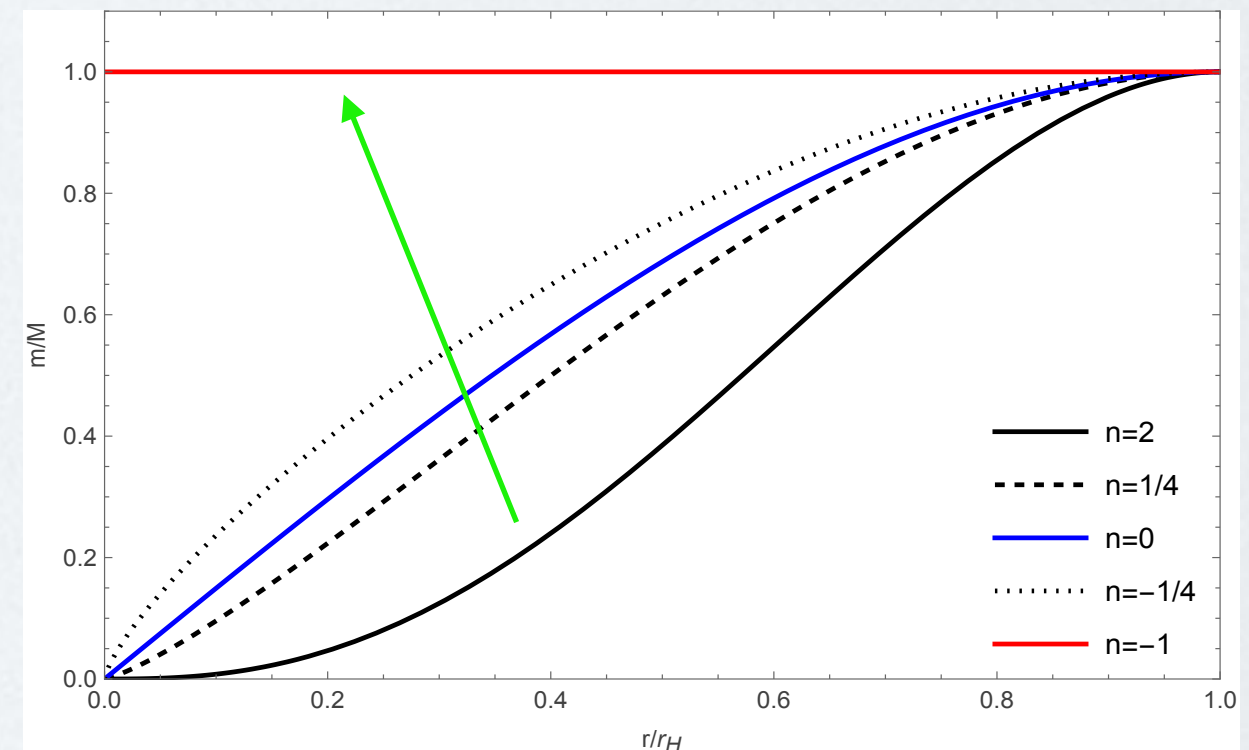
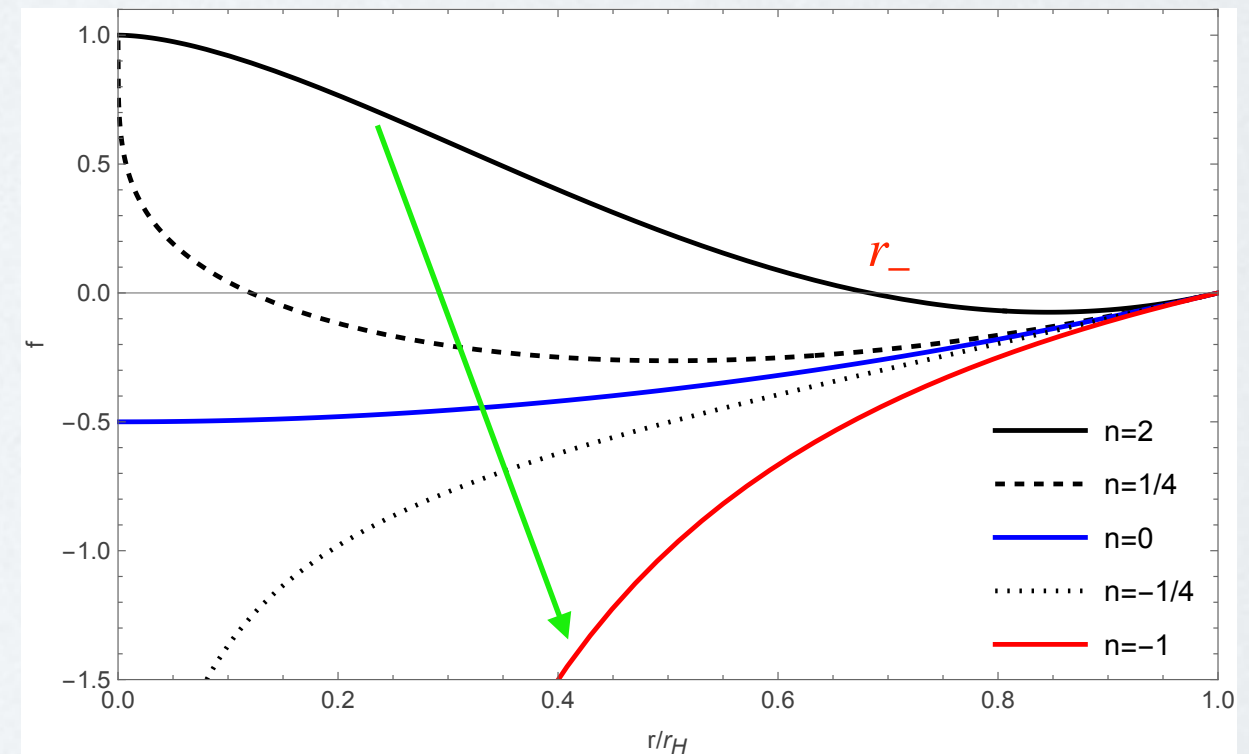
$$f = \begin{cases} 1 - \frac{2 G_N m(v, r)}{r} & 0 \leq r \leq r_H = 2 G_N M \\ 1 - \frac{2 G_N M}{r} & r \geq r_H \end{cases}$$

