

# The threshold of gravitational collapse in vacuum

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## Gravity at the Edge

We wish to find and understand solutions to the (vacuum) field equations  $R_{ab} = 0$  at their most extreme.

Highly dynamical, non-linear PDEs. Strategy?

- ▶ Mathematical relativity; the use of rigorous mathematics to understand the solution space, where possible demonstrating specific, sharp, estimates.
- ▶ Numerical relativity; use of numerical methods to obtain *in hand* approximate solutions that converge to the continuum solution as you use more computational resources.

The two compliment each other's strengths and can be combined to unearth the truth.

# The cosmic censorship conjectures

The most important open conjectures in GR for asymptotically flat spacetimes are those of cosmic censorship:

- ▶ Weak: generically singularities lie inside black holes. (Global existence outside BHs).
- ▶ Strong: maximal Cauchy developments are inextendible. (Global uniqueness).

(Proper formulation needs care).

How can we attack this?

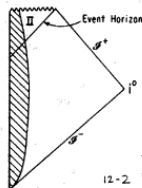


Fig. 12.2. Another representation of the closure,  $\bar{M}$ , of the physical spacetime depicted in Figure 12.1. As in Figure 12.1, the angular dimensions are suppressed so each point in this diagram (except those at  $r = 0$  and the point  $i^0$ ) represents a 2-sphere.

Wald 1984

# Weak cosmic censorship: spherical symmetry

- ▶ Christodoulou 1987-1997.
- ▶ Choptuik, 1992, Collapse of spherical scalar field:
  - Small blackholes can be created,  $M \propto (p - p_*)^\gamma$ .
  - $\gamma$  universal.
  - $p = p_*$  self-similar, unique.

*Critical phenomena in gravitational collapse!*

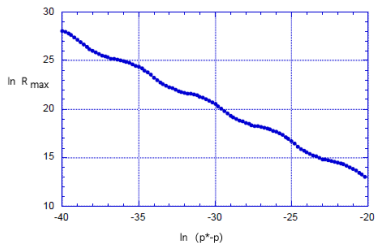


FIG. 1.  $\ln R_{\max}$  is plotted vs.  $\ln(p_* - p)$ . The result is a line with slope  $-2\gamma$  and a periodic wiggle with period  $\Delta/(2\gamma)$

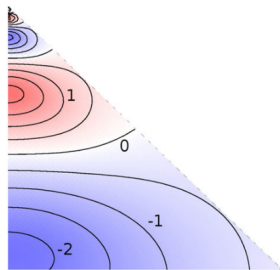
Garfinkle & Duncan 1998: clever coordinates. No AMR.



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Reiterer & Trubowitz 2012: Choptuik spacetime exists (and is real analytic)!

# Weak cosmic censorship conjecture: beyond spherical

Vacuum. Consider initial data with no trapped surface.

- ▶ Small waves disperse.  
Christodoulou & Klainerman, 1993.
- ▶ Proof: trapped region can form.  
Christodoulou, 2008.
- ▶ Numerics: Abrahams & Evans, 1993.

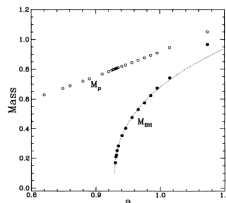


FIG. 2. Critical behavior of black-hole mass. Black-hole masses  $M_{\text{BH}}$  derived from a sequence of simulations are plotted (filled circles) as a function of initial wave packet amplitude. The quasilocal (Brill) masses  $M_p$  of the wave packets are shown (open circles) for comparison. The dotted curve represents the best-fit power law  $M_{\text{BH}} = C(a - a^*)^\beta$  with values  $a^* = 0.928$ ,  $C = 1.750$ , and  $\beta = 0.369$ . Wave packet width  $L = 2\pi$  is chosen to normalize the mass scale.

Despite multiple attempts, nobody has reproduced the latter.

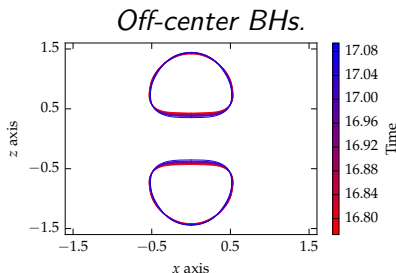
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DH, A. Weyhausen & B.Brüggemann [Phys.Rev.D93, (2016) 063006, Phys.Rev.D96, (2017) 104051].

Method *vastly* superior than earlier approaches. But problems persist.

Brill Waves. Key difficulties:

- ▶ No evidence of scaling in first peak. Hard to pass.
- ▶ Null-infinity?
- ▶ Resolving fine features.
- ▶ Coordinate singularities.



