

# Painlevé equations and anharmonic oscillators

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**Integrable Systems around the world**

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- 1 The problem relating:
  - Rational solutions of **Painlevé II**
  - Degenerate spectrum of **anharmonic oscillator**
- 2 Link between PII and the anharmonic oscillator
- 3 Our approach
- 4 Exact WKB Method

All in  $\leq 30$  minutes!

$$\text{PII} : v'' = 2v^3 + tv + \alpha, \quad \alpha \in \mathbb{C}$$

## Key property:

- $\exists!$  rational solutions  $\iff \alpha = N \in \mathbb{Z}$ :

$$v(t) = \frac{d}{dt} \log \left( \frac{Y_{N-1}(t)}{Y_N(t)} \right)$$

where  $Y_n(t)$  are **Yablonskii-Vorob'ev polynomials**:  
 $Y_0 = 1, Y_1 = t$  and

$$Y_{N+1} = \frac{tY_N^2 + 4(Y_N')^2 - Y_N Y_N''}{Y_{N-1}}$$

(YES, they're polynomials!)

**Poles** of rational sols of PII  $\longleftrightarrow$  **Roots** of Yablonski-Vorob'ev poly

**Question:** What do the roots of  $Y_n(t)$  look like?

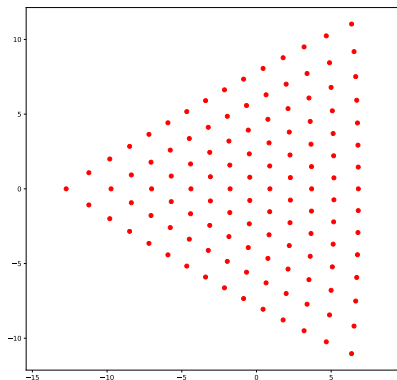


Figure: Roots of  $Y_{15}(t)$

- Buckingham, Miller (2013): Large  $N$  asymptotic analysis via JM Lax pair of PII
- Bertola, Bothner (2015): Hankel determinant expression for  $Y_N^2$ , RHP analysis of (pseudo) orthogonal polynomials

# Anharmonic oscillator

Eigenvalue problem:

$$y'' - \left( \frac{x^4}{4} - \frac{ax^2}{2} - (N+1)x \right) y = \lambda y, \quad y(re^{\pm i\pi/3}) \xrightarrow{r \rightarrow \infty} 0.$$

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If  $N \in \mathbb{N}$  there exists **quasi-polynomial** solutions:

$$y(x) := p(x)e^{-\frac{x^3}{6} + \frac{ax}{2}}, \quad p(x) \text{ polynomial of } \deg(p) \leq N.$$

**Question:** for which  $a \in \mathbb{C}$  are these eigenvalues  $\lambda$  degenerate?

**Answer:** Obtain such  $a \in \mathbb{C}$  as zeros of a discriminant :

$$D_n(a) = \text{poly of } \deg N(N+1)/2.$$

# Roots of the discriminant

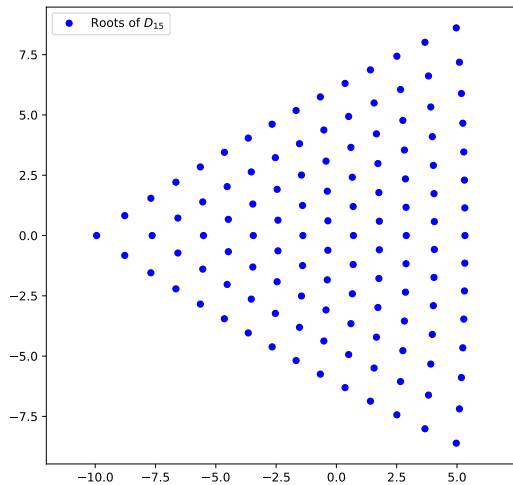


Figure: The answer seems familiar...

