

Plasma-driven superradiant instabilities

Numerical investigation of the stability of accreting Kerr black holes

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Black holes, gravitational waves and fundamental physics
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Superradiance with photons

The question:

Superradiance triggered by plasma surrounding astrophysical black holes?

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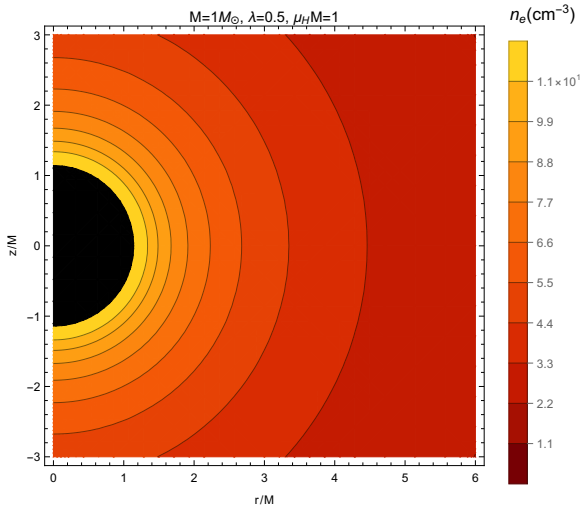
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¹ J. Conlon, C. Herdeiro, Phys. Lett. B780 (2018) 169-173

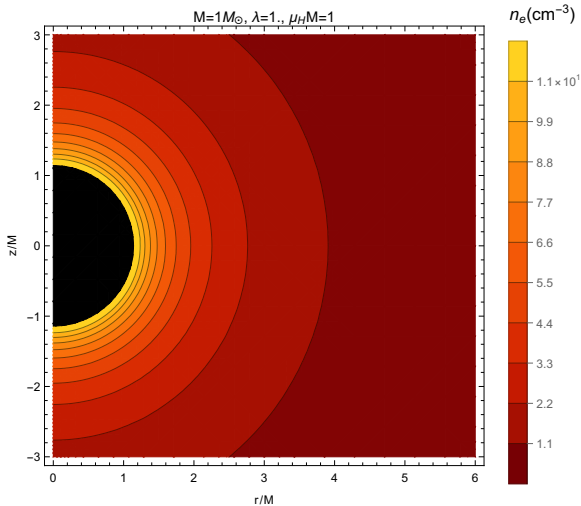
Density profiles

Spherically symmetric density distribution (Bondi accretion)



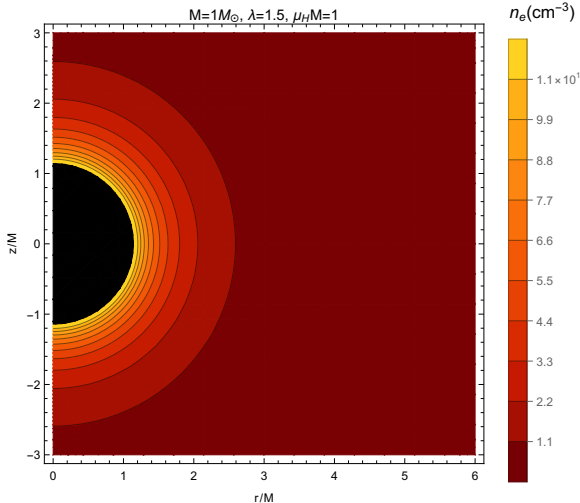
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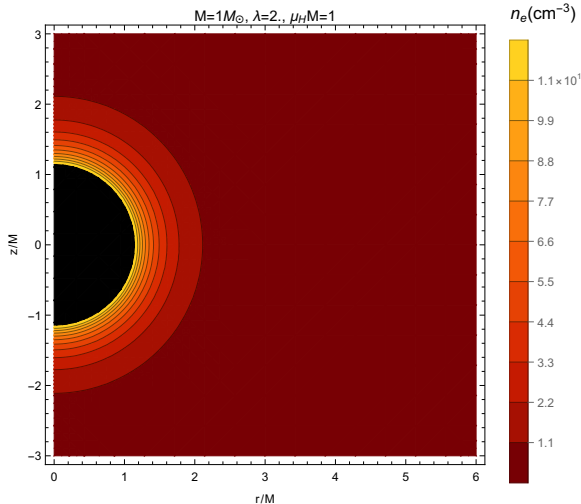
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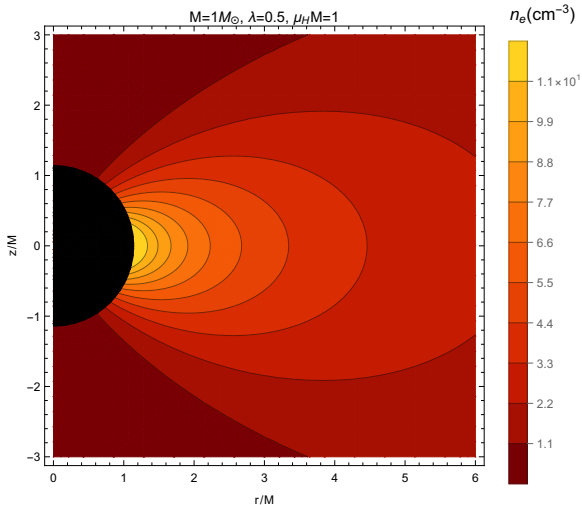
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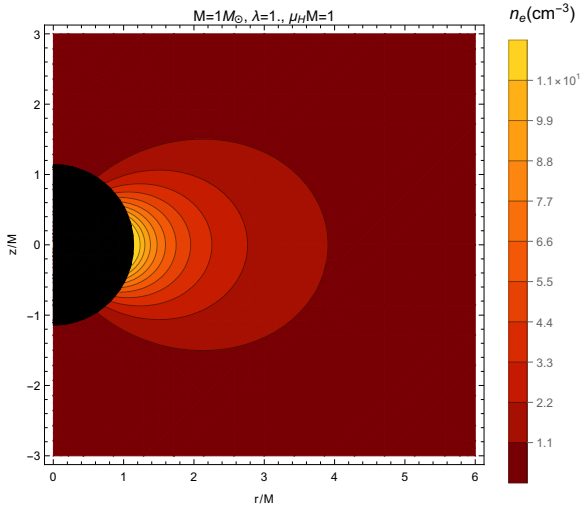
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Axisymmetric density distribution (ADAF I)



