

Turbulent plasmas as high-energy particle accelerators: (some) theory and applications

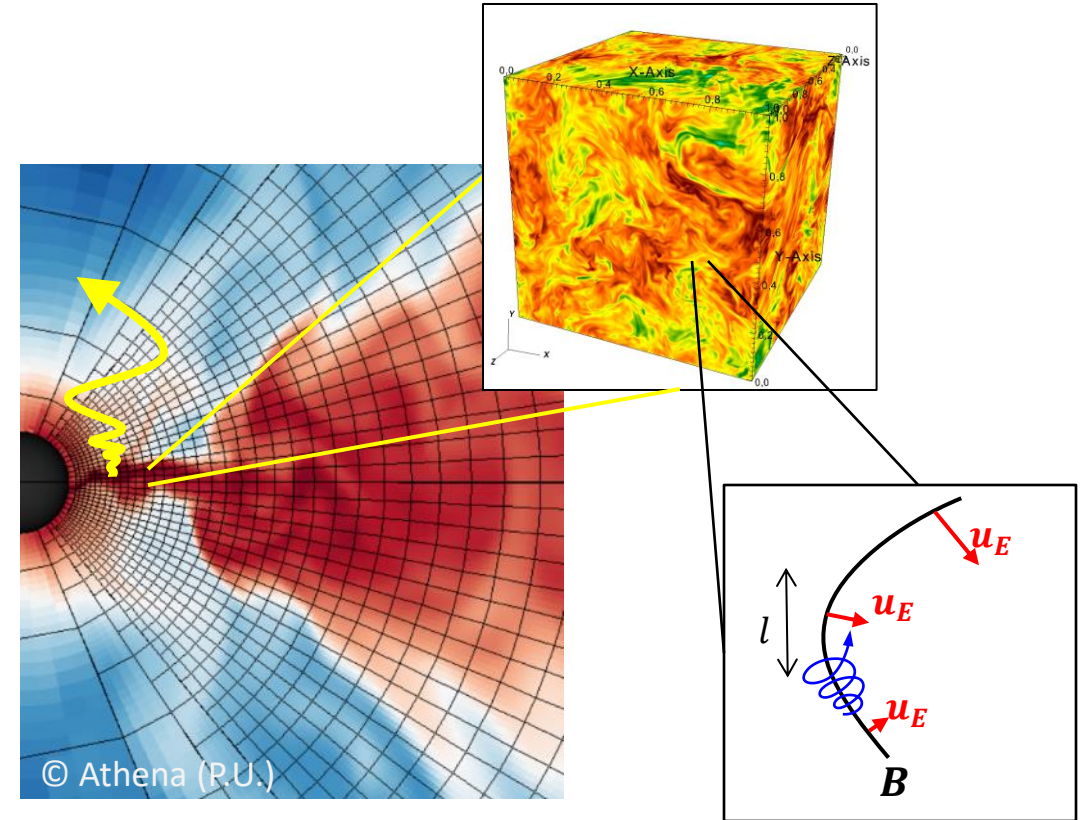
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Outline/summary:

- Theoretical modeling of stochastic particle acceleration:
... *generalized Fermi mechanism in sharp field bends and compression regions*
- Application to PWNe and Crab flares
... *relativistic turbulence as an extreme particle accelerator, radiation at synchrotron burn-off limit*
- Application to AGN coronae as neutrino sources
... *turbulent co-energization of electrons and protons in BH corona*
→ *natural account of X-ray (leptonic) and neutrino (TeV) spectra*



Stochastic particle acceleration in turbulence: a universal scheme in astrophysics

→ dissipation/acceleration in any turbulent plasma,
with efficient acceleration at large magnetization σ , large amplitude $\delta B/B$

magnetization: $\sigma = \text{magnetic energy dens.} / \text{plasma energy dens.}$
 $\sim (\text{rms eddy velocity}/c)^2$

e.g, ISM: $\sigma \sim 10^{-9}$... but in VHE sources, $\sigma \gtrsim 0.01$

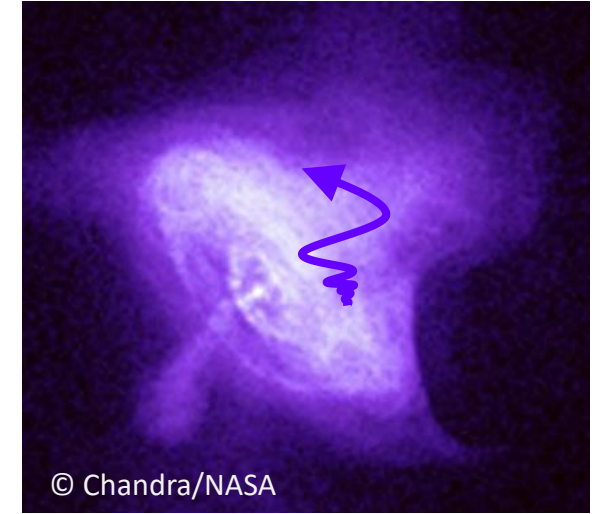
here: relativistic turbulence: $\sigma \gtrsim 0.1$ meaning eddy velocity $\sim c$

→ a stochastic process of acceleration:

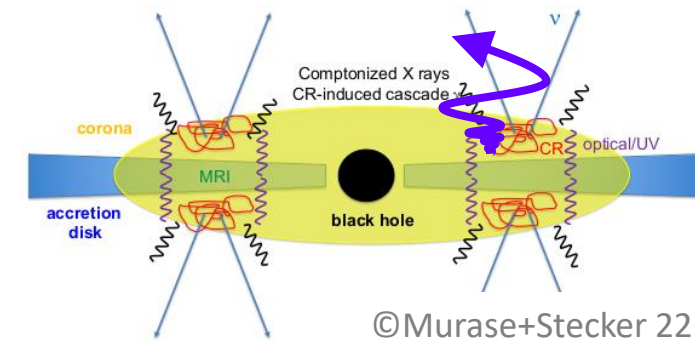
... particles move in energy space by energy gain- or loss- interactions in randomly oriented electric fields

... specific virtues: efficient over broad range of σ + hard spectra expected

... long-standing question: acceleration physics?



