

Generalized Gibbs Ensembles

of the Calogero Fluid

GGE



Herbert Spohn

TUMünchen

1. Generalized Gibbs GGE

ideal gas

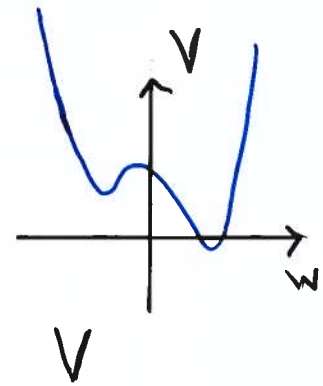
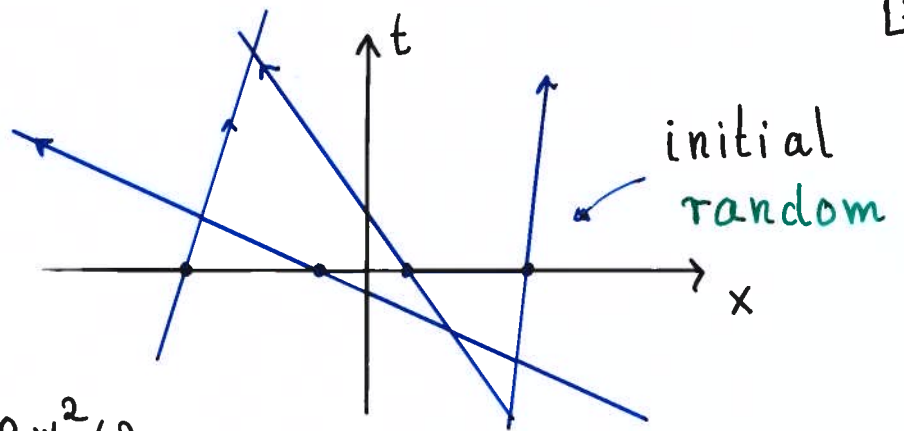
- thermal equilibrium Poisson intensity $p dx \frac{1}{\sqrt{2}} e^{-\beta w^2/2} dw$
- isolated // NO thermalization!

- long time limit $t \rightarrow \infty$

GGE Poisson

$$p dx \frac{1}{\sqrt{2}} e^{-V(w)} dw$$

parameter $\rho > 0$, confining potential



Theorem (Kallenberg 1978)

space time stationary, $\rho > 0$, $\langle \rho_0^2 \rangle < \infty$, finite entropy / length !!

→ mixture of GGE

2. Integrable and interacting

ideal gas $\sum_j p_j$ conserved, no interaction

- 1D fluids

$$H_N = \sum_{j=1}^N \frac{1}{2} p_j^2 + \sum_{i < j=1}^N V_{mec}(q_i - q_j)$$

|| integrable ||

$$I_1, \dots, I_N, I_2 = H_N$$

$$\{I_m, I_n\} = 0$$

meta theorem

ONLY solutions

better: N local conservation laws

$$V_{mec} = 0, V_{mec}(x) = \begin{cases} |x| \leq a \\ 0 & |x| > a \end{cases} \text{ hard rods}$$

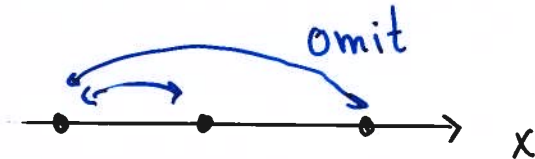
Calogero 1975

$$V(x) = \frac{1}{\sinh^2(x)} + \dots$$

(Sutherland - Calogero - Moser)

$$V(x) = \frac{1}{x^2} \text{ high density limit}$$

- dilute limit



$$H_{\text{total}} = \sum_j \left(\frac{1}{2} p_j^2 + e^{-(q_{j+1} - q_j)} \right)$$