Minkowski breaking [2,3]

MSH mass function [1,2,3]:

$$m = \frac{r}{2 G_N [n(v) - 2]} \left\{ [n(v) + 1] \frac{r^2}{r_H^2} - 3 \left(\frac{r}{r_H} \right)^{n(v)} \right\}$$

Time evolution $\sim \dot{n}(v) < 0$: $n > 2 \rightarrow -1$



[1] J. Ovalle, PRD 109 (2024) 104032 [arXiv:2405.06731]; J. Ovalle, arXiv:2509.00816

[2] J. Ovalle, R.C. A. Kamenshchik, PRD113 (2026) 064042 [arXiv:2603.06451]

[3] R.C., A. Giusti, A. Kamenshchik, J. Ovalle, to appear in EPJC [arXiv:2605.01808]

Collapsing Schwarzschild mimickers

EM tensor [1,2]:

$$ds^2 = -f(v, r) dv^2 + 2 dv dr + r^2 d\Omega^2$$

$$f = 1 - \frac{2 G_N m(v, r)}{r} \quad \Rightarrow \quad G_{ab} = 8 \pi G_N T_{ab}$$

Tetrad components:

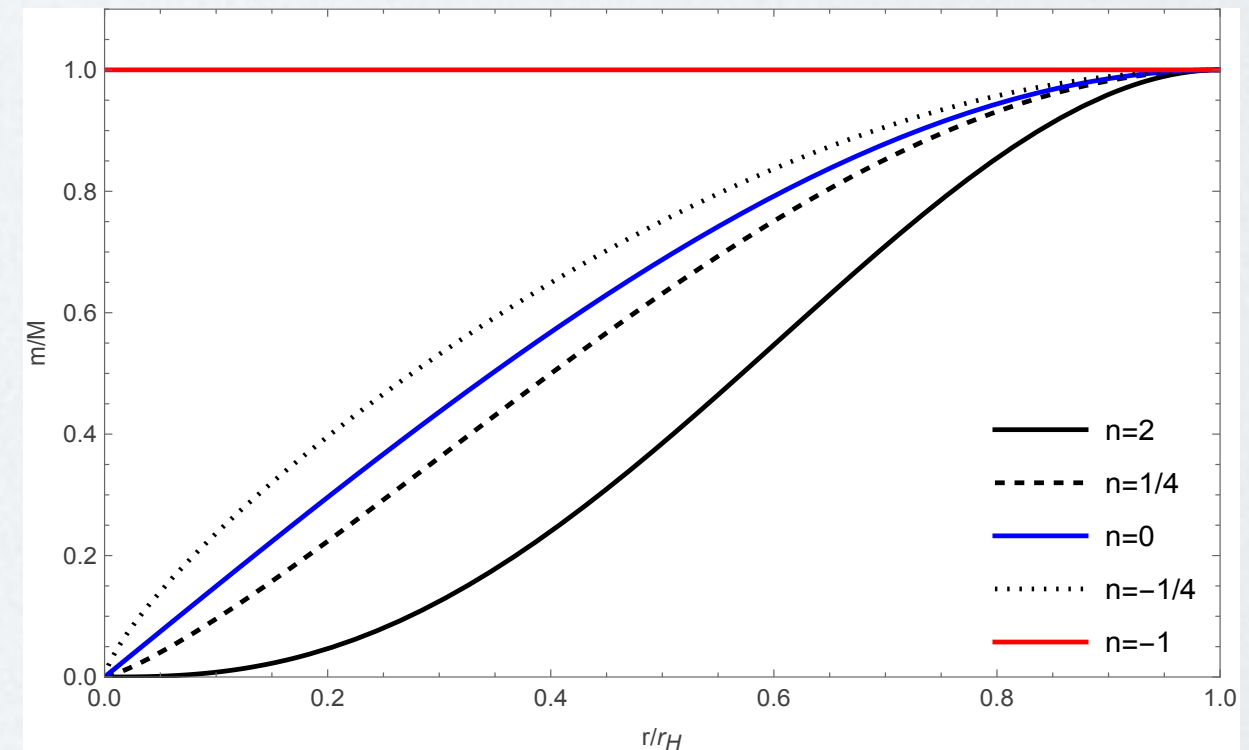
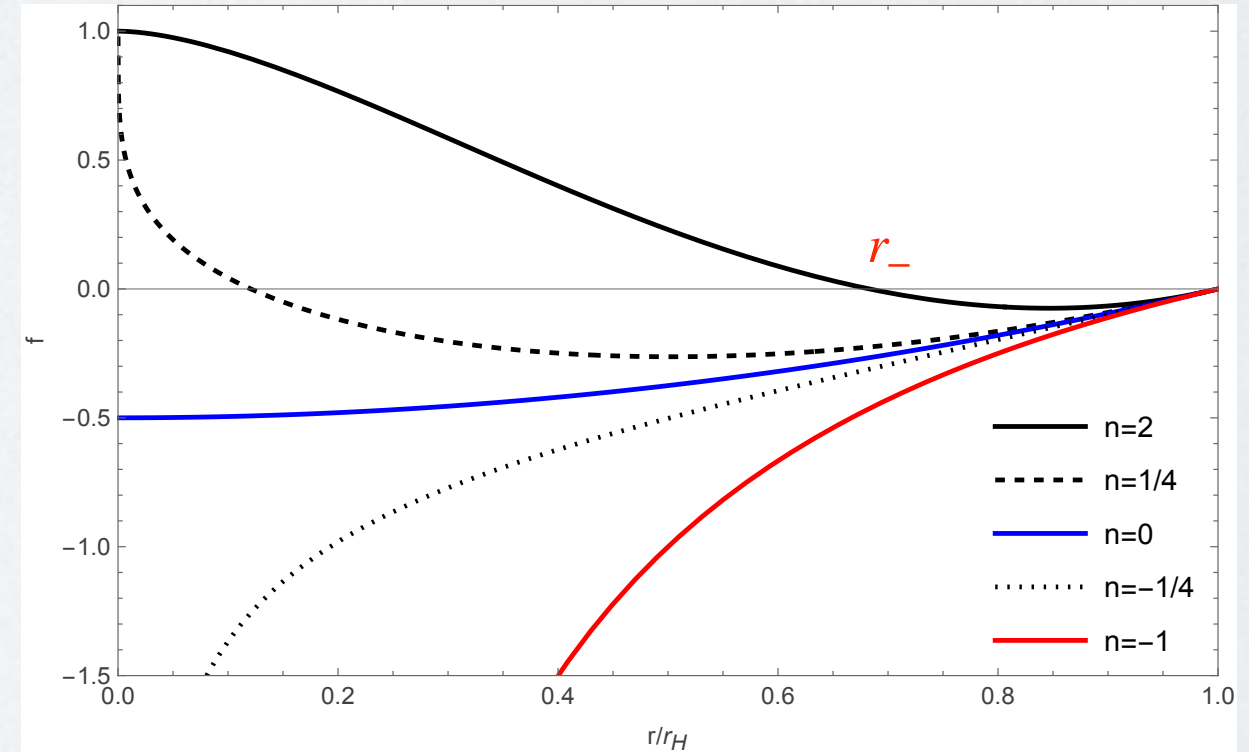
$$g_{ab} = \eta_{ab}$$