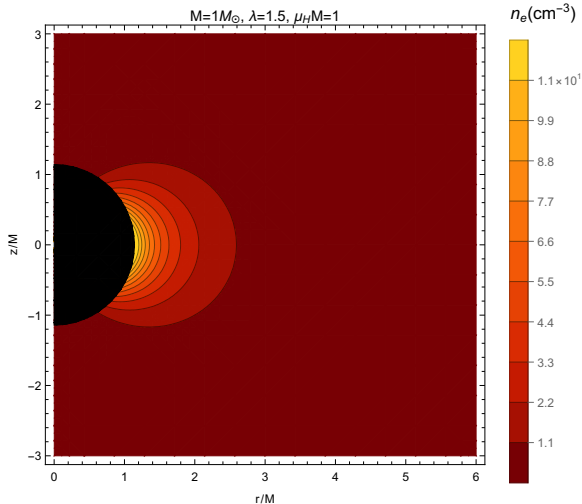
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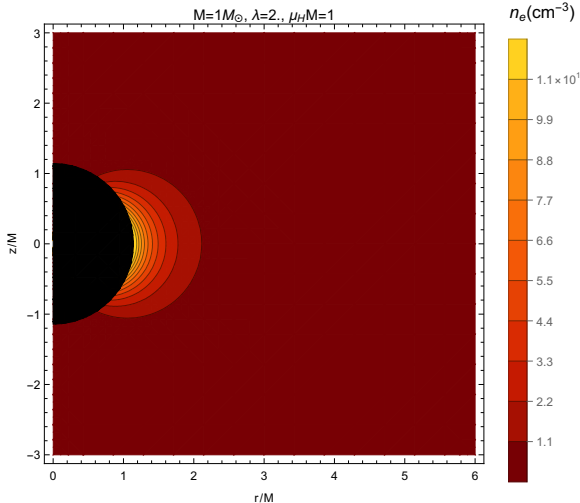
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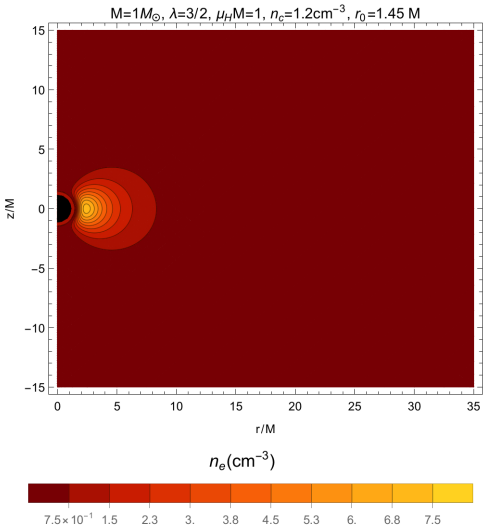
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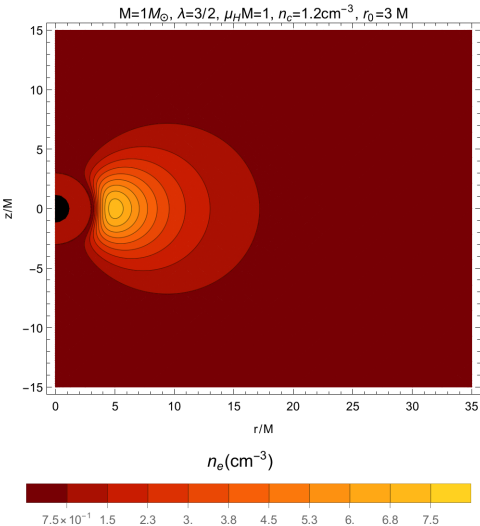
Density profiles

Distribution with inner edge and “corona” (ADAF II)



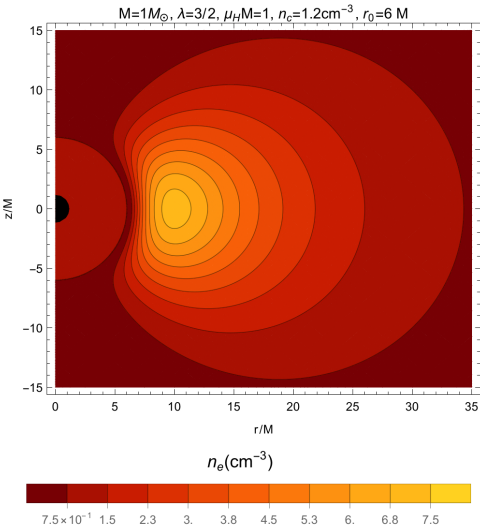
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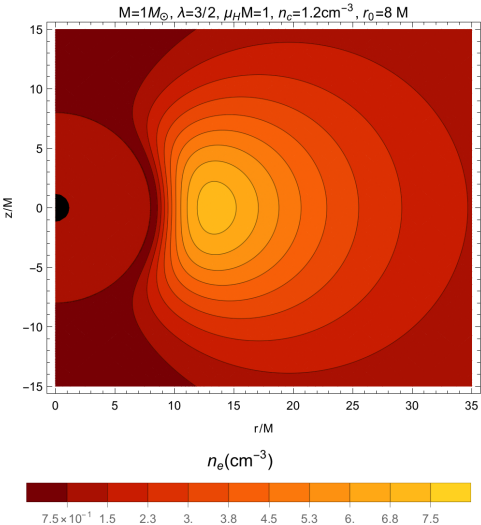
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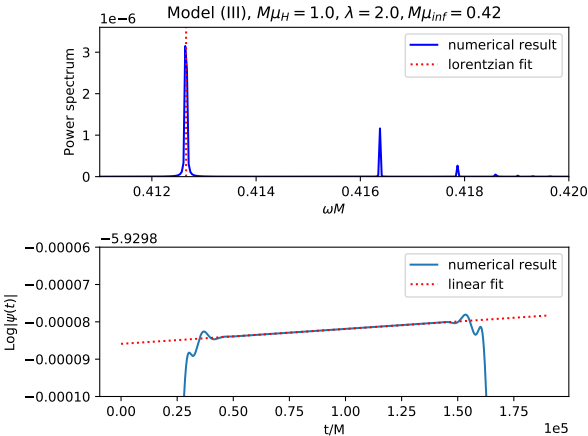
Numerical setup