3. Calogero fluid

phase space $\mathbb{R}_{(+)}^N \times \mathbb{R}^N = \Gamma_N$

conserved fields Lax matrix $N \times N$

$$L_{ij} = \delta_{ij} p_j + i (1 - \delta_{ij}) \frac{1}{\sinh(q_i - q_j)} \quad L = L^*$$

Lax pair $L(q, p), M(q, p)$

$$\frac{d}{dt} L = [L, M]$$

$\Rightarrow L \psi_\alpha = \lambda_\alpha \psi_\alpha$ eigenvalues are conserved **NON LOCAL**

$$\left(\sum_{j=1}^N p_j \right)^2$$

\Rightarrow local fields

$$Q^{[n]}(x) = \sum_{j=1}^N \delta(x - q_j) (L^n)_{jj}, \quad Q^{[n]} = \int dx Q^{[n]}(x) = \text{tr } L^n$$

density

total

$Q^{[0]}$ particle number, $Q^{[1]}$ momentum, $Q^{[2]}$ energy, $Q^{[3]}$?, ...

4. GGE on ring



periodize on $[0, l]$

$$f: \mathbb{R} \rightarrow \mathbb{R}$$

$$f_l(x) = \sum_{m \in \mathbb{Z}} f(x + ml)$$

$$\rightsquigarrow H_l \quad L_l$$

still integrable V_{mec} double periodic Weierstrass

$$\rightsquigarrow GGE \quad Q_l^{[n]}$$

$$e^{-\sum_{n=0}^{\infty} \mu_n Q_l^{[n]}} = e^{-\text{tr} V(L_l)}$$

relative to $\frac{1}{N!} d^N q d^N p$
on $[0, l]^N \times \mathbb{R}^N$

parameters volume l , number N
• confining potential V fixed

spacetime stationary

canonical

goal infinite volume $\frac{\ell}{N} = \nu$ fixed

free energy/length $\lim_{\ell \rightarrow \infty} -\frac{1}{\nu} \log Z_N(\nu, V) = F(\nu, V)$

physically more relevant

density of states DOS of Lax

$$\rho_{Q,N}(\omega) = \frac{1}{N} \sum_{j=1}^N \delta(\omega - \lambda_j)$$

↙ eigenvalues

expect

LLN $\lim_{N \rightarrow \infty} \rho_{Q,N} = \rho_Q$ a.s.

• also CLT relates to average currents

NO IDEA

$$V_{\text{mec}}(x) = \frac{1}{x^2}$$

P. Choquard 2000

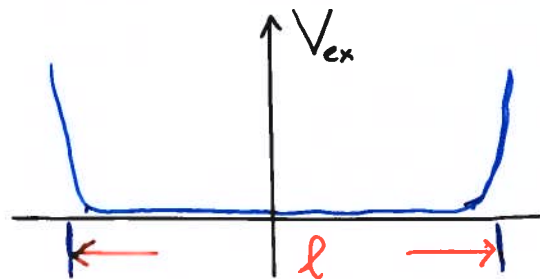
$$\left(\frac{1}{x^2}\right)_\ell = \frac{1}{\sin^2(x/\ell)}$$

5. External potential l

$$T_N = W_N \times \mathbb{R}^N$$

↑

$$\{x_1 < \dots < x_N\}$$



$$Z_N(\nu, V) = \int_{T_N} d^N q d^N p e^{-\text{tr} V(L)} e^{-\sum_{j=1}^N V_{ex}(q_j)}$$