$$T_{ab} = \begin{bmatrix} \rho + \Phi/2 & \Phi/2 & 0 & 0 \\ \Phi/2 & p_r + \Phi/2 & 0 & 0 \\ 0 & 0 & p_t & 0 \\ 0 & 0 & 0 & p_t \end{bmatrix}$$

$$\rho = -p_r = \frac{3}{8 \pi G_N r_H^2} \left(\frac{n+1}{n-2} \right) \left[1 - \left(\frac{r}{r_H} \right)^{n-2} \right]$$

$$p_t = -\frac{3}{8 \pi G_N r_H^2} \left(\frac{n+1}{n-2} \right) \left[1 - \frac{n}{2} \left(\frac{r}{r_H} \right)^{n-2} \right]$$

$$\Phi = -\frac{3 r \dot{n}}{8 \pi G_N r_H^2 (n-2)^2} \left\{ 1 - \left(\frac{r}{r_H} \right)^{n-2} \left[1 + (n-2) \ln \left(\frac{r_H}{r} \right) \right] \right\}$$



[1] A. Wang, A. Wu, Gen. Rel. Grav. 31 (1999) 107 [arXiv:gr-qc/9803038]

[2] R.C., A. Giusti, A. Kamenshchik, J. Ovalle, to appear in EPJC [arXiv:2605.01808]

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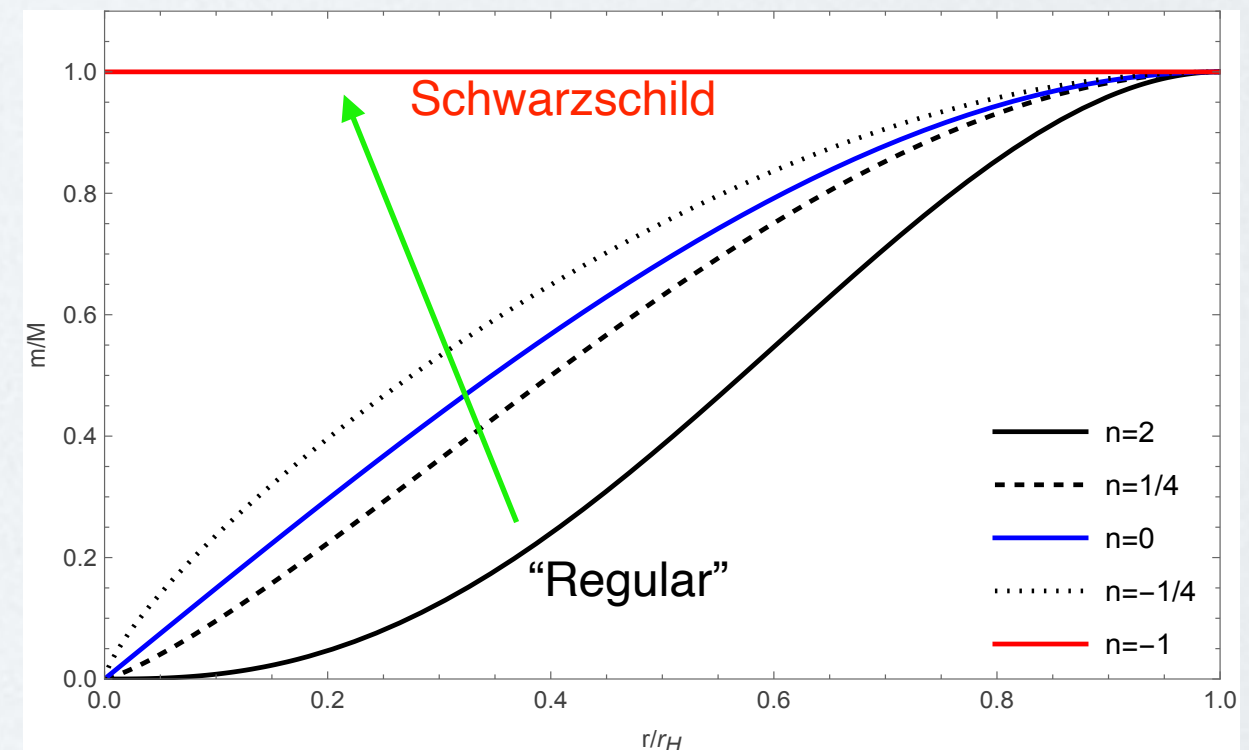
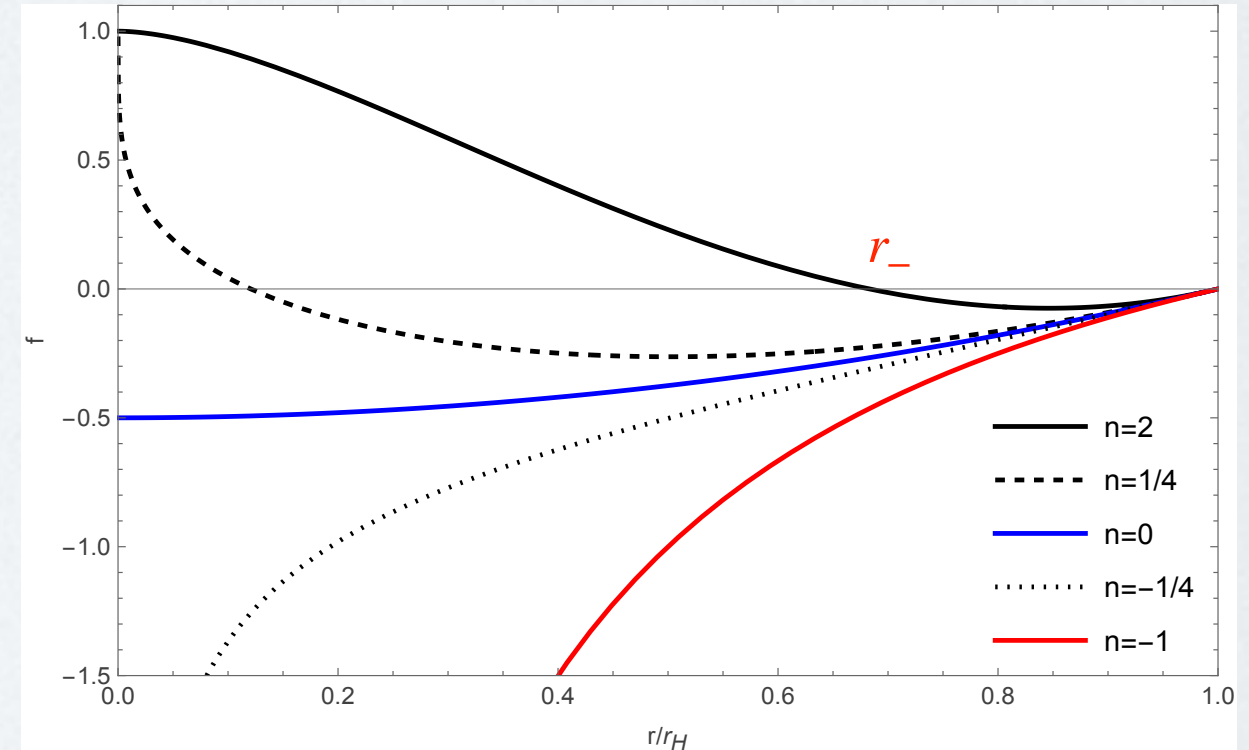
$$m(r) \sim r^3 \quad (n > 2)$$