# Suspicious coincidence

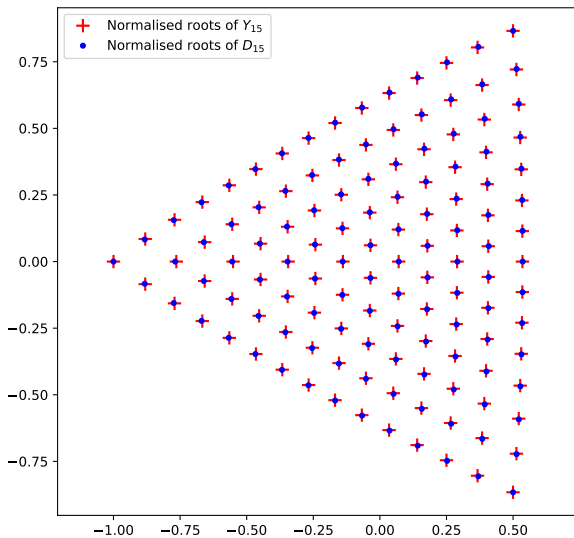


Figure: This image should surprise you



## Shapiro-Tater conjecture (2014)

(After appropriate scaling) the sets

$$\{\text{roots of } D_N(a) = 0\} \text{ and } \{\text{roots of } Y_N(t) = 0\}$$

coincide as  $N \rightarrow \infty$ .

Shapiro-Tater results:

- Support of counting measure for the 'algebraic' eigenvalues.
- Partial results on monodromy of eigenvalues.
- Only numerical evidence towards conjecture.

**Coincidence? Unlikely!**

Lax pair at the pole  $t = a$  of transcendent  $\longrightarrow$  Anharmonic oscillator

$$\begin{cases} \Phi_x = A(x, t)\Phi \\ \Phi_t = B(x, t)\Phi \end{cases} \longrightarrow y'' = Q(x)y$$

**Previous works:**

- Its, Novokshenov (1986)

$$\text{PII}(0) - \text{Flashcka-Newell} \longrightarrow Q(x) = 16x^4 + 8ax^2 + \lambda$$

- Masoero (2010)

$$\text{PI} - \text{tritonquée solution} \longrightarrow Q(x) = 4x^3 - 2ax - 28b$$

where  $a = \text{pole}$ .

Analysis of anharmonic oscillators leads to **asymptotic description of poles**.

# Jimbo-Miwa Lax pair

Jimbo-Miwa (1981) gave a Lax pair for PII:

$$\begin{aligned}\Phi_x &= \left( x^2 \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} + x \begin{bmatrix} 0 & u \\ -2u^{-1}w & 0 \end{bmatrix} + \begin{bmatrix} w + \frac{t}{2} & -uv \\ -2u^{-1}(vw + \theta) & -w - \frac{t}{2} \end{bmatrix} \right) \Phi \\ \Phi_t &= \left( \frac{x}{2} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} + \frac{1}{2} \begin{bmatrix} 0 & u \\ -2u^{-1}w & 0 \end{bmatrix} \right) \Phi, \quad \theta \in \mathbb{C}\end{aligned}$$

where  $u = u(t)$ ,  $v = v(t)$ ,  $w = w(t)$  are functions of  $t$ .

If  $A_t - B_x + [A, B] = 0$ , they satisfy:

$$u_t = -vu, \quad v_t = v^2 + w + \frac{t}{2}, \quad w_t = -2vw - \theta.$$

and  $v(t)$  solves  $\text{PII}(-\theta + 1/2)$ .

**Question:** How to get anharmonic oscillator from here?

## Reduction at pole

Transform the differential equation:  $\Phi_x(x, t) = A(x, t)\Phi(x, t)$  by setting

$$\widehat{\Phi}(x, t) = G(x, t)\Phi(x, t)$$

such that the system becomes:

$$\widehat{\Phi}_x(x, t) = \begin{bmatrix} 0 & 1 \\ Q(x, t) & 0 \end{bmatrix} \widehat{\Phi}(x, t).$$

$\widehat{\Phi} = [y(x, t), y_x(x, t)]^\top \rightarrow$  scalar ODE:  $y'' = Q(x, t)y$ .

Near a pole of PII  $t = a$ , the potential simplifies:

$$\lim_{t \rightarrow a} Q(x; t) = x^4 - ax^2 + 2\theta x - \left(-10b + \frac{7}{36}a^2\right) =: Q_{\text{JM}}(x).$$

Found link between PII and the anharmonic oscillator:

$$Q_{\text{ST}} = \frac{x^4}{4} - \frac{ax^2}{2} - (N+1)x - \Lambda$$
$$\updownarrow$$
$$Q_{\text{JM}} = x^4 - ax^2 + 2\theta x - \lambda$$

Only beginning of the story.