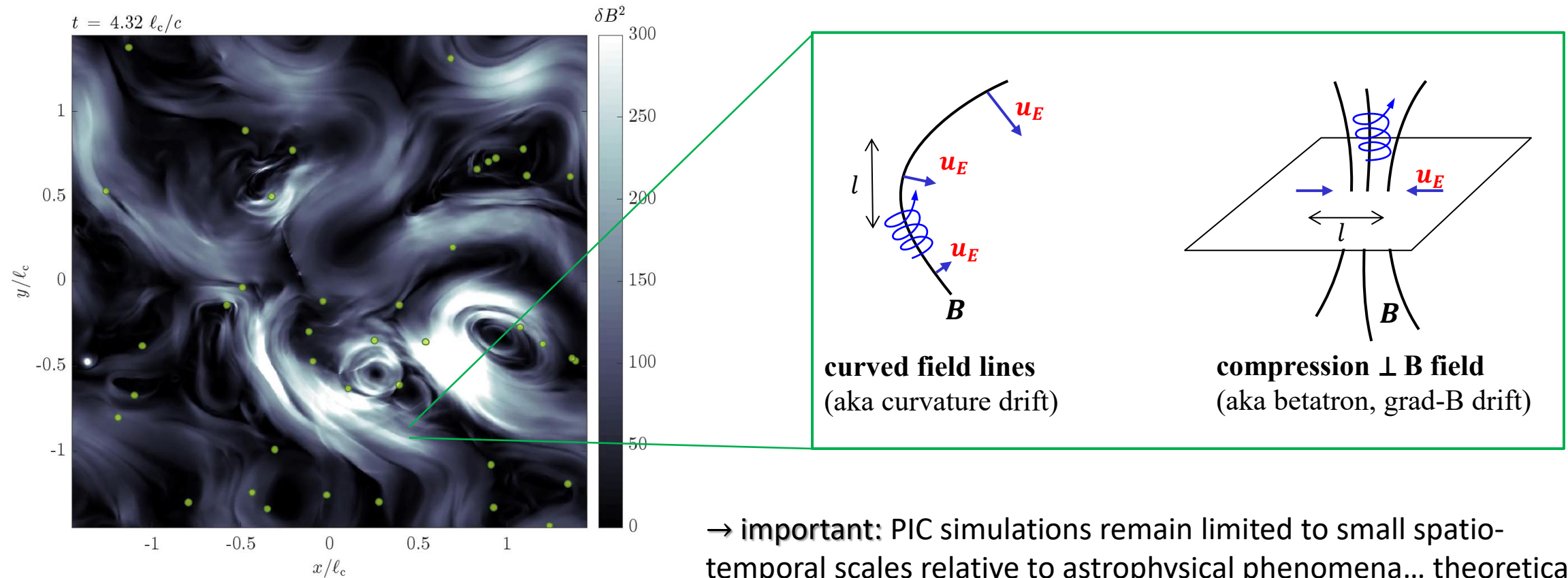
... in pulsar wind nebulae



... in black hole environments

Microphysics of particle in turbulent plasmas: simulations & theory

→ recent progress: MHD + kinetic (PIC) numerical simulations¹ offer direct probes of particle acceleration in turbulence... raising new questions and a direct measurement of acceleration rate (diffusion coefficient)



© V. Bresci, M.L., L. Gremillet (22)

→ important: PIC simulations remain limited to small spatio-temporal scales relative to astrophysical phenomena... theoretical extrapolation needed: theory²

- Refs: 1. Dmitruk+03, ..., Isliker+17, Pecora+18, Zhdankin+2017, Comisso+Sironi 18,19, Wong+20, Trotta+20, Pezzi+22, Pugliese+23, Wong+25,...
2. M.L. 19 [PRD 99, 083006 (2019)], 21 [PRD 104, 063020 (2021)], 22 [PRL 129, 215101 (2022)], 25 [PRE 112, 015205 (2025)];
+ previous works by Webb 85, Bykov+Toptygin 83, Ptuskin 88, Chandran+Maron 04, Cho+Lazarian 06, Petrosian 12, Ohira 13, Brunetti+Lazarian 16, ...

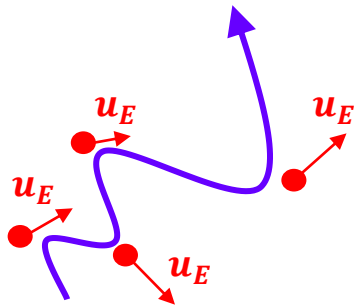
Standard implementation of stochastic acceleration in VHE sources... + outline

→ standard modeling: Fokker-Planck approximation for diffusive processes, fully characterized by diffusion coefficient $D_{\varepsilon\varepsilon}$

$$\partial_t n_\epsilon = \partial_\epsilon (D_{\varepsilon\varepsilon} \partial_\epsilon n_\epsilon) + \dots \quad \text{with } n_\epsilon \equiv dn/d\varepsilon$$

→ functional form of $D_{\varepsilon\varepsilon}(\varepsilon, v_A, \delta B/B, \ell_c, \dots)$ ← acceleration theory... but, which theory?

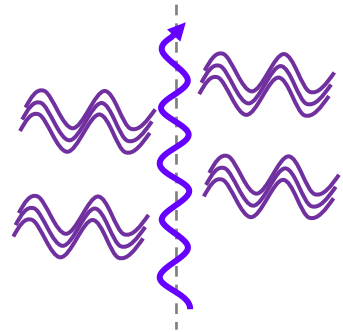
the Fermi pinball¹?



- ... discrete, point-like interactions
- ... $\mathbf{E} = \mathbf{0}$ in scattering center rest frame
- ... net gain: head-on vs tail-on (?)
- ... however, generalization to turbulence?

resonant wave-particle interactions²?

- ... perturbative calculation of $D_{\varepsilon\varepsilon}$



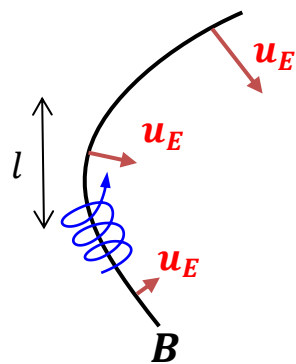
- ... however, waves or structures?