$$m(r) \sim r^3 \ln(r) \quad (n = 2)$$

$$m(r) \sim r^{1+n} \quad (-1 < n < 2)$$

$$m(r) = M \quad (n \rightarrow -1^+)$$

$$4 \pi r^2 \Phi \sim \dot{n} r^{1+n} \ln(r) \quad \Rightarrow \quad \dot{n} \rightarrow 0 \quad (n \rightarrow -1^+)$$



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[2] R.C., A. Giusti, A. Kamenshchik, J. Ovalle, to appear in EPJC [arXiv:2605.01808]

Collapsing Schwarzschild mimickers

Classical static end-state (before Schwarzschild singularity) [1]?

Geodesic 4-velocity: $u^\mu \nabla_\mu u^\nu = 0$

$$u^\mu u_\mu = -1$$

$$u^\mu = \left(-\frac{1}{1 + \sqrt{1-f}}, \sqrt{1-f}, 0, 0 \right)$$

$f = f(r)$



$$h_{\mu\nu} = g_{\mu\nu} + u_\mu u_\nu$$

$$B_{\mu\nu} = \nabla_\mu u_\nu$$

$$\theta = B^\mu{}_\mu = 0$$

(expansion)

$$\sigma_{\mu\nu} = \frac{1}{2} (B_{\mu\nu} + B_{\nu\mu}) - \frac{1}{3} \theta h_{\mu\nu} = 0$$

(shear)

$$\omega_{\mu\nu} = \frac{1}{2} (B_{\mu\nu} - B_{\nu\mu}) = 0$$

(rotation)



Raychaudhuri equation: $\frac{d\theta}{d\tau} = -\frac{\theta^2}{3} - \sigma_{\mu\nu} \sigma^{\mu\nu} + \omega_{\mu\nu} \omega^{\mu\nu} - R_{\mu\nu} u^\mu u^\nu = -R_{\mu\nu} u^\mu u^\nu = 0$



$$R_{\mu\nu} = 0$$

“Regular”



Schwarzschild vacuum!

Collapsing Schwarzschild mimickers

Quantum (Madelung) approximation [1]:

$$\mu |\psi|^2 \simeq \rho = \frac{3}{8\pi G_N r_H^2} \left(\frac{n+1}{n-2} \right) \left[1 - \left(\frac{r}{r_H} \right)^{n-2} \right]$$