Ruijsenaars 1988 - 1995

scattering coordinates

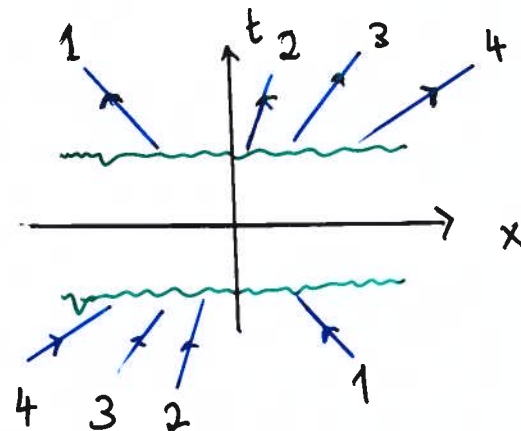
$\lambda \in W_N$ eigenvalues of L

$$\lim_{t \rightarrow \infty} p_j(t) = \lambda_j, \quad \lim_{t \rightarrow \infty} q_j(t) - \lambda_j t = \phi_j \in \mathbb{R}$$

$\Phi: (\lambda, \phi) \mapsto (q, p)$ one-to-one, canonical
generic

Result (R 1988) algebraic construction of Φ

$$Z_N(\nu, V) = \int_{T_N} d^N \lambda d^N \phi e^{-\sum_{i=1}^N V(\lambda_i)} e^{-\sum_{i=1}^N V_{ex}(q_i(\lambda, \phi))}$$



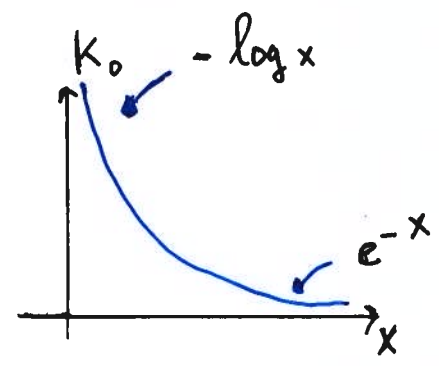
• special choice $V_{ex}(x) = e^{-l/2} \cosh x$

$$\sum_{j=1}^N V_{ex}(q_j) = \sum_{j=1}^N e^{-l/2} \gamma_j \cosh \phi_j$$

$$\gamma_j = \prod_{\substack{m=1 \\ m \neq j}}^N \left(1 + \frac{1}{(\lambda_m - \lambda_j)^2} \right)^{1/2}$$

confining!

modified Bessel



⇒

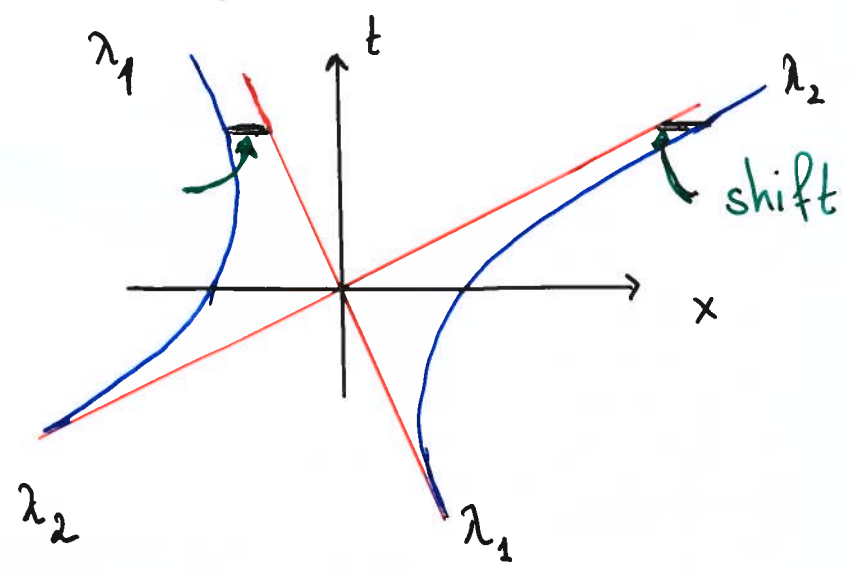
$$Z_N(\lambda, V) = \frac{1}{N!} \int_{\mathbb{R}^N} d\lambda \prod_{j=1}^N e^{-V(\lambda_j)} \prod_{j=1}^N 2 K_0(2 e^{-l/2} \gamma_j)$$