→ here: all about “generalized Fermi” = generalizing Fermi idea to continuous random flow for turbulent plasmas
→ new, interesting phenomenological signatures tested on simulations
+ applications to cases of interest

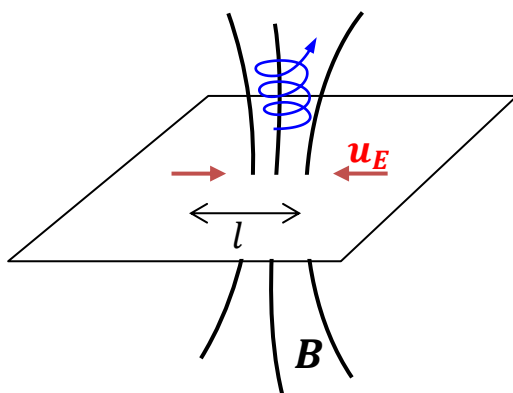
Generalized Fermi acceleration: stochastic acceleration in a continuous random flow

→ an effective theory, covariant & non-perturbative, for modelling stochastic particle acceleration

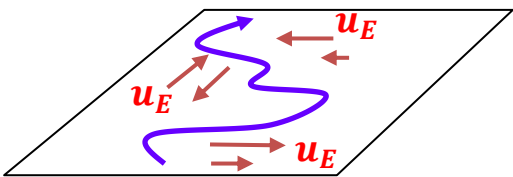
⇒ main interaction diagrams:



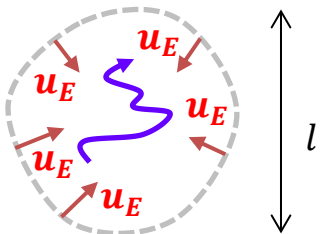
dynamic, curved field line
(aka curvature drift)



dynamic compression $\perp$ B field
(aka betatron)



turbulent sheared flow



large-scale compressions
(aka magnetic pumping)

close to original intuition by Fermi!
(cf Fermi 1949)

in large-amplitude turbulence $\delta B/B \gtrsim 1$, curvature + $\perp$ compression provides dominant contribution to acceleration due to strong non-gaussianity (intermittency) of random forces

Refs:

1. M.L. 19 [PRD 99, 083006 (2019)], M.L. 21 [PRD 104, 063020 (2021)], M.L. 22 [PRL 129, 215101 (2022)], M.L. 25 [PRE 112, 015205 (2025)]
based on previous work by G. Webb (1985, 1989)

Application to large-amplitude turbulence: $\delta B \sim B$

→ key features:

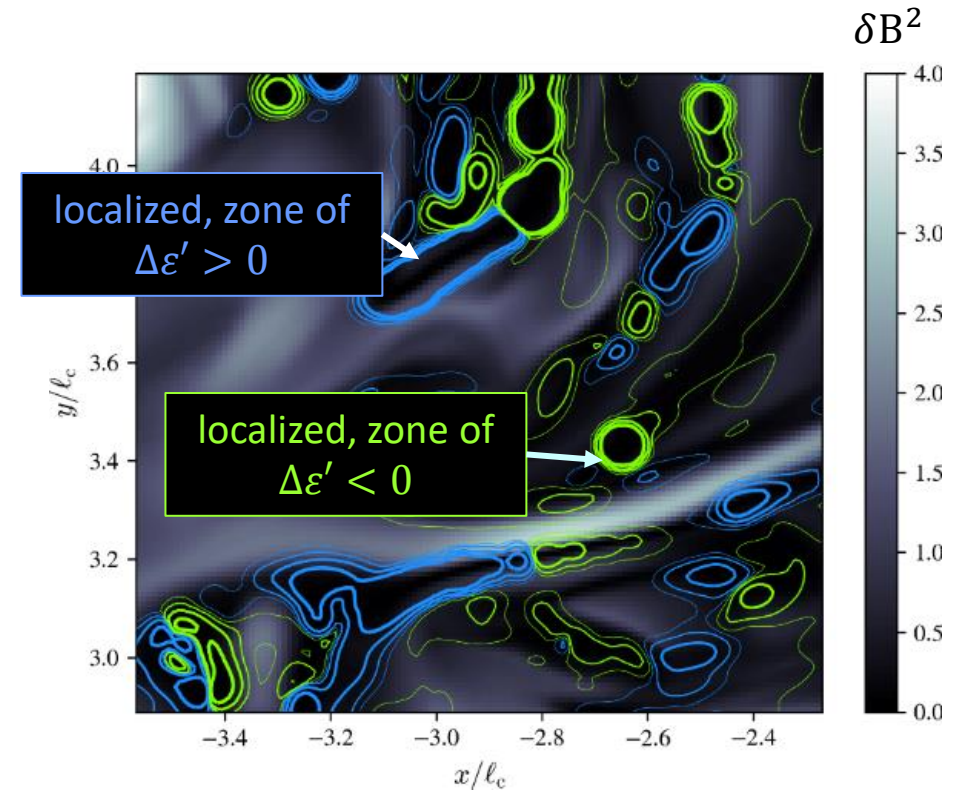
- ... energy loss/gain from localized regions of strong gradients of $\mathbf{u}_E$ (on some scale l)
- ... exponential energy loss/gain while in an accelerating/decelerating structure

$$\frac{d\varepsilon'}{dt'} \approx \pm \Gamma_l \varepsilon' \quad \text{with} \quad \Gamma_l \sim \mathcal{O}(\Delta u_E / l)$$

- ⇒ acceleration = non-Brownian diffusion, succession of intense episodes of gain/loss
- ⇒ diffusion coefficient = insufficient statistics for predicting particle energy distribution¹

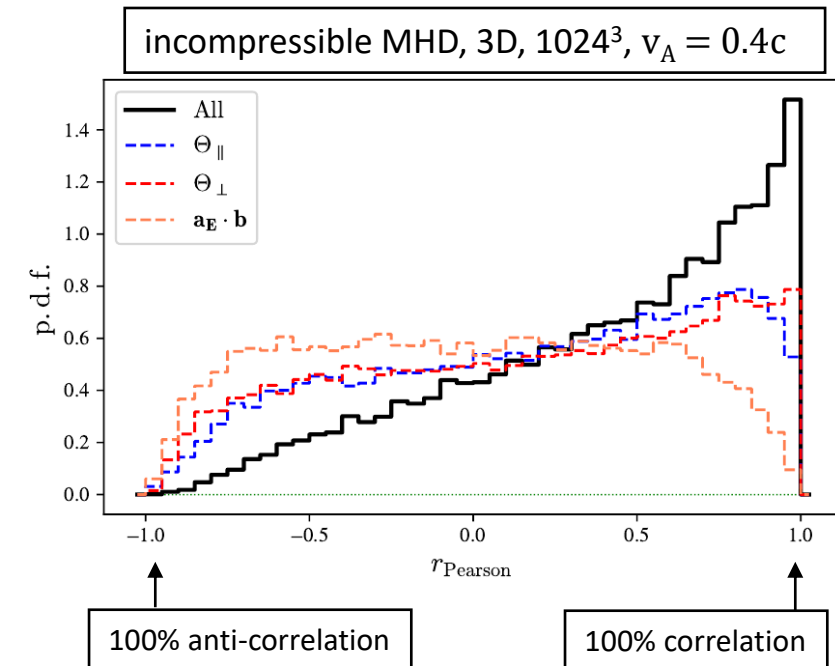
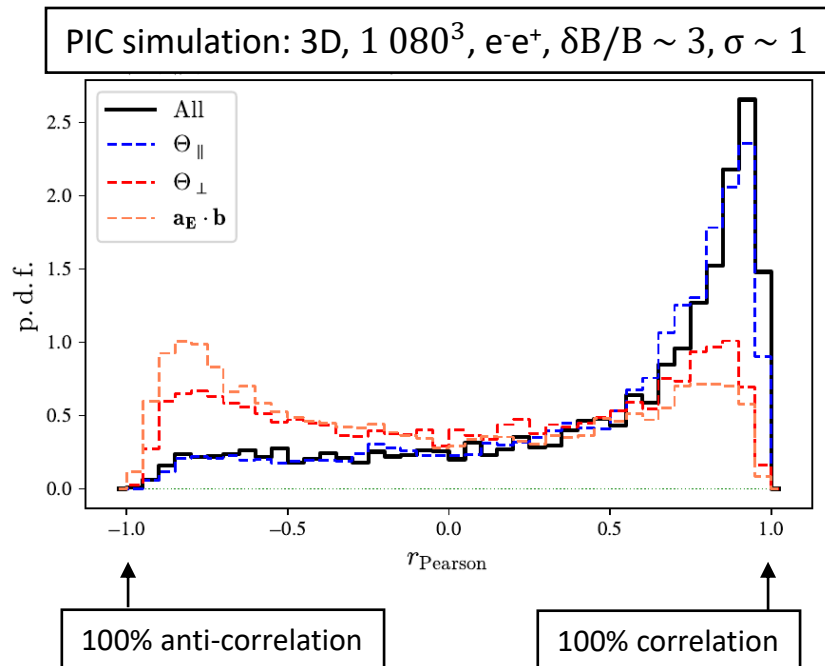
→ remarks:

- ... a generalized transport equation accounting for pdf(Γ_l): Ref.² + comparison to test particles in MHD sim.
- ... role of field line curvature in acceleration (+transport) supported by other studies³



Generalized Fermi acceleration: test on numerical experiments (PIC + MHD)

→ test¹: for each particle history in a simulation, reconstruct $\varepsilon'_{\text{th}}(t)$ using theoretical model and velocity gradients measured in the simulation at $\mathbf{x}, t$...
... then measure degree of correlation r_{Pearson} between the observed $\varepsilon'_{\text{obs}}(t)$ and reconstructed $\varepsilon'_{\text{th}}(t)$



[Note: $\Theta_{||}$ = field line curvature, $\Theta_{\perp}$ = perp. gradient (mirror), $\mathbf{a_E} \cdot \mathbf{b}$ = field line acceleration]

⇒ model captures the dominant contribution to particle energization

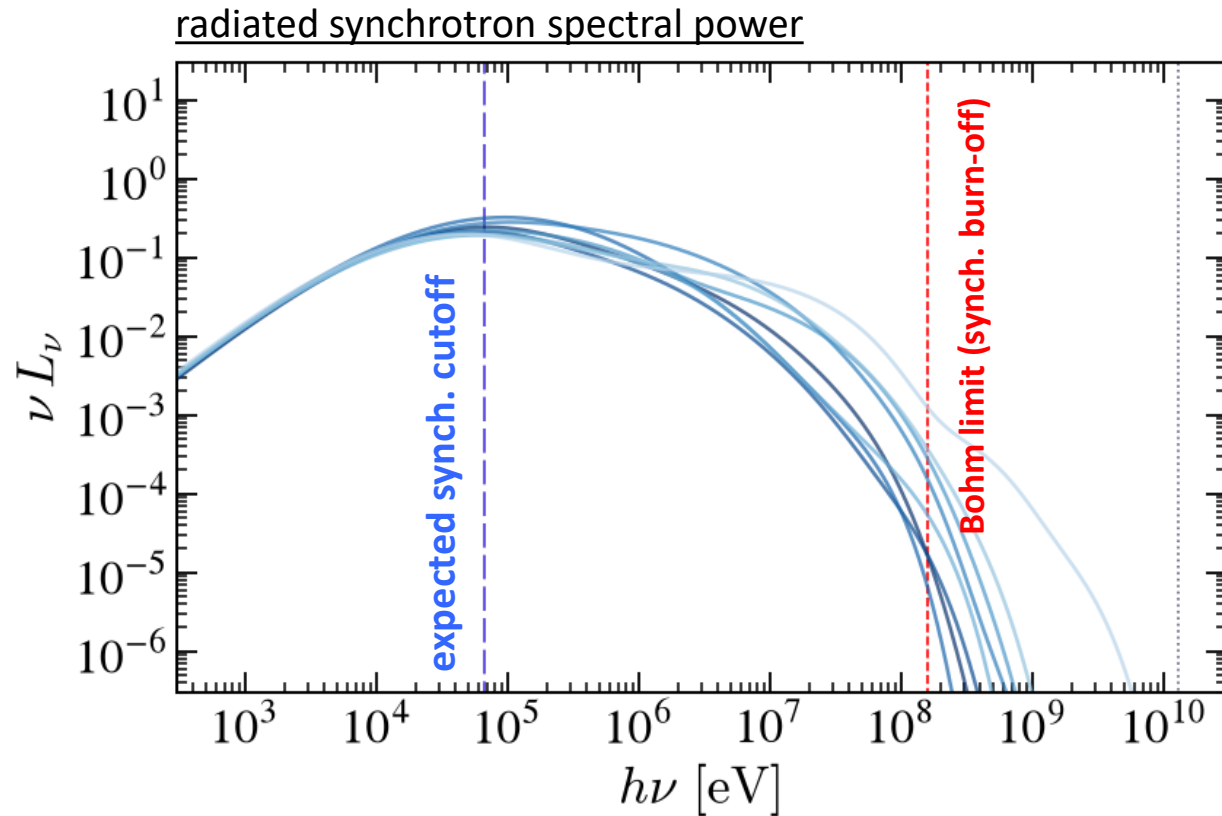
+ note: in wave turbulence w/ resonant wave-particle interactions, no apparent correlation seen (as expected)

Impact of energy losses: turbulence as an extreme particle accelerator?

→ maximum energy w/ synchrotron cooling:

... on average, turbulence is a slow accelerator¹: $D_{\varepsilon\varepsilon} \simeq 0.1 \sigma \varepsilon^2 c / \ell_c \Rightarrow t_{\text{acc}} \sim l_c / c$ (**turn-around time!**)

... vs Bohm limit: $t_B \sim r_g / c \ll t_{\text{acc}} \Rightarrow$ expect sharp cut-off ε_{cut} well below Bohm limit?