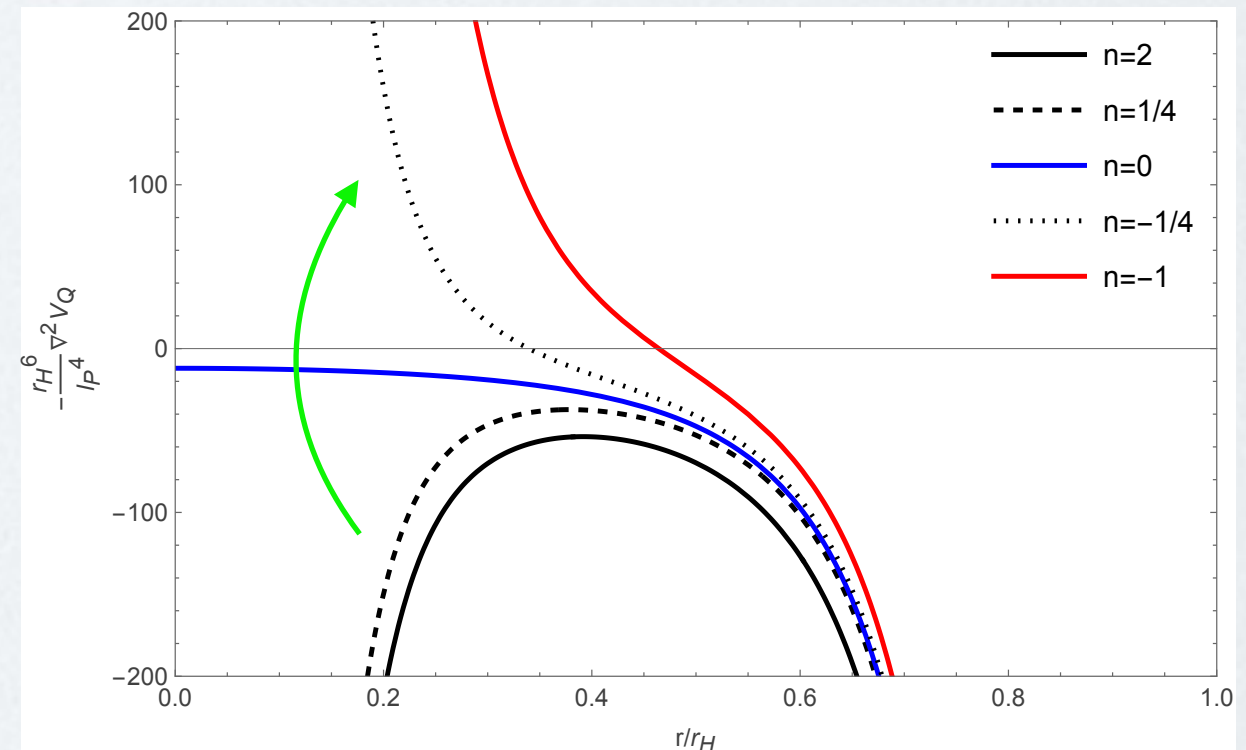
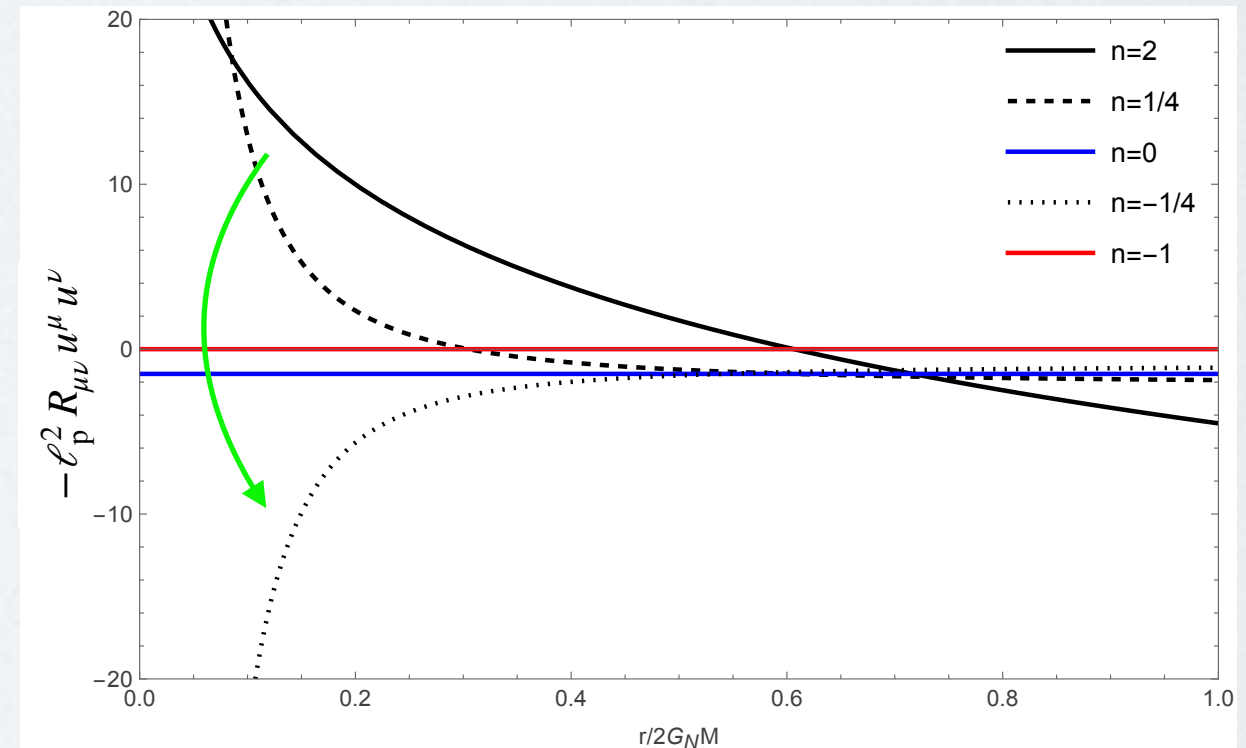
Quantum (Madelung) potential [2,3]:



$$\frac{d\theta}{d\tau} = -R_{\mu\nu} u^\mu u^\nu - \nabla^2 V_Q$$

$$V_Q \simeq -2 \frac{\ell_p^2}{r_H^2} \frac{\nabla^2 |\psi|}{|\psi|} \simeq -\frac{\ell_p^2}{r_H^2} \frac{\nabla^2 \rho}{\rho}$$

Minkowski breaking [3]



$$-\nabla^2 V_Q \simeq -\frac{n \ell_p^4}{r_H^2 r^4}$$

$$-\nabla^2 V_Q \simeq -\frac{\ell_p^4 / r_H^2}{(r - r_H)^4}$$

[1] E. Madelung, Z. Phys. 40 (1927) 322

[2] S. Das, PRD 89 (2014) 084068 [arXiv:1311.6539]