Time evolution of Klein-Gordon equation on Kerr spacetime with method of lines

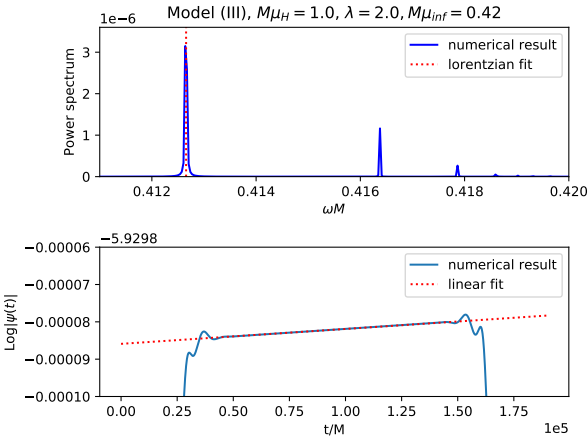
- Decomposition: spherical harmonics
- Stepper algorithm: 4th order Runge-Kutta (RK4)
- Boundary Conditions:
- At infinity: outgoing BC
- Horizon: ingoing BC + artificial damping region (PML)
- Tests: Convergence, E/J Conservation, Kerr QNMs, BH-bomb instabilities, constant mass QBSs, separable mass functions²

²V. Cardoso, I. P. Carucci, P. Pani, and T. P. Sotiriou, Phys. Rev. D 88, 044056 (2013)

Results: Bondi

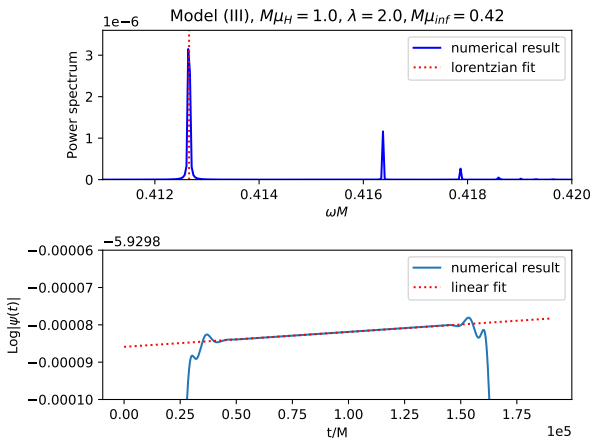


Results: Bondi



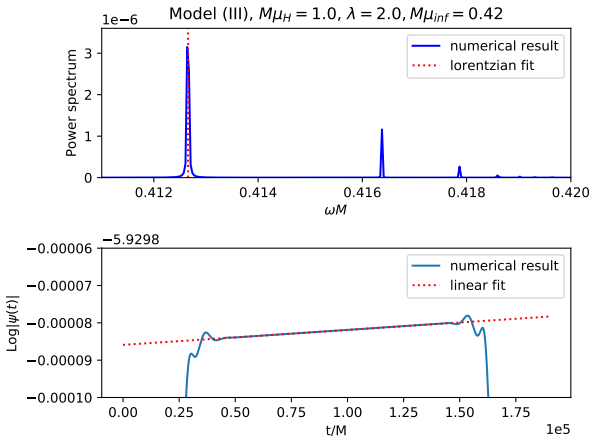
- $M = 10M_{\odot}$, $n_H \sim 0.1 \text{ cm}^{-3}$

Results: Bondi



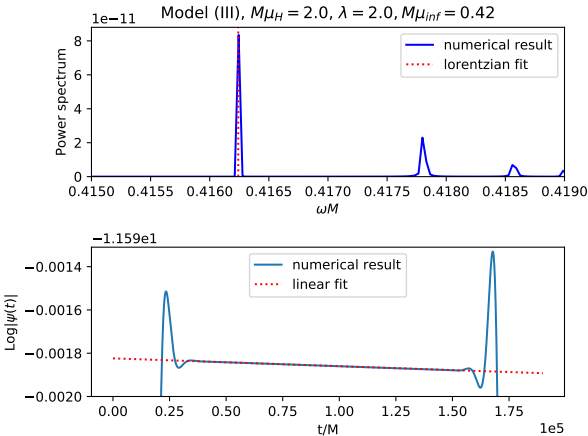
- $M = 10M_{\odot}$, $n_H \sim 0.1 \text{ cm}^{-3}$
- $\omega M = 0.413 + i 4 \cdot 10^{-11}$

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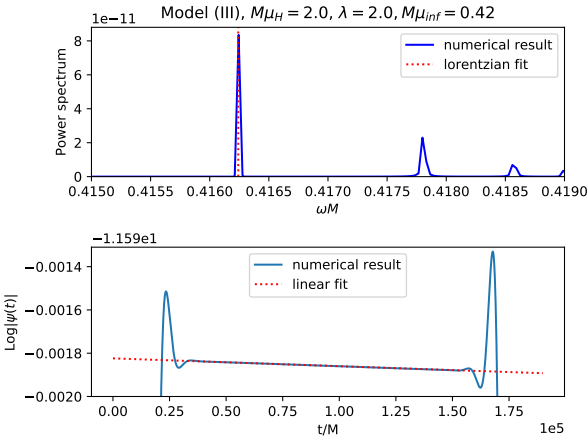


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- $\omega M = 0.413 + i 4 \cdot 10^{-11}$
- **INSTABILITY!**