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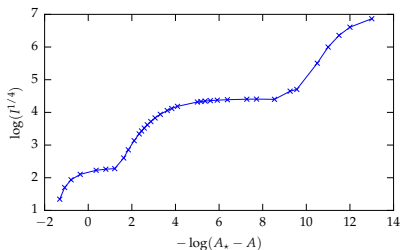
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*Scaling of Kretschmann.*



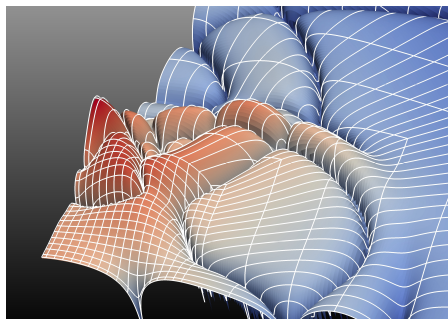
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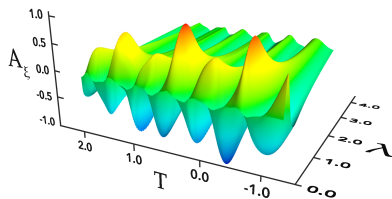


# Aspherical toy models I: SphGR - Maxwell

Baumgarte, Gundlach & DH. [Phys.Rev.Lett.123, (2019) 171103].

To what extent are our difficulties caused by the lack of symmetry? In scalarfield collapse with large asphericity the Choptuik solution suffers from a second decay channel. What else can we do?

- ▶ Maxwell minimally coupled to GR: dynamical solutions aspherical.
- ▶ Curvature quantities still scale near the threshold.
- ▶ But the threshold solution seems no longer unique!



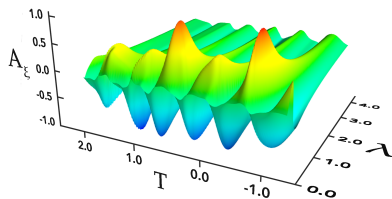
Spacetime plot of field variable at threshold.

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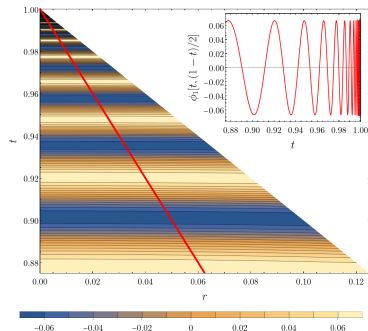
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# Aspherical toy models II: Semi-linear wave equations

Suárez Fernández, Vicente & DH. [Phys.Rev.D103, (2021) 044016].

What is the most stupid model that we can concoct? Surprisingly just applying a *deformation* function  $\phi = \mathcal{D}[\varphi]$  to solutions of the wave equation solves many riddles!

- ▶ Spherical: complete critical phenomenology recovered.
- ▶ General: power law scaling persists.
- ▶ General: uniqueness of threshold solutions lost!



Spherical threshold solution.

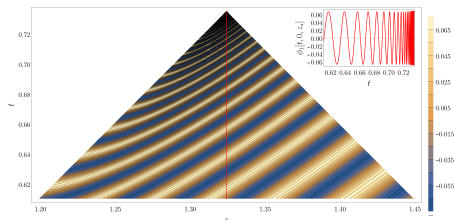


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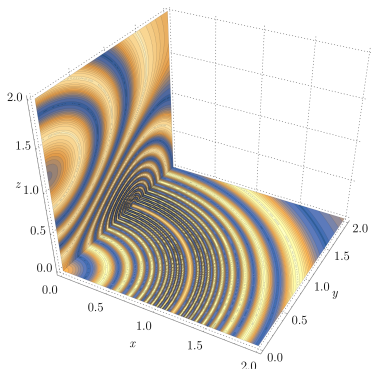
Axisymmetric threshold solution.

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Another axisymmetric (near) threshold solution!

# Prague - universal curvature features

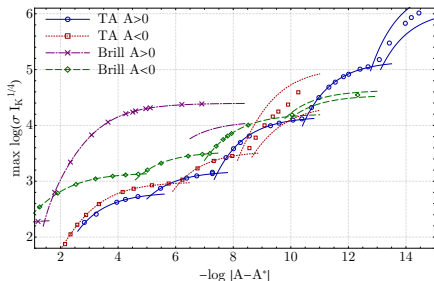
Khirnov, Ledvinka. [Class.Quant.Grav.35, (2018) 215003], Ledvinka, Khirnov, [Phys.Rev.Lett.127, (2021) 011104]

The Prague code:

- ▶ Cartoon, FD within ET.
- ▶ 'Quasimaximal' slicing.
- ▶ Optimized in assembly!