[3] R.C., A. Giusti, A. Kamenshchik, J. Ovalle, to appear in EPJC [arXiv:2605.01808]

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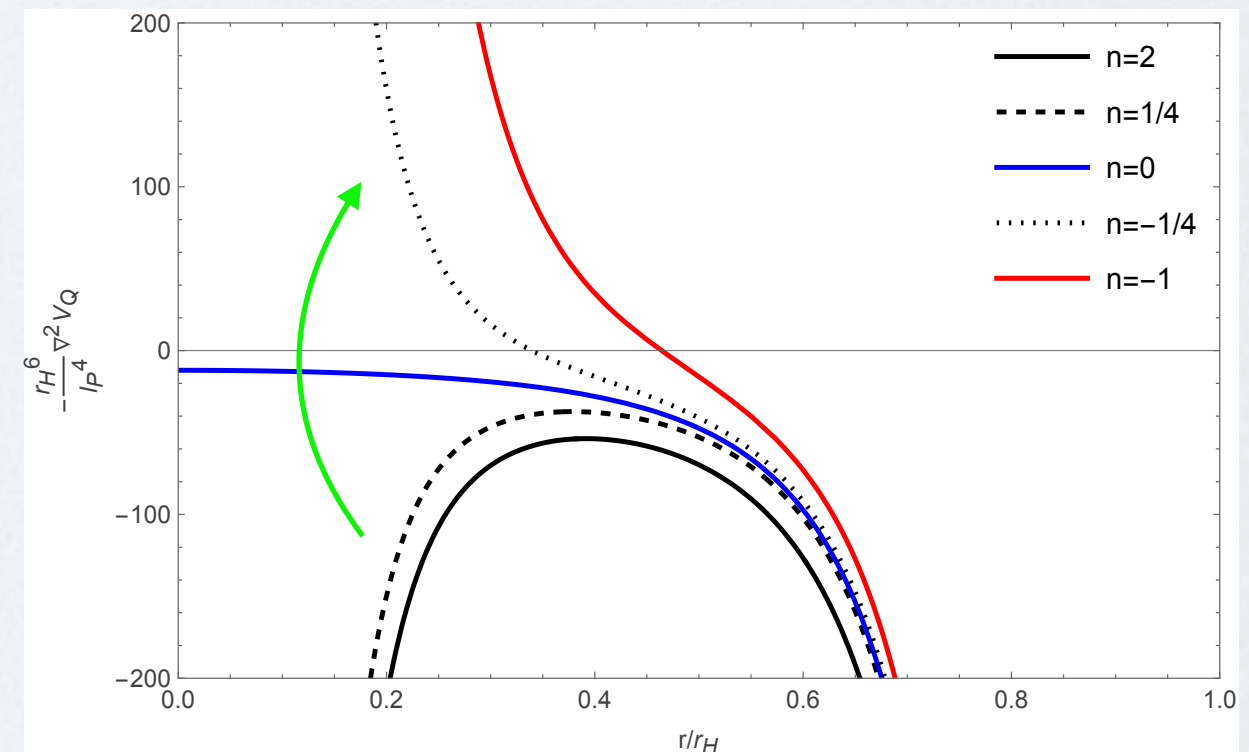
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Minkowski breaking [3]

Outward quantum pressure
(after Minkowski breaking)

Unstable horizon
(Outer Boulware vacuum
Hawking radiation?)



$$-\nabla^2 V_Q \simeq -\frac{n \ell_p^4}{r_H^2 r^4} \quad (n \neq 0)$$

$$-\nabla^2 V_Q \simeq -\frac{\ell_p^4 / r_H^2}{(r - r_H)^4}$$

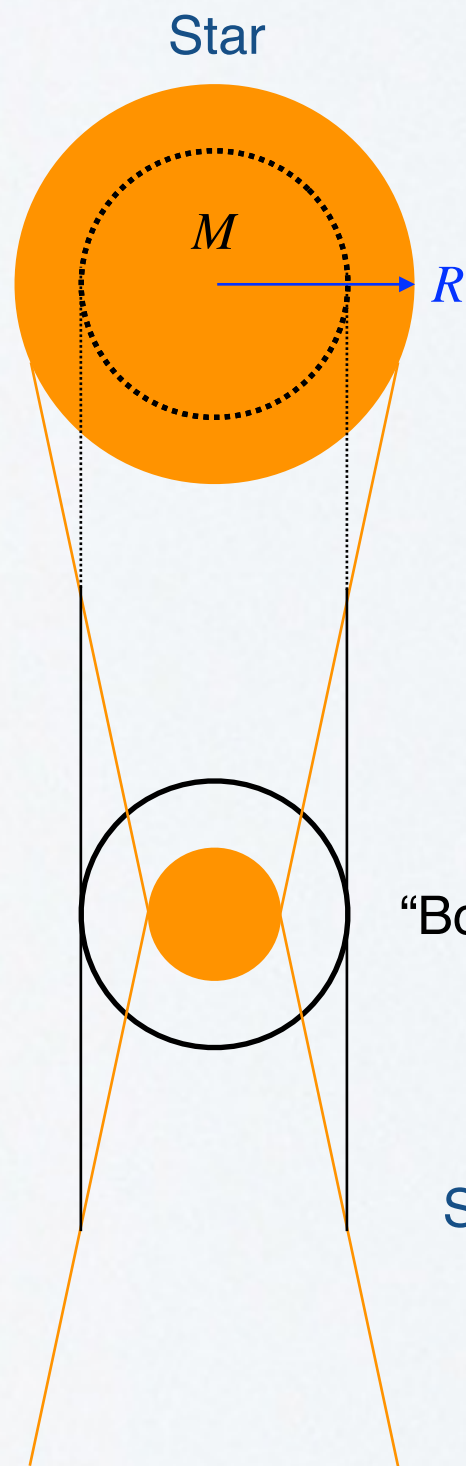
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QG and gravitational collapse

BH form from collapse - the semiclassical view:



- $|\text{matter}\rangle \sim$ very large number of SM particles ($M_{\odot} \sim 10^{57}$ neutrons)
- $|\text{gravity}\rangle \sim$ very large number of gravitons ($N_G \sim M_{\odot}^2/m_p^2 \sim 10^{76}$)



Lemaître-Tolman-Bondi-Oppenheimer-Snyder model

Canonical quantisation of R (and M) [1]



“Bounce” $\sim$ BH to WH tunneling

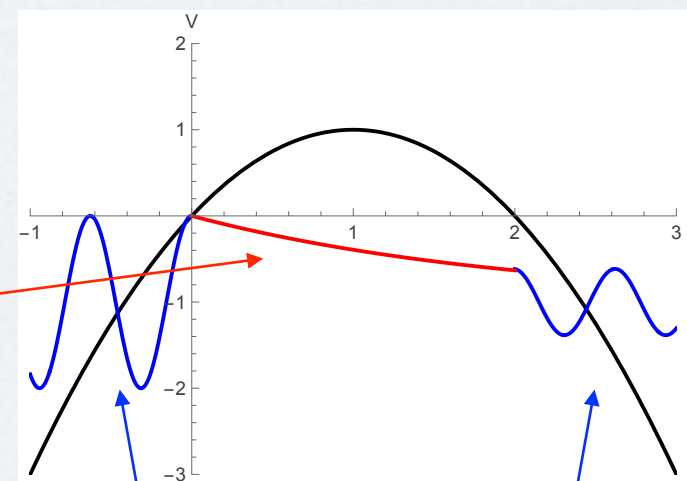


Semiclassical approximation [1,2]



Size of “bounce” ($R < R_H?$)

“Quantum” geometry?



Classical geometry

[1] Kuchar, Hajicek, Vaz, Kiefer, Husain, Giesel, Rovelli...

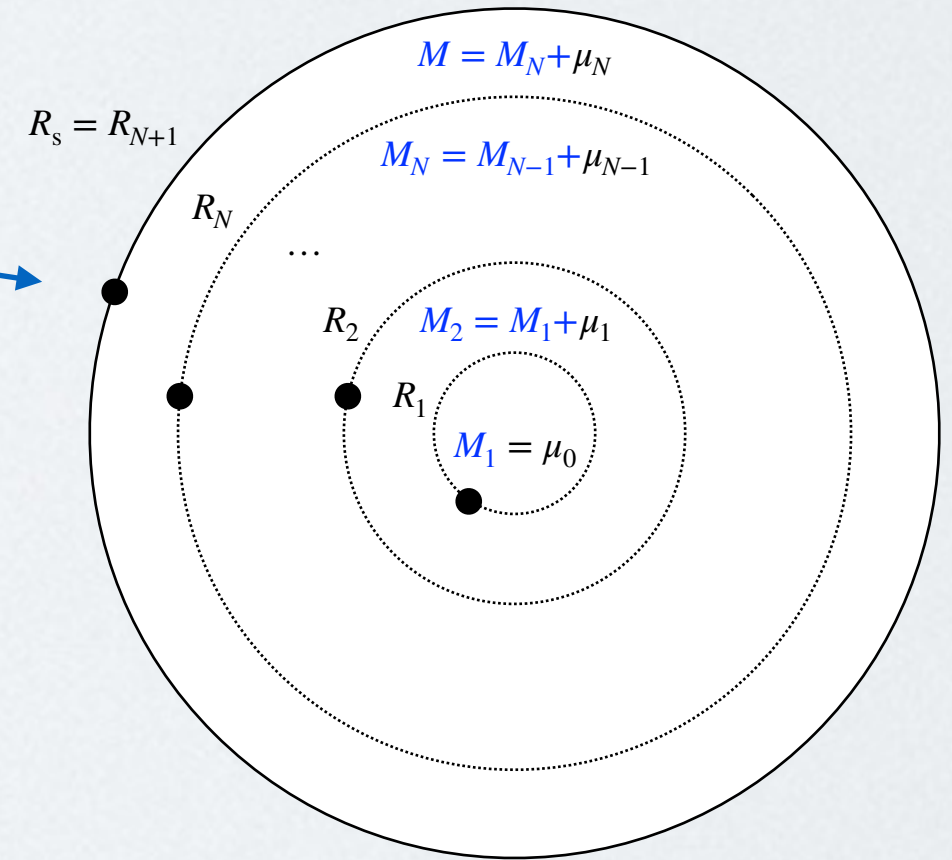
[2] R.C. IJMPD 9 (2000) 511 [arXiv:gr-qc/9810073]

Quantum dust core

Collapsing ball of dust [2]

$$ds^2 = - \left(1 - \frac{2 G_N M_i}{r} \right) dt^2 + \left(1 - \frac{2 G_N M_i}{r} \right)^{-1} dr^2 + r^2 d\Omega^2$$

$$\left(\frac{dR_i}{d\tau} \right)^2 + 1 - \frac{2 G_N M_i}{R_i} \simeq \frac{E_i^2}{\mu^2}$$



Dust particle Hamiltonians:

$$H_i \equiv \frac{P_i^2}{2\mu} - \frac{G_N \mu M_i}{R_i} = \frac{\mu}{2} \left(\frac{E_i^2}{\mu^2} - 1 \right) \equiv \mathcal{E}_i$$

Schrödinger equation:

$$\hat{H}_i \Psi_{n_i} = \mathcal{E}_{n_i} \Psi_{n_i}$$

Spectrum of bound states ($n_i \geq 1$):

$$\frac{\mathcal{E}_{n_i}}{\mu} \simeq - \frac{G_N^2 \mu^3 M_i}{2 \hbar^2 n_i^2} = - \frac{1}{2 n_i^2} \left(\frac{\mu M_i}{m_p^2} \right)^2 = \frac{1}{2} \left(\frac{E_i^2}{\mu^2} - 1 \right)$$

“GR” [2]

$$\bar{R}_{n_i} \equiv \langle \Psi_{n_i} | \hat{R}_i | \Psi_{n_i} \rangle \simeq n_i^2 \ell_p \left(\frac{m_p^3}{\mu^2 M_i} \right)$$