Results: Bondi

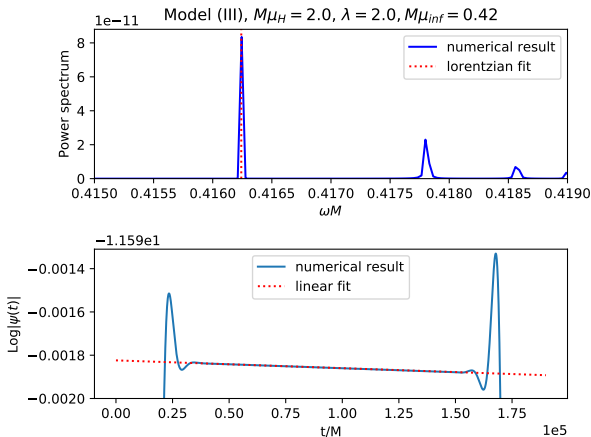


Results: Bondi



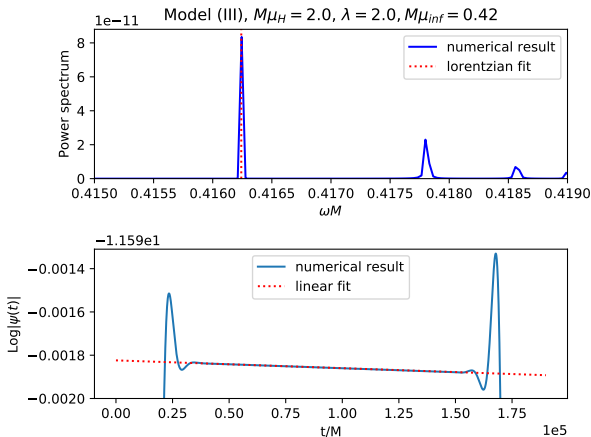
- $n_H \sim 0.5 \text{ cm}^{-3}$

Results: Bondi



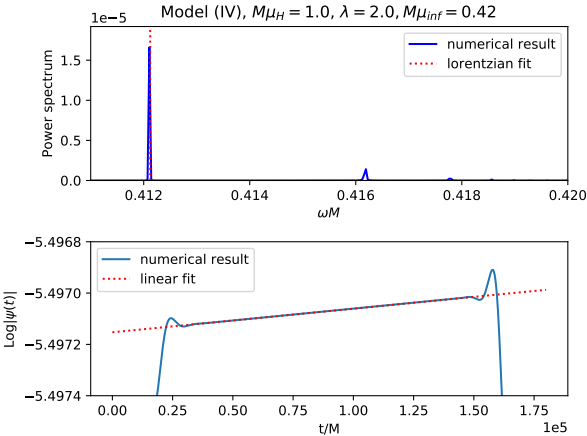
- $n_H \sim 0.5 \text{ cm}^{-3}$
- $\omega M = 0.416 - i 3.6 \cdot 10^{-10}$

Results: Bondi

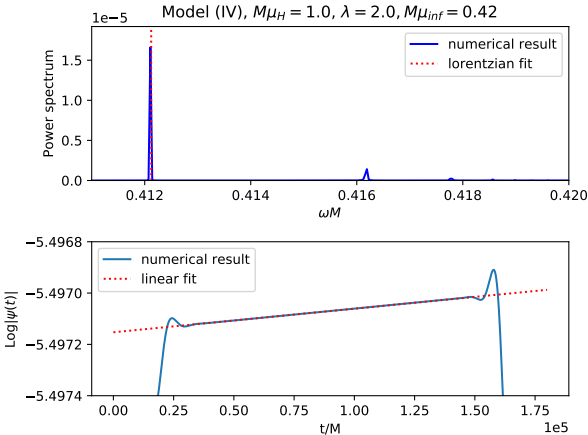


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Results: ADAF I

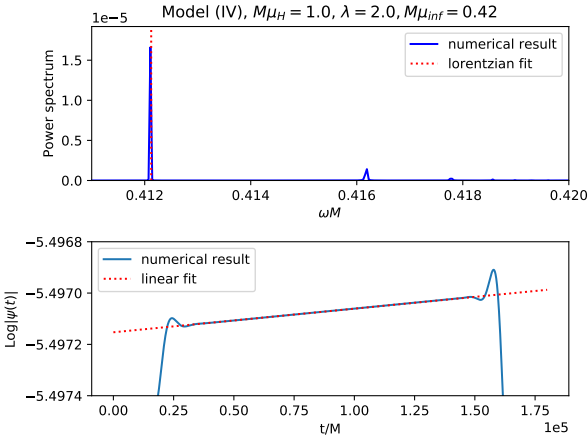


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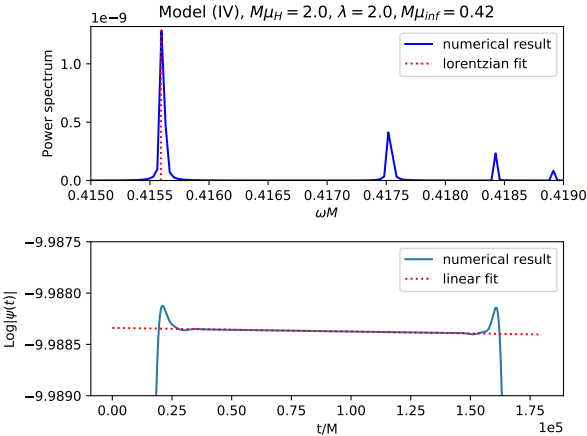
- $\omega M = 0.412 + i 9.2 \cdot 10^{-10}$