... intermittency → particles are accelerated at different rates in different regions²

⇒ extension of spectrum beyond expected ε_{cut} , up to Bohm limit³

... ε_{cut} particles: exponential acc. in gradient $\Delta u_E \sim c$ on scale r_g (generalized Fermi type)

+ characteristic features: steepening of spectrum, strong time variability at radiative limit ($\sim 100\text{MeV}$)

→ a specific radiative signature of relativistic, large-amplitude turbulence

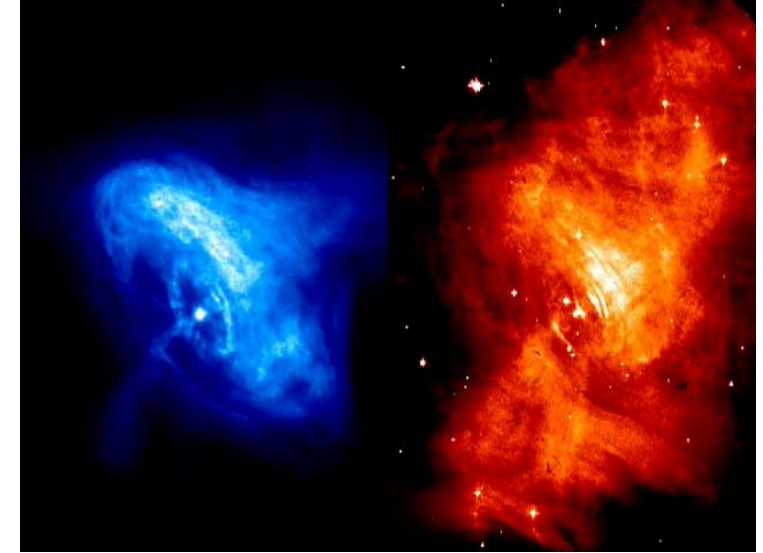
Turbulence in the Crab nebula as an extreme particle accelerator

→ application to the Crab nebula, and its flaring behavior?

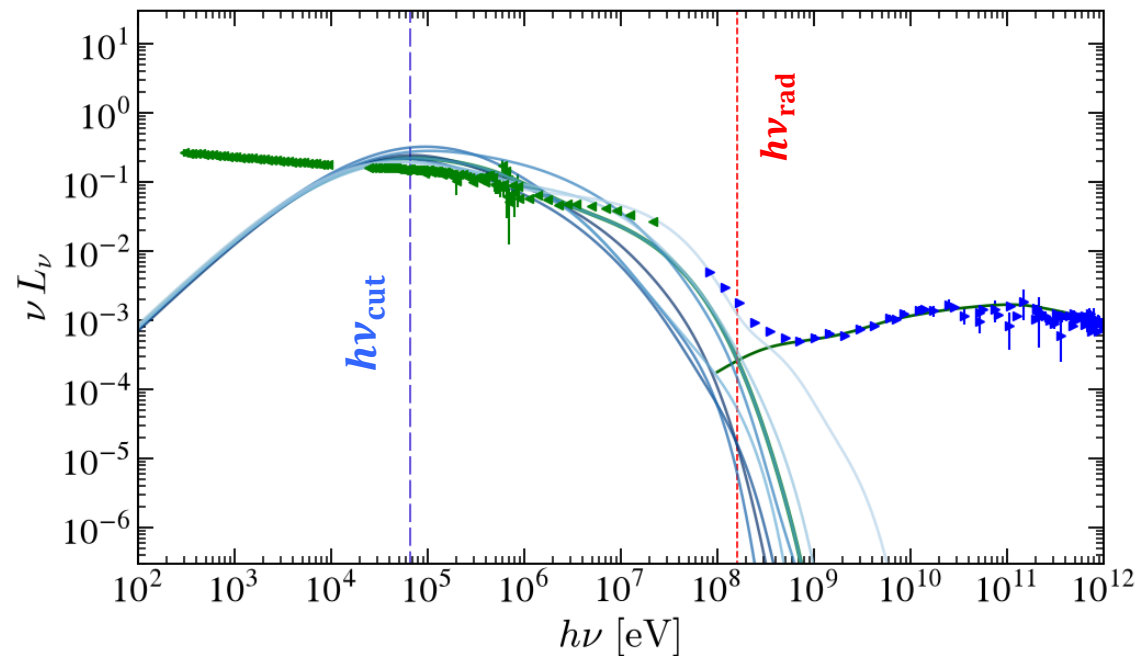
... assuming a region of relativistic turbulence downstream of the termination shock, e.g.¹, reasonable estimates indicate that turbulence can produce PeV particles:

$$B \sim 200\mu\text{G}, \ell_c \sim 0.1\text{pc}, \sigma \sim 1$$

$$\Rightarrow h\nu_{\text{cut}} \sim 100\text{keV}, \quad h\nu_{\text{rad}} \sim 100\text{MeV}$$