Need to understand the **quasi-polynomials** and **repeated eigenvalue** conditions.

Implement conditions for existence of

- ① quasi-polynomials (QP),
- ② repeated eigenvalues (RE)

asymptotically for large  $N$ , to obtain quantization conditions:

$$\pi i(2k + 1) = N \oint_{\gamma} \sqrt{x^4 - ax^2 + 2x - \lambda} dx$$

which implicitly determine  $(a, \lambda)$  such that we have QP and RE. Then compare with similar quantization conditions for Yablonski-Vorob'ev.

What are these conditions?

# Quasi-polynomial condition

A necessary condition to have quasi-polynomial solutions of

$$y''(x) = \left( \frac{x^4}{4} - \frac{ax^2}{2} - (N+1)x - \Lambda \right) y(x).$$

is the **vanishing of the Stokes matrices**. The solutions  $[P(x), Q_i(x)]$  give explicit Stokes phenomenon:

$S_3 = \begin{bmatrix} 1 & 0 \\ s_3 & 1 \end{bmatrix}$

$S_1 = \begin{bmatrix} 1 & s_1 \\ 0 & 1 \end{bmatrix}$

$S_5 = \begin{bmatrix} 1 & s_{-1} \\ 0 & 1 \end{bmatrix}$

where the Stokes factors are:

$$s_i = \int \frac{1}{P_N(x)^2} e^{x^3/3 - ax} dx$$

## Repeated eigenvalue condition

If  $p(x)e^{\theta(x)}$  is a quasi-polynomial solution, a necessary condition for the existence of repeated eigenvalues is:

$$\int_{\Gamma} p(x)^2 e^{2\theta(x)} dx = 0,$$

where  $\Gamma$  is the contour:

**Question:** how to turn these *exact* conditions into *asymptotic* conditions?



The **exact WKB method** in the spirit of Voros (94) and Kawai, Takei (05).

$$y'' = n^2 \left( \frac{x^4}{4} - \frac{ax^2}{2} - x - \Lambda \right) f = n^2 Q(x)y$$

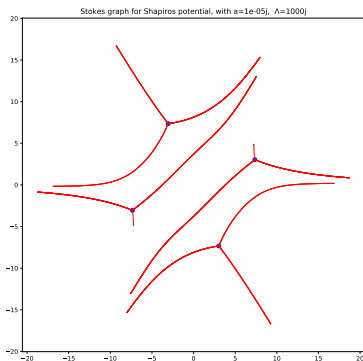
Ansatz  $y(x) = e^{\int^x S(u,n)du}$  where  $S(x, n) = nS_{-1} + S_0 + n^{-1}S_1 + \dots$

$$\begin{aligned} \implies \psi_{\pm}^{(\tau)}(x, n) &= \frac{1}{S_{\text{odd}}(x, n)^{1/2}} \exp \left( \pm \int_{\tau}^x S_{\text{odd}}(u, n) du \right) \\ &= \frac{n^{-1/2}}{V(x)^{1/4}} \exp \left( \pm n \int_{\tau}^x \sqrt{Q(u)} du \right) (1 + \mathcal{O}(n^{-1})) \end{aligned}$$

The WKB solutions  $\psi_{\pm}$  are asymptotic to actual solutions **in certain regions**.

# Generic Stokes graphs

Stokes lines are the level set  $\text{Im} \int_{\tau}^x \sqrt{Q(u)} du = 0$ ,  $\tau = \text{root of } Q(u)$

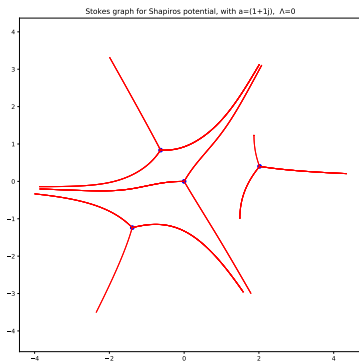


## Theorem

For every region  $\mathcal{D}$  and turning point  $\tau$ , there exists a unique vector solution  $\mathbf{y}$  to ODE s.t.  $\mathbf{y}(x) \sim [\psi_+^{(\tau)}(x, n), \psi_-^{(\tau)}(x, n)]$  in  $\mathcal{D}$  as  $n \rightarrow \infty$ .