Results: ADAF I

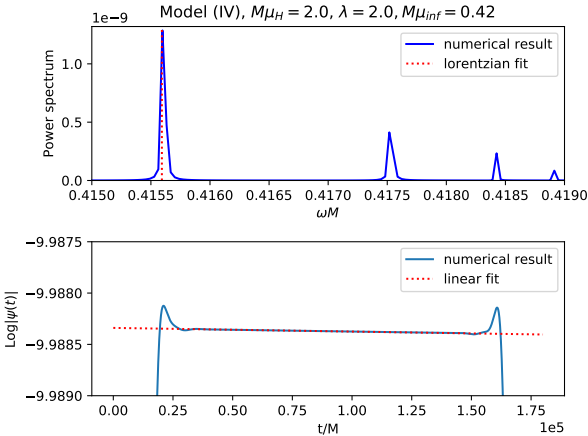


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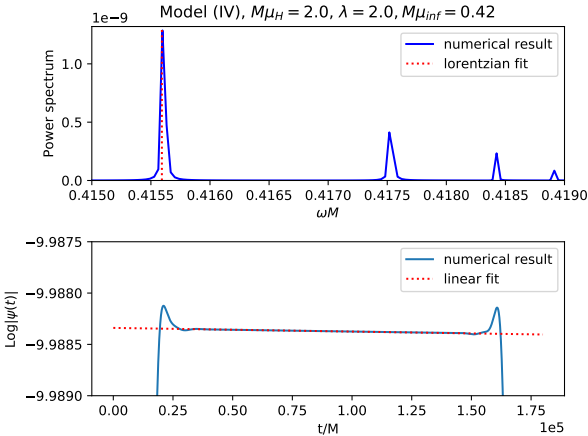


Results: ADAF I



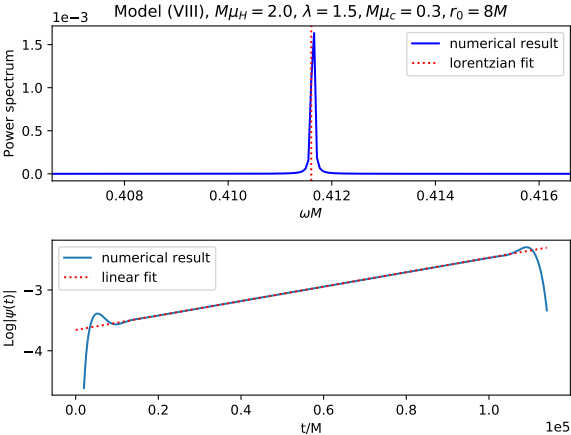
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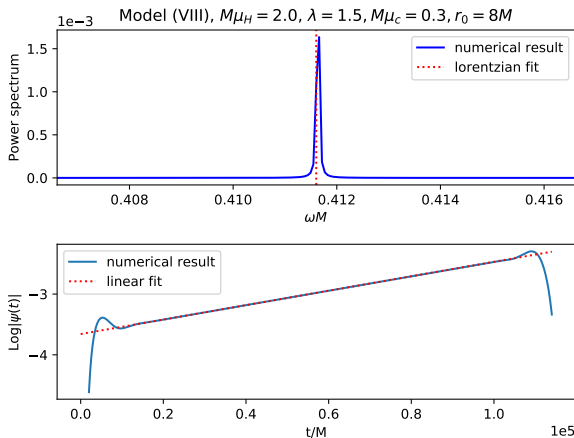


- $\omega M = 0.416 - i \, 3.6 \cdot 10^{-10}$
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Results: ADAF II

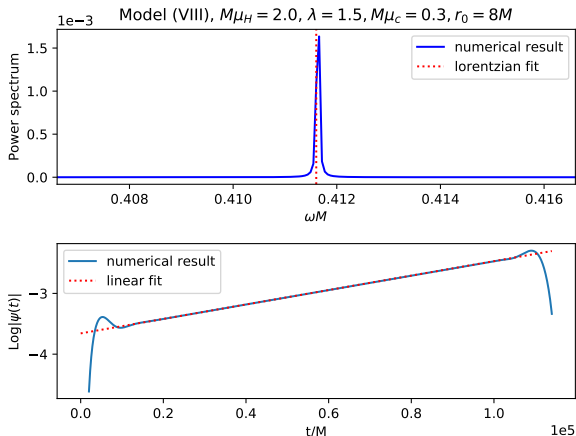


Results: ADAF II



$$M = 10M_{\odot}, n_c \simeq 0.01 \text{ cm}^{-3}: \quad \omega M = 0.412 + i 1.2 \cdot 10^{-5} !$$

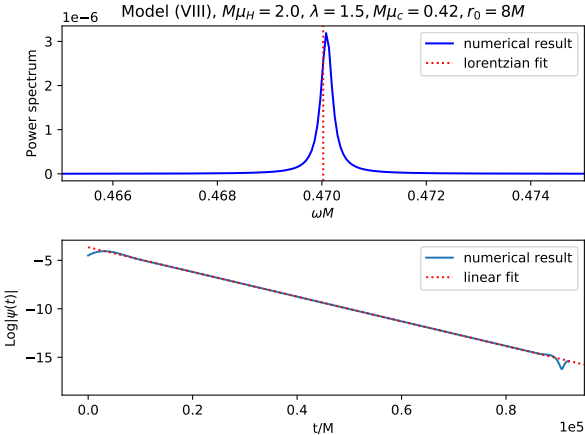
Results: ADAF II



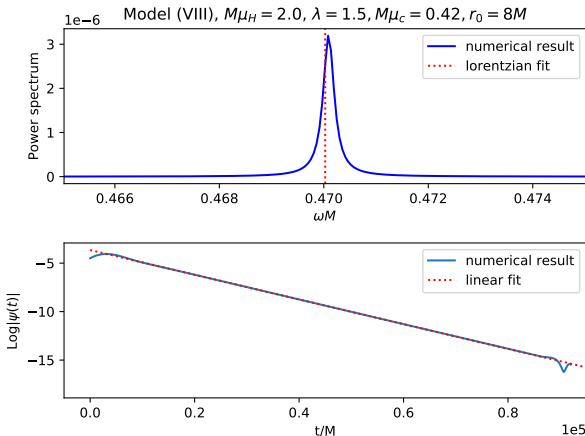
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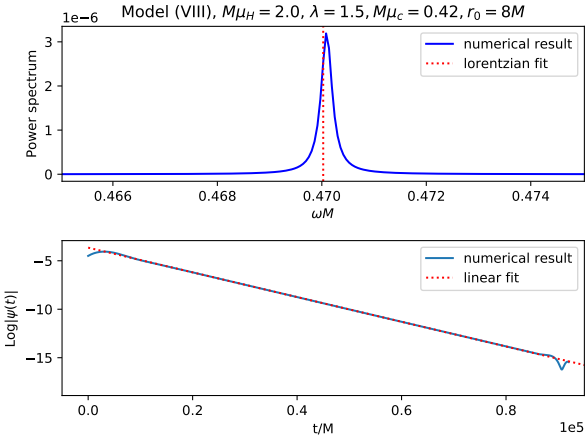
Results: ADAF II