direct overlay (no tuning beyond flux norm.) of simulation to Crab



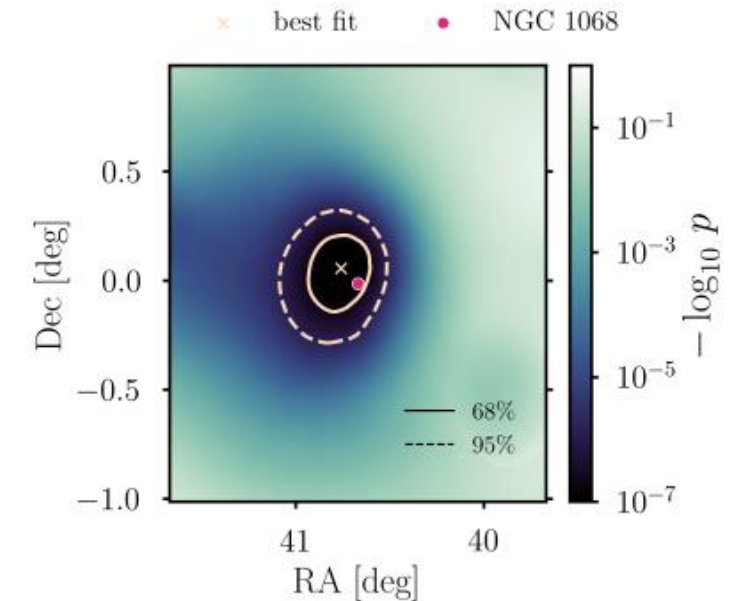
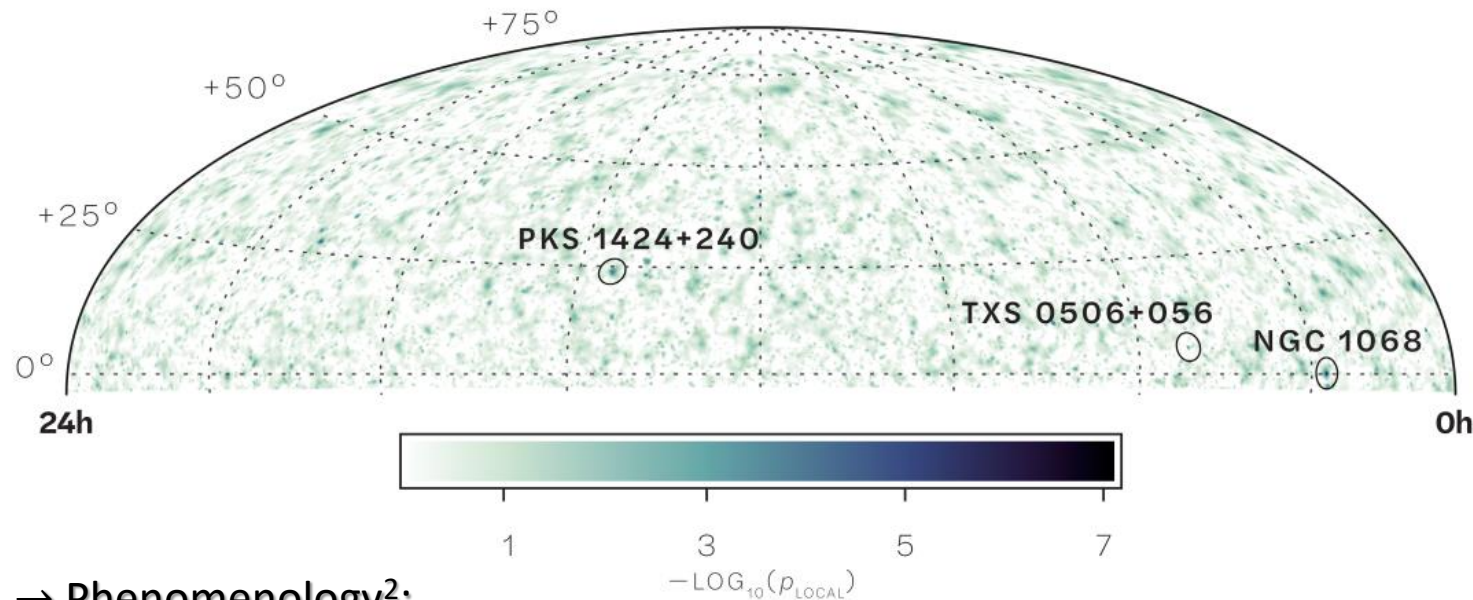
→ additional motivations:

... Crab flux at 100MeV is variable in time² → here interpreted as fluctuating energy cut-off

... Crab flares: rare events of stochastic acceleration at/above synchrotron burn-off?

Application to BH coronae: stochastic proton acceleration & IceCube neutrinos

→ IceCube 22... 26¹: excess of high-energy (3 – 30 TeV) neutrinos from nearby AGN (Seyfert) NGC 1068... and possibly NGC4151, NGC7469, CGCG 420-015



→ Phenomenology²:

... origin of neutrinos: proton acceleration to $\gtrsim 30\text{TeV}$ + hadronic interactions

... lack of (comparable) gamma-ray flux requires compact zone for reprocessing through $\gamma\gamma \rightarrow e^+e^-$

⇒ likely region: BH corona, within $\lesssim 30r_g$

⇒ likely acceleration agent: strong, semi-relativistic turbulence ($\sigma \sim 0.1$) energizing electrons, accelerating protons

i.e., IceCube opening a window on hadronic counterpart of leptonic dissipation channels seen in X-ray?

IceCube neutrinos = window on hadronic processes in black hole coronae

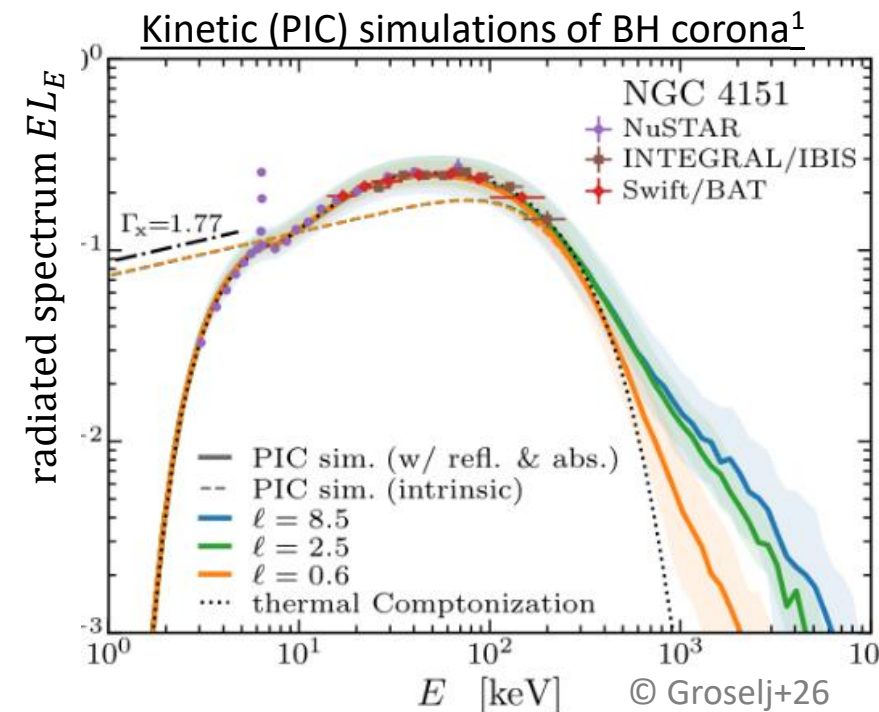
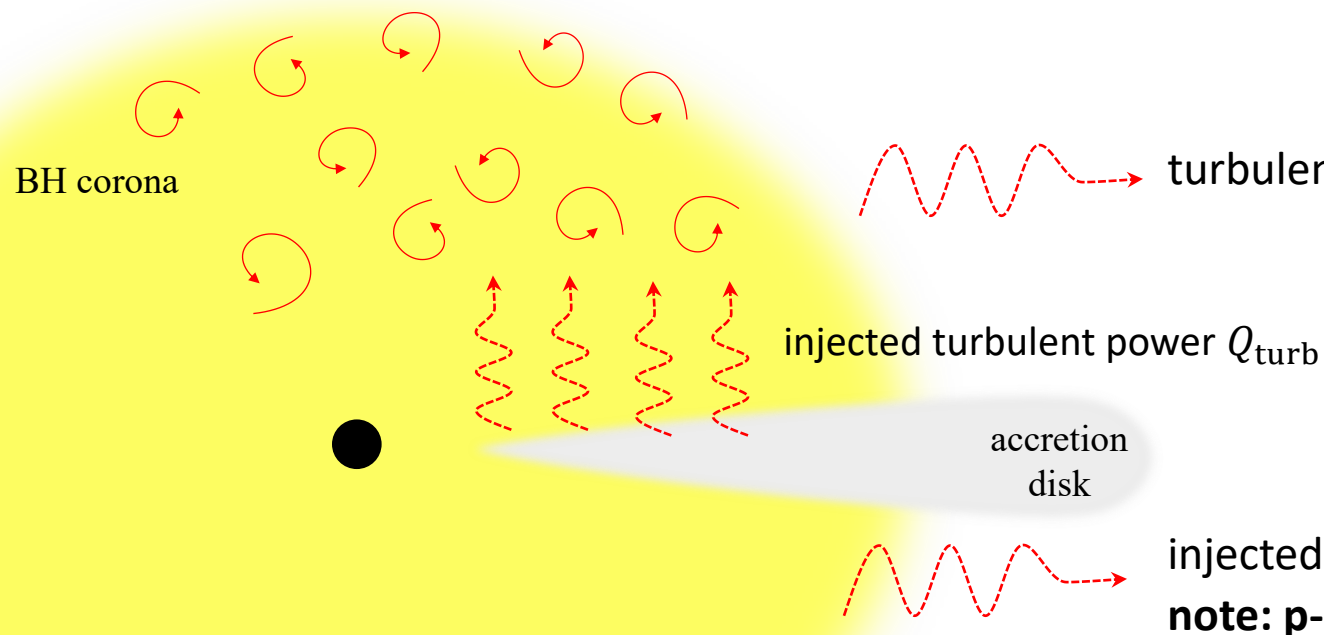
→ corona: $B \sim O(\text{kG}), v_A \sim O(0.3c), R_c \sim O(10 r_g)$