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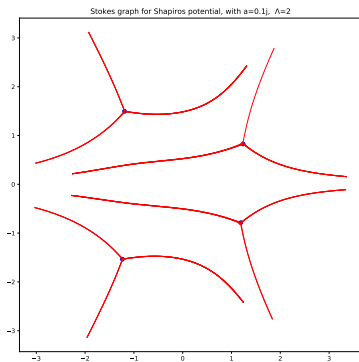


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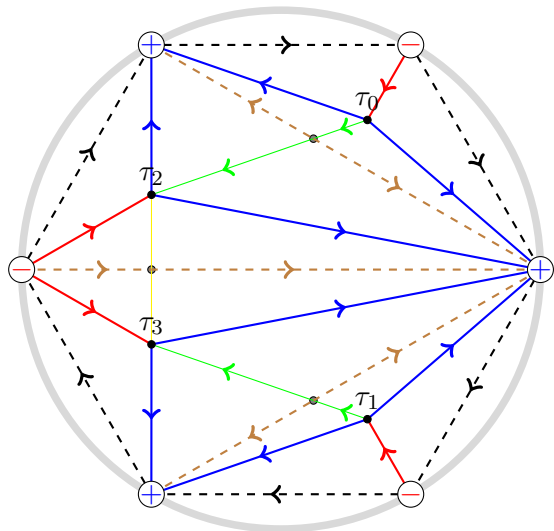


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# RHP for WKB solutions

Jumps across Stokes lines are known (Voros 83) so we can write a RHP.



Computation of **Stokes matrices** using:

$$\begin{bmatrix} 1 & 0 \\ -i & 1 \end{bmatrix}, \quad \begin{bmatrix} 1 & i \\ 0 & 1 \end{bmatrix}$$

$$\begin{bmatrix} 0 & i \\ i & 0 \end{bmatrix}, \quad \begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix},$$

$$e^{\sigma_3 V_{a,b}}, \quad e^{\sigma_3 W_a},$$

$$V_{i,j} = \frac{1}{2} \int_{\tau_i}^{\tau_j} S_{\text{odd}}(x, n) dx,$$

$$W_i = \frac{1}{2} \int_{\tau_i}^{\infty} S_{\text{odd}}^{\text{reg}}(x, n) dx.$$

# Stokes matrices

$$\begin{aligned} S_0 &= \begin{bmatrix} 1 & 0 \\ -ie^{2W_0}(1 + e^{-2V_{1,0}}) & 1 \end{bmatrix}, & S_3 &= \begin{bmatrix} 1 & -ie^{-2W_3}X \\ 0 & 1 \end{bmatrix}, \\ S_1 &= \begin{bmatrix} 1 & -ie^{-2W_2}(1 + e^{2V_{0,2}}) \\ 0 & 1 \end{bmatrix}, & S_4 &= \begin{bmatrix} 1 & 0 \\ -ie^{2W_3} & \end{bmatrix}, \\ S_2 &= \begin{bmatrix} 1 & 0 \\ -ie^{2W_2} & \end{bmatrix}, & S_5 &= \begin{bmatrix} 1 & -ie^{-2W_1}(1 + e^{2V_{3,1}}) \\ 0 & 1 \end{bmatrix}. \end{aligned}$$

where  $X = 1 + e^{-2V_{3,1}} + e^{-2(V_{3,1}+V_{1,0})} + e^{-2(V_{3,1}+V_{1,0}+V_{0,2})}$ .

Stokes matrices vanish if and only if

$$1 + e^{2V_{0,2}} = 0 \iff \pi i (2k + 1) = V_{0,2} = \int_{\gamma_{0,2}} S_{\text{odd}}(x, n) dx,$$

$$1 + e^{2V_{1,3}} = 0 \iff \pi i (2l + 1) = V_{1,3} = \int_{\gamma_{1,3}} S_{\text{odd}}(x, n) dx,$$

**What's next?** Implement (RE) condition asymptotically with exact WKB. *Work in progress*

- ① The conjecture of Shapiro-Tater
- ② Connection between Lax pairs near poles  $\longleftrightarrow$  anharmonic oscillators
- ③ Study of quasi-polynomials and repeated eigenvalues.
- ④ Implement asymptotically via exact WKB
- ⑤ Work in progress - stay tuned!



THANK YOU FOR LISTENING!