Increase to $n_c \simeq 0.02 \text{ cm}^{-3}$:

$$\omega M = 0.471 - i 1.4 \cdot 10^{-4} !$$

Results: ADAF II



Increase to $n_c \simeq 0.02 \text{ cm}^{-3}$: $\omega M = 0.471 - i 1.4 \cdot 10^{-4}$!

STABLE!

Conclusions

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- Plasma-driven superradiant instabilities: found, but fragile...
- ...highly sensitive to effective mass in ergoregion!
- Likely to be suppressed by realistic densities in the ergoregion ($n_e \gg 1 \text{ cm}^{-3}$)

THANK YOU FOR YOUR ATTENTION!

BONUS SLIDES

Tests

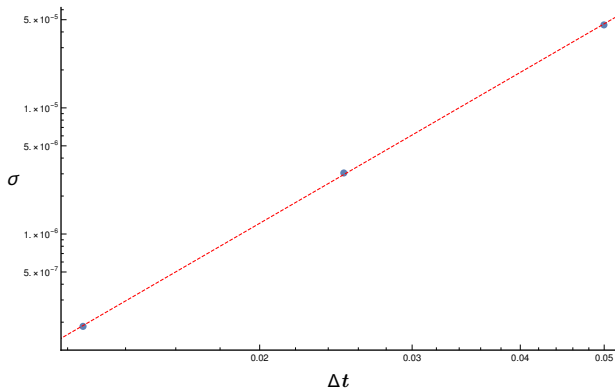
- Convergence test: x - and t -error scaling: check!
- Conserved quantities test: check!
- Comparison with original 1+1D code³: check!
- Comparison with approximated analytical formulas⁴: check!
- Comparison with numerical ω -domain⁵ results: check!

³S. Dolan

⁴V. Cardoso et al., PRD 70,044039 (2004)

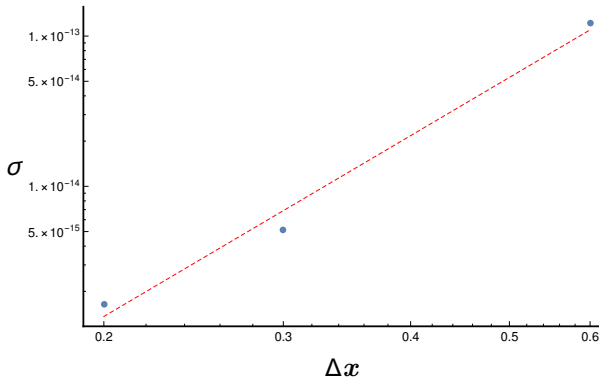
⁵E. Berti, <https://pages.jh.edu/~eberti2/ringdown/>

Numerical convergence: Δt resolution



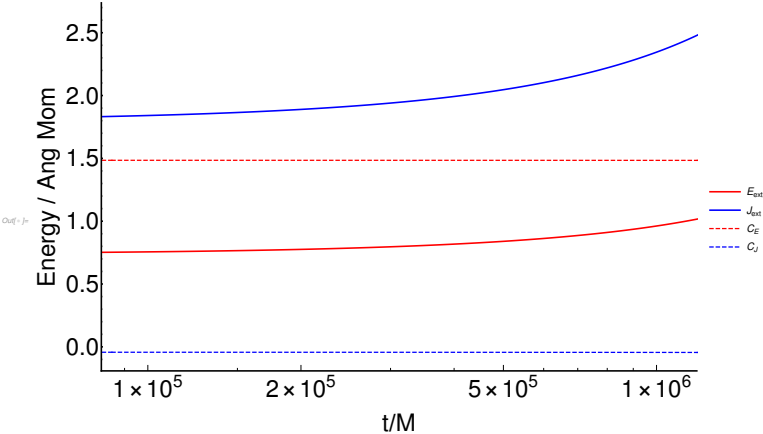
$$\log \sigma = (1.93 \pm 0.15) + (3.97 \pm 0.04) \log \Delta t$$