... electron energization through turbulent heating:

$kT_e \sim 100\text{keV} \rightarrow$ hard X-ray continuum by Comptonization
on short times scales $\ll R_c/c$!

recent first-principles modelling through PIC simulations¹

... proton acceleration: long timescales $\gg R_c/c$ (beyond PIC!)

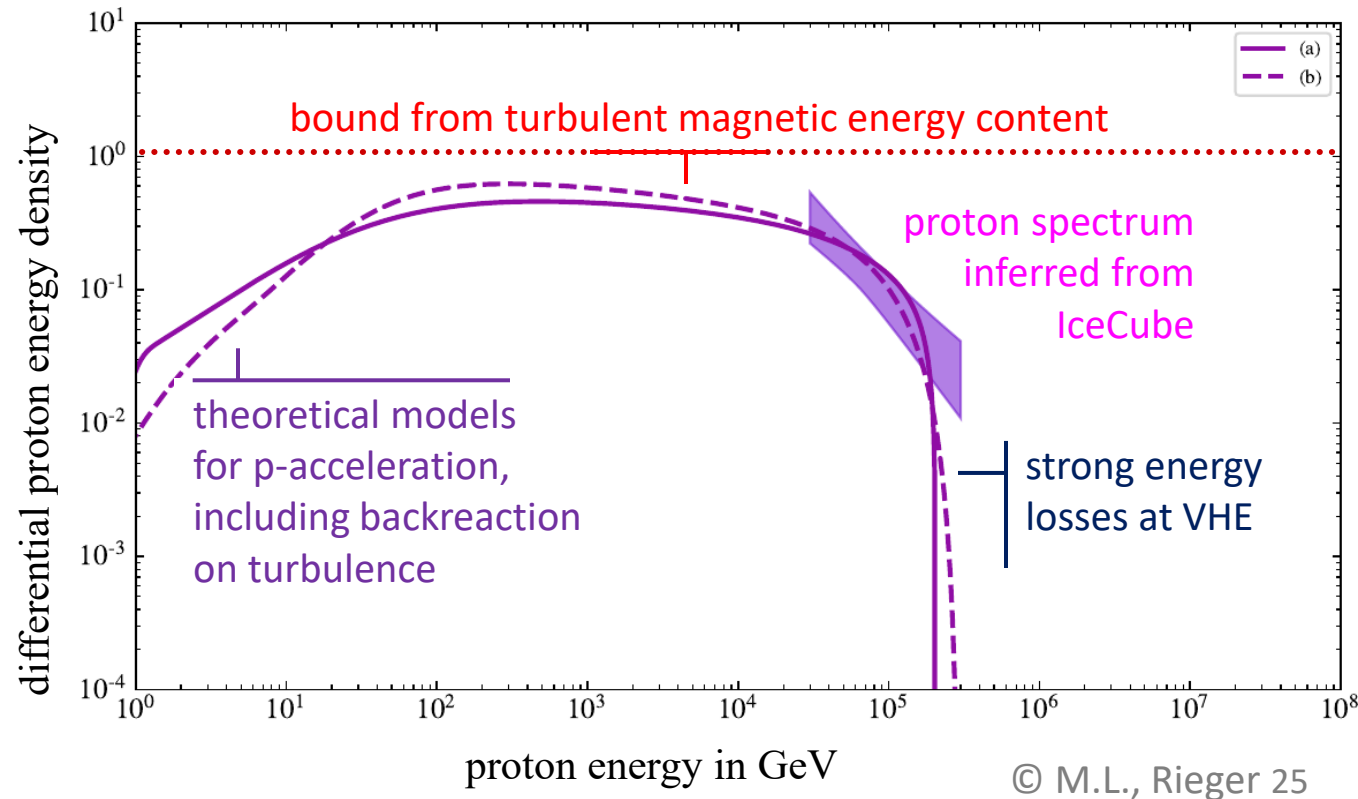


Nonlinear stochastic acceleration: p energy budget dictated by turbulence driving

→ energy budget:

$$\dots \text{ for NGC1068: } L_\nu \sim 10^{42} \text{ erg/s} \sim 10^{-3} L_{\text{edd}} \Rightarrow u_p \sim L_\nu t_{\text{loss}}/V_{\text{cor}} \sim O(10^5) \text{ ergs/cm}^3 \sim O(u_B) \sim O(p_{th})$$

⇒ a natural explanation for observed L_ν : proton energy fraction limited by rate at which turbulence is driven ^{1,2}
(then, damping of turbulence self-regulates acceleration process, shaping flat energy spectrum¹)



Detailed modeling of hadronic processes in AGN cores by Sébastien Le Bihan et al.

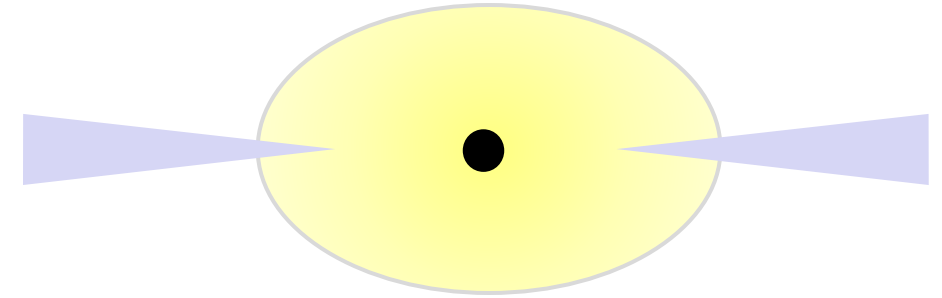
→ Simplified model:

spherical, size $R_{co} = 15 r_g$

... note, key parameters are:

advection through the corona (here, std radial inflow)

diffusive escape from corona (here, std diffusion coefficient)



→ Transport equations:

$$\partial_t n_s(p, t) = [\mathcal{L}_{\text{stoch. acc.}} n_p(p, t)]_{s=p} + -\partial_\varepsilon (\varepsilon \nu_{\text{loss}} n_s) - \frac{n_s}{\tau_{\text{esc}}} + Q_s$$

... for species: protons (s=p), pairs, photons, pions, neutrinos

... describing stochastic acceleration of protons, electromagnetic cascades, escape processes, Comptonization

$$\partial_t \mathcal{E}_{B_k}(t) = \text{cascade}(k) - \text{acceleration}(k) - \text{heating}$$

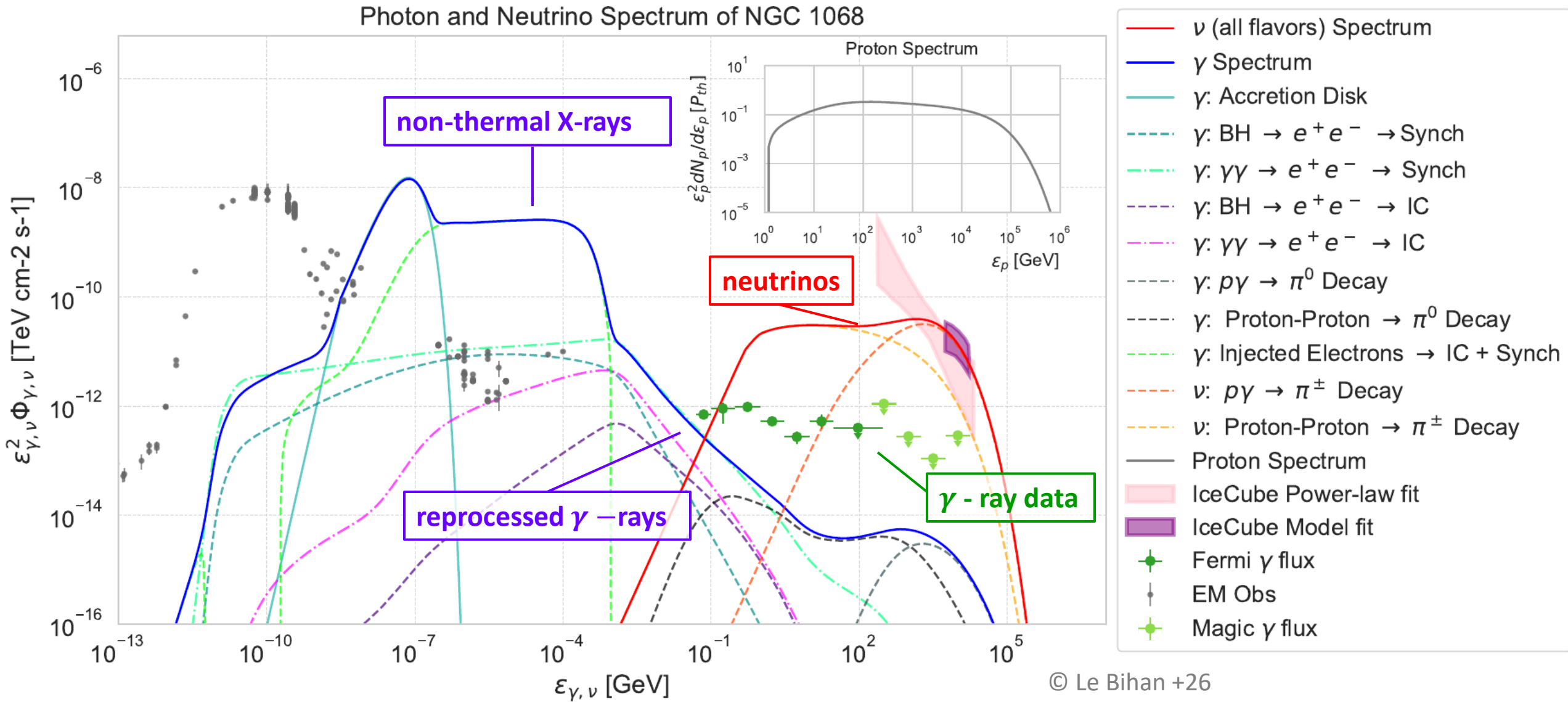
... describing evolution of turbulent magnetic energy in k-space, including damping from accelerated protons

→ Numerical code (developed by S. Le Bihan):

Turb-AM3¹: couples stochastic acceleration solver + turbulence feedback with AM3 (multi-messenger radiative code²)

... NB: applicable to various other types of sources!

Detailed model (S. Le Bihan et al.): multi-messenger (stationary) spectrum for NGC1068



⇒ standard coronal + turbulence parameters provide a natural fit to IceCube (+photon) data

Summary

- Turbulent plasmas as a modern particle acceleration scheme in VHE multi-messenger astrophysics
 - ... the Fermi picture is well alive: a generalized Fermi effective theory for modelling acceleration*
 - ... main energization agents (large-amplitude): sharp field line bends, compression $\perp B$, acceleration of field lines*
 - ... satisfactory tests on simulations*
- Many open questions on microphysics of particle acceleration:
 - ... extension to low $\delta B/B$ and role of intermittency, structures, waves?*
 - ... extension to low v_A (requires long timescales)?*
 - ... impact of turbulence anisotropy at gyroradius $\ll$ outer scale?*
- Signatures and applications to VHE multi-messenger astrophysics:
 - ... w/ synchrotron cooling: extension beyond expected cut-off due to non-homogeneity of acceleration rates + acceleration up to synchrotron burn-off: an origin for Crab flares?*
 - ... on long timescales, in BH coronae: p-spectra self-regulated by feedback on turbulence, provides natural account of observed IceCube neutrino flux*