Newtonian spectrum

Quantum dust core

Allowed spectrum [1,2]:

$$0 \leq \frac{E_i^2}{\mu^2} \simeq 1 - \frac{1}{n_i^2} \left(\frac{\mu M_i}{m_p^2} \right)^2 \longrightarrow n_i \geq N_i = \frac{\mu M_i}{m_p^2} \longrightarrow \bar{R}_{N_i} \simeq \frac{3}{2} G_N M_i$$

Layer thickness [1]:

$$\frac{\Delta \bar{R}_{n_i}}{\bar{R}_{n_i}} = \frac{\sqrt{n_i^2 + 2}}{3 n_i} \simeq \frac{1}{3} \longrightarrow \bar{R}_{N_{i+1}} \simeq \bar{R}_{N_i} + \Delta \bar{R}_{N_i} \simeq \frac{4}{3} \bar{R}_{N_i} \longrightarrow M_{i+1} \simeq \frac{4}{3} M_i$$

Bekenstein-like area law:

$$N_G = \frac{M}{\mu} N_N \simeq \frac{M^2}{m_p^2}$$

$$R_s \simeq \frac{3}{2} G_N M = \frac{3}{4} R_H$$

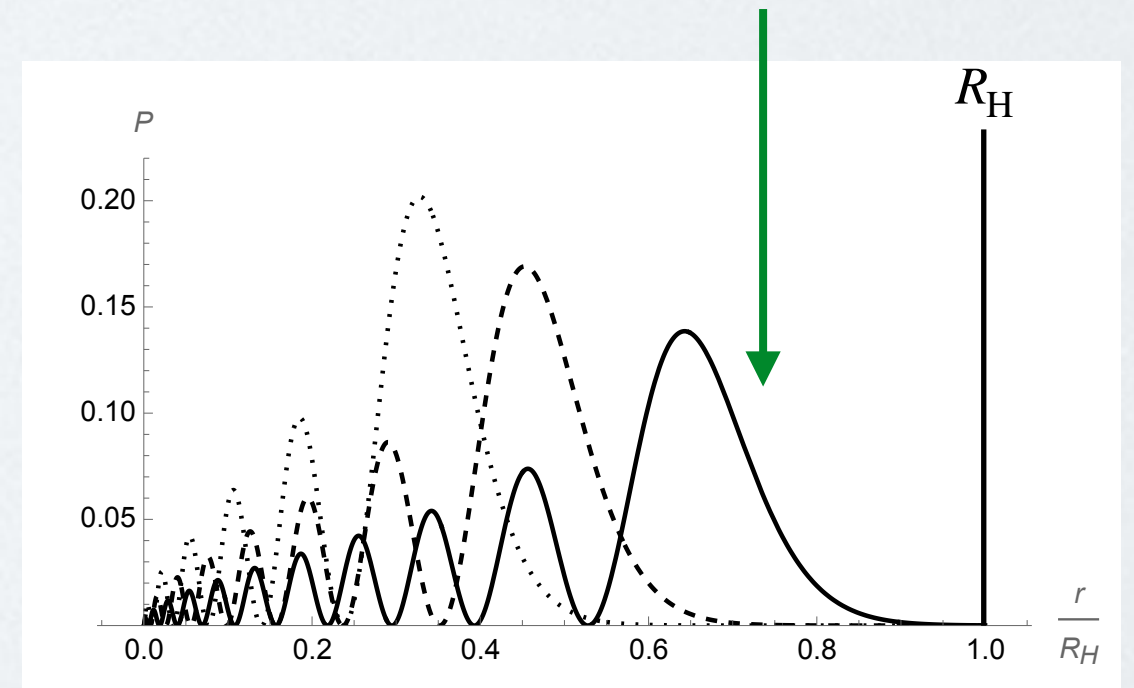
Effective mass and energy density [1]*:

$$m(r) \sim r$$



Integrable singularity ($n = 0$):

$$\rho(r) \sim r^{-2}$$

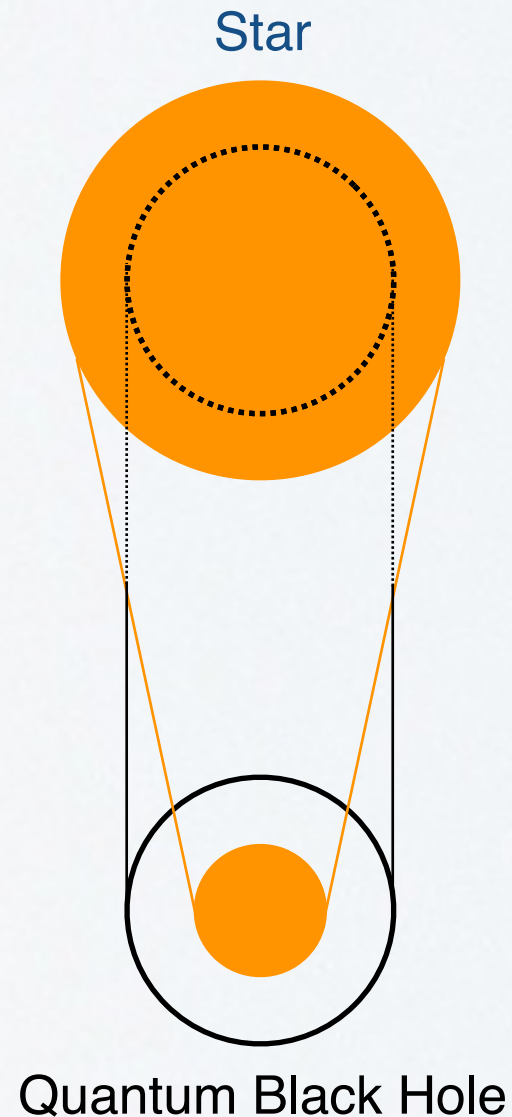


[1] R.C., *Quantum dust cores of black holes*, PLB 843 (2023) 138055 [arXiv:2304.06816]

[2] R.C., *A quantum bound for the compactness*, EPJC 82 (2022) 1 [arXiv:2103.14582]

Conclusion - QG and gravitational collapse

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