Numerical convergence: Δx resolution

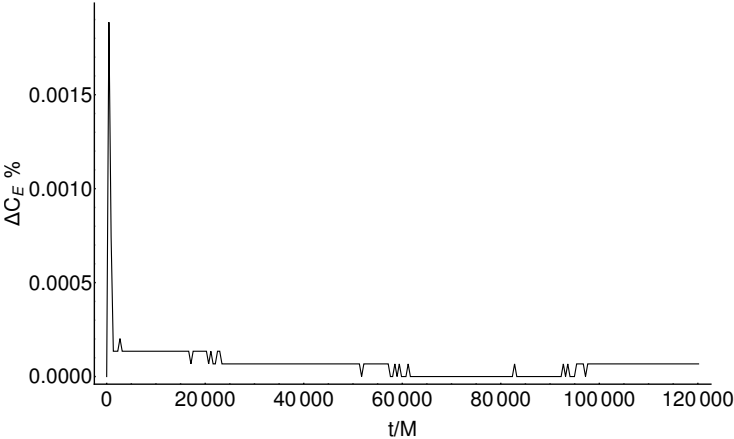


$$\log \sigma = (-27.80 \pm 0.56) + (3.99 \pm 0.47) \log \Delta x$$

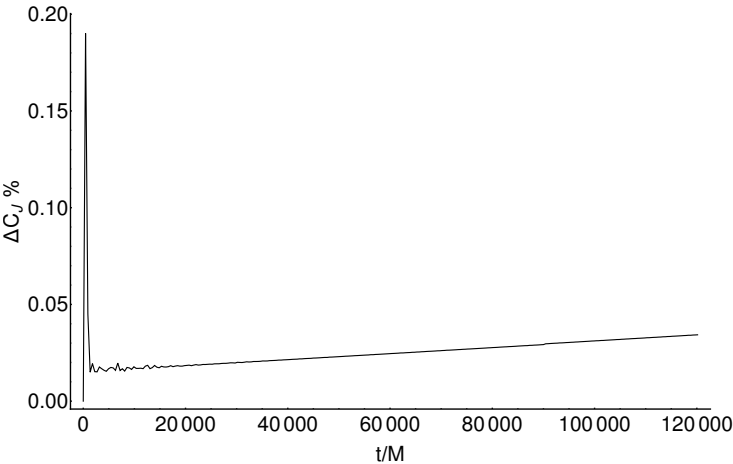
E/J conservation



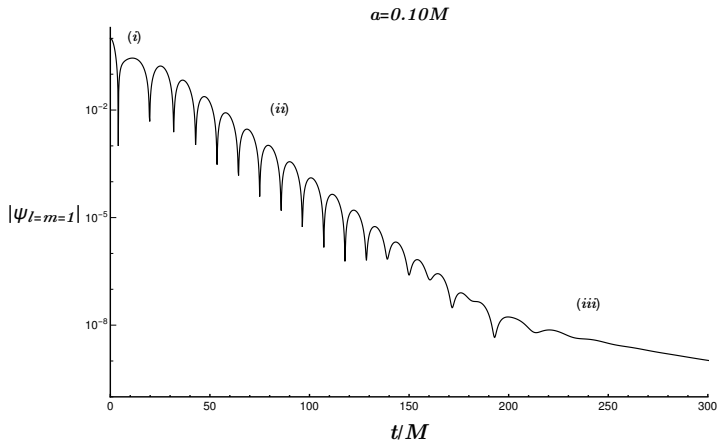
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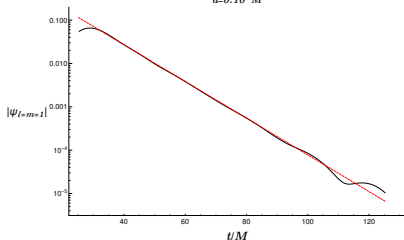
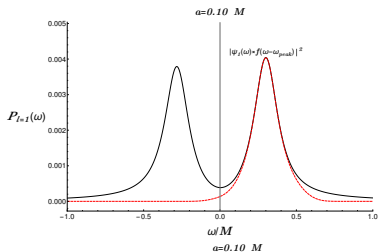


QNM



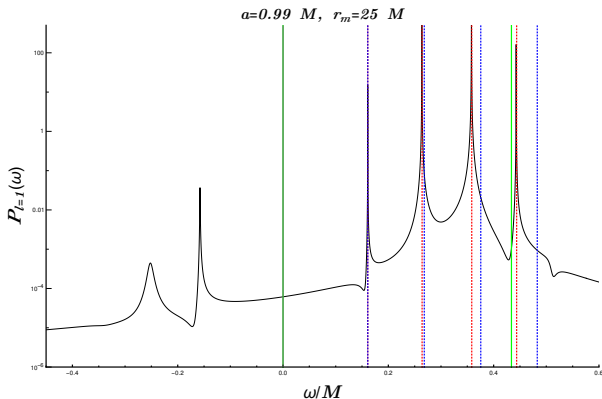
- ❶ Prompt-response: depends on the specific initial data
- ❷ Ringdown: quasinormal modes
- ❸ Power-law tail

QNM test



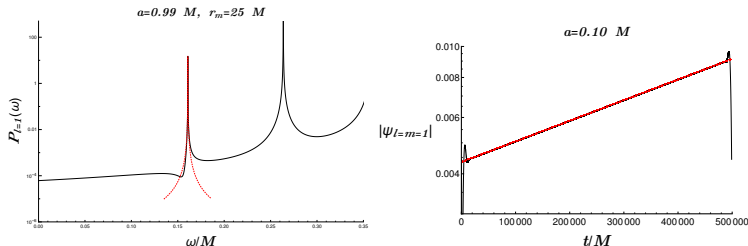
- Peak finding
- Fit with Lorentzian: $\frac{(\pi\nu)^{-1}}{(\omega - \omega_{peak})^2 + \nu^2} \rightarrow \hat{\omega} = \omega_{peak}$
- Filter: $f(\omega - \hat{\omega}) = \exp(-\frac{(\omega - \hat{\omega})^4}{(nd\omega)^4})$
- linear fit: $\log |\psi| = q + \nu_{peak} \cdot t \rightarrow \nu = \nu_{peak}$

Mirror test



- green: superradiant regime
- red: results from reference 1+1D code
- blue: approximated formula $\omega_{nlm} \simeq \frac{(n+\frac{\ell}{2}+1)}{r_{*,\text{mirror}}}$

Extraction of superradiant modes



$$\omega_{l=m=1} = 0.160877(\pm 0.03\%) + i7.42416 \cdot 10^{-6}(\pm 0.2\%)^6$$

⁶relative errors w.r.t. results of Dolan(2013), in agreement with Cardoso et al.(2004) analytical and numerical results.