|| mean-field ||

• Calogero 2-particle scattering shift

$$\phi_{ca}(w) = -\log\left(1 + \frac{i}{w^2}\right)$$

$$W = \lambda_1 - \lambda_2$$



6. free energy functional (1-particle)

$$\rho \geq 0, \int_{\mathbb{R}} dw \rho(w) = \frac{1}{\nu}$$

$$\mathcal{F}(\rho) = \int_{\mathbb{R}} dw \rho(w) \left(V(w) - 1 + \log \rho(w) - \log \left(1 + \int_{\mathbb{R}} dw' \rho(w') \phi_{ca}(w-w') \right) \right)$$

minimizer ρ^* (unique)

$$\mathcal{F}(\rho^*) = \mathcal{F}(\nu, V) = \lim_{N \rightarrow \infty} -\frac{1}{N} \log Z_N(\nu, V) \quad \parallel \quad \nu = \frac{l}{N}$$

Lax DOS $\rho_a = \nu \rho^*$

TBA formalism (Yang, Yang 1969)

Lagrange multiplier μ minimizer $\rho_p(\mu)$

$$T\psi(w) = \int dw' \phi_{ca}(w-w') \psi(w')$$

TBA equation $\epsilon = V - 1 - \mu - T e^{-\epsilon}$

quasi-energy $\rho_n = e^{-\epsilon}$

$$\rho_p(\mu) = \frac{1}{1 - \rho_n T} \rho_n$$

adjust $\mu \rightsquigarrow \rho^*$

7. Bethe equations

- Lieb-Liniger δ -Bose gas

phase shift θ_{ll} , scattering shift $\phi_{ll} = \theta'_{ll}$

$$\underbrace{2\pi I_j}_{\text{input}} = v N k_j + \sum_{i=1}^N \theta_{ll}(k_j - k_i)$$

output (k_1, \dots, k_N)

$$\frac{2c}{w^2 + c^2}$$

$$I_1 < \dots < I_N$$

- DOS

$$\rho_{Q,N}(w) = \frac{1}{N} \sum_{j=1}^N \delta(w - k_j)$$

I_j : integer

counting measure

weight $e^{-\sum_{j=1}^N V(k_j)}$

$$\lim_{N \rightarrow \infty} \rho_{Q,N} = \rho_Q$$

Dorlas, Lewis, Pule' 1989

- Calogero fluid

$$y_j = v N \lambda_j + \sum_{i=1}^N \theta_{ca}(\lambda_j - \lambda_i)$$

$$V(w) = \frac{1}{2} \beta w^2 - \mu \beta$$

thermal

$$y \in W_N$$

weight $e^{-\sum_{j=1}^N V(\lambda_j)}$

DOS

Lebesgue

|| asymptotic BA ||

~ agrees with previous result ~
for $N \rightarrow \infty$

Outlook

HS 2020, Guionnet, Memin 2021

10

- Toda lattice $\sum_{j=1}^N \frac{1}{2} p_j^2 + \sum_{j=1}^{N-1} e^{-(q_{j+1} - q_j)}$ $\phi_{to}(w) = \log w^2$ $w \rightarrow \infty$ of $-\log(1 + \frac{1}{w^2})$

|| same method ||

- Calogero-Moser $V_{mec}(x) = \frac{1}{x^2}$ | classical + quantum on ring μ of w
TBA exact

|| $\phi_{cm} = 0$ || \Rightarrow

$$\partial_t \rho(x, t; w) + \partial_x (w \rho(x, t; w)) = 0$$

non-interacting

↑
local DOS

Calogero fluid

$$\partial_t \rho(x, t; w) + \partial_x \left(\underbrace{v^{eff}(x, t; w)}_{\text{local nonlinear functional of } \rho} \rho(x, t; w) \right) = 0$$

|| generalized hydrodynamics ||

interacting

|| local nonlinear functional of ρ ||