Tuning to threshold:

- ▶ Power-laws; not universal.
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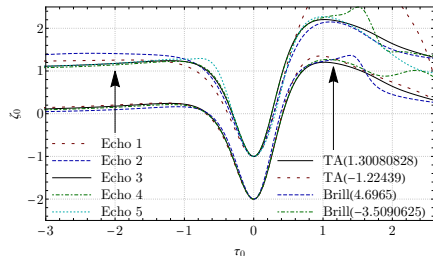
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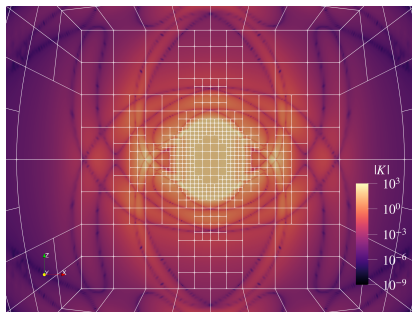
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hp-refinement bamps upgrade  
reduces cost by  $O(10)$  factor.  
Complete rewrite of AH finder.

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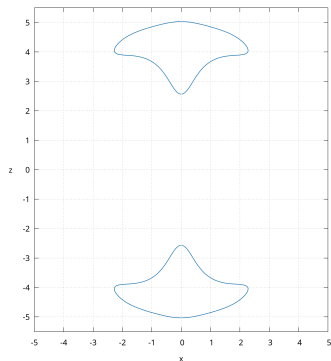
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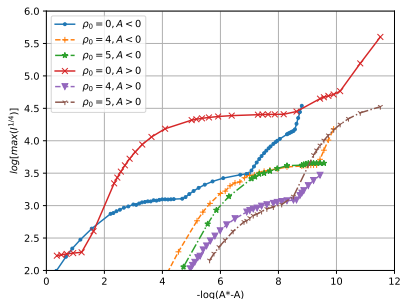
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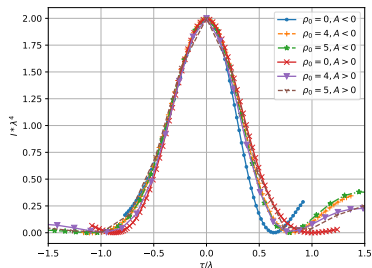
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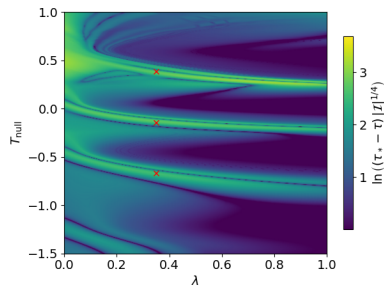
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Shock-avoiding slicing condition allows better tuning with SphGR.

Families with accumulation at center?

- ▶ Yes, with *approximate* DSS.
- ▶ Observe power-law with wiggle.



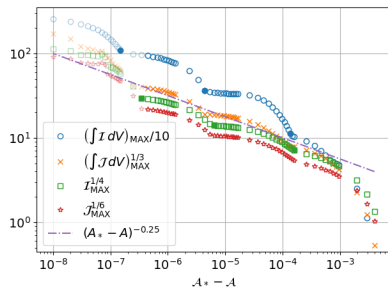
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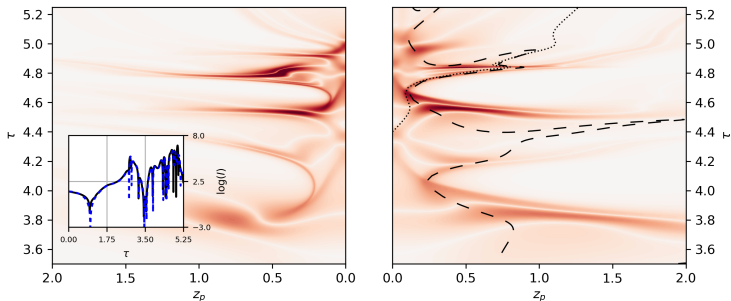


# The end of universality

Baumgarte, Brüggemann, Cors, Gundlach, DH, Renkhoff, Khirnov, Ledvinka, Suárez Fernández. [arXiv:2305.17171].

Code comparison {bamps, prague, sphGR}:

- ▶ Quantitative agreement in canonical coordinates.
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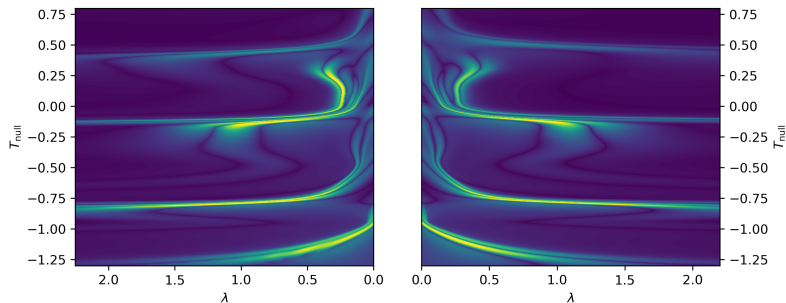


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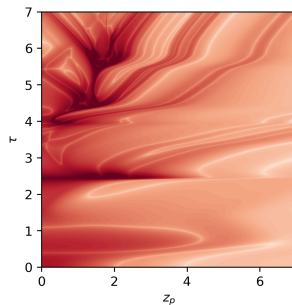