Superradiance

Superradiance with rotating BHs

Classical extraction of the BH's rotational energy in favor of bosonic degrees of freedom

Superradiant scattering

Waves coming from infinity scattering off rotating BHs get amplified if

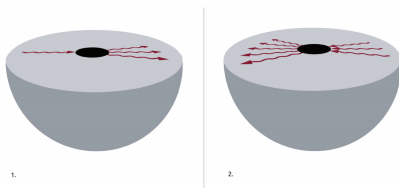
$$0 < \omega \leq m\Omega_H \implies |\mathcal{R}|^2 = |\mathcal{I}|^2 - \frac{\omega - m\Omega_H}{\omega} |\mathcal{T}|^2 \geq |\mathcal{I}|^2$$

Superradiant instabilities

Confinement leads to instability

Confined superradiant modes are forced to be repeatedly amplified, resulting in an exponential growth: **instability**

- Black-hole bomb: BH enclosed by a mirror



- Mass term, μ : natural confinement of modes with $\omega \leq \mu$

$$\Psi(r \rightarrow \infty) \sim \frac{e^{-qr}}{r} \quad q = \sqrt{\mu^2 - \omega^2}, \operatorname{Re}(q) \geq 0$$

Spectrum of the perturbations

How to study the (in)stability?

Solve for characteristic oscillations:

$$(\square - \mu^2)\psi = 0 \rightarrow \psi_{nlm}$$

$$(a, \mu M) \rightarrow \{\omega_{nlm}, \nu_{nlm}\}$$

Quasi-bound states

$$\psi_{QBS}(t) \sim e^{-i(\omega + i\nu)t}$$

for $\nu > 0$, instability with associated growth rate $\tau = \frac{1}{\nu}$

Frequency-domain vs Time-domain approach

- Analytical approximations (Detweiler, 1980; Zouros & Eardley, 1979; eikonal approximation; Pösche-Teller potential)
- Numerical ω -domain techniques (WKB, shooting, continued fraction method)

Non-separability requires time-domain approach

Space-time dependent mass term in perturbation equation ruins separability, on which these techniques rely

Effective mass from non-linear plasma effects

In unmagnetized, cold plasma: charges (e^-) oscillate with characteristic frequency:

Photon propagation in plasma

$$\omega^2 = k^2 + \omega_p^2, \quad \omega_p \equiv \sqrt{\frac{4\pi\alpha}{m_e} n_e}$$

Non-uniform plasma distribution: effective mass is space-time dependent

$$m_{\text{eff}}^2 \equiv \omega_p^2(t, \vec{x}) \sim n_e(t, \vec{x})$$

Can BH superradiance be induced by galactic plasmas?

J. Conlon, C. Herdeiro, PHYS LETT B 780 (2018):

- Superradiance most efficient for

$$M\mu = \left(\frac{M}{M_{\odot}}\right)\left(\frac{\mu}{10^{-10} \text{ eV}}\right) \sim 1$$

- Conditions in ISM seem ideal for *stellar mass BHs*:

$$n_e \sim 10^{-3} - 10 \text{ cm}^{-3} \iff \mu \sim 10^{-12} - 10^{-10} \text{ eV}$$

- PROBLEM: accretion!

The briefest review of accretion theory

Based on the accretion rate ($\dot{M}$ in units of $\dot{M}_{Ed}$), three scenarios:

- $\dot{M} \gtrsim 1$: hot, dense, “thick” disk
- $0.01 \lesssim \dot{M} \lesssim 1$: cold, geometrically “thin” disk
- $\dot{M} \lesssim 0.01$: hot, low density, advection dominated accretion flow (ADAF)
- $\dot{M} \ll 1$: spherically symmetric Bondi accretion

Superradiance in plasma: a toy model

Background: Kerr

Perturbations: massive scalar

Mass profile:

(I)	$\mu_0^2(r) = \mu_H^2 \left(\frac{r_\pm}{r}\right)^\lambda$	Bondi (sph. sym.)
(II)	$\mu_0^2(r) \sin^2 \theta$	ADAF ("thick" disk)
(III)	$\mu_0^2(r) + \mu_\infty^2$	Bondi + μ_∞
(IV)	$\mu_0^2(r) \sin^2 \theta + \mu_\infty^2$	ADAF + μ_∞
(V)	$\mu_1^2(r) = \mu_H^2 \Theta(r - r_0) \left(1 - \frac{r_0}{r}\right) \left(\frac{r_0}{r}\right)^\lambda$	Bondi + inner edge
(VI)	$\mu_1^2(r) \sin^2 \theta$	ADAF + inner edge
(VII)	(V) + μ_∞^2	Bondi + inner edge + μ_∞
(VIII)	(VI) + μ_∞^2	ADAF + inner edge + μ_∞

Numerical setup

Details in *Sam Dolan*, arXiv:1212.1477:

- 1 + 1D decomposition:

$$\Phi(t, r, \theta, \phi) = \sum_{\ell=|m|}^{\infty} \frac{\psi_{\ell m}(t, r)}{r} e^{im\phi} \mathbf{Y}_{\ell m}(\theta)$$

- System of coupled PDEs (ℓ -mode couples to $\ell \pm 2$)
- $O(4)$ finite difference scheme for spatial derivatives: system of coupled ODEs
- Evolve with Runge-Kutta $O(4)$ time step
- Boundary Conditions:
- At infinity: outgoing BC
- Horizon: artificial damping region (perfectly-matched layers, PML) introduced close to left boundary to kill spurious reflections

Boundary conditions: Spurious reflections

- Implementation of BC on finite grids: spurious reflections
- Artificial mirror-like instability from right boundary:

$$\nu_{\ell, mirror} \sim r_{mirror}^{-2(\ell+1)}$$

$$r_{\infty} = 10^3 M \rightarrow \tau_{\infty} \sim 10^{12} M \gg (\nu_{l=m=1})^{-1} \sim 10^6 M$$

- Left boundary: PML

Boundary conditions: Perfectly-matched layers

Introduce artificial damping that suppresses spurious reflections from left boundary (approximated horizon)

$$\ddot{u} = u''$$

First-order rearrangement: $\dot{u} = v'$, $\dot{v} = u'$

Solution: $u \sim \exp -i\omega(t \pm x)$

Introduce x -dependent damping term:

$$\dot{u} = v' - \gamma(x)u$$

$$\dot{v} = u' - \gamma(x)v$$

Solution:

$$u \sim \exp(-i\omega(t \pm x) \pm \int_y^x \gamma(y)dy)$$

Exponential attenuation for both left and right going waves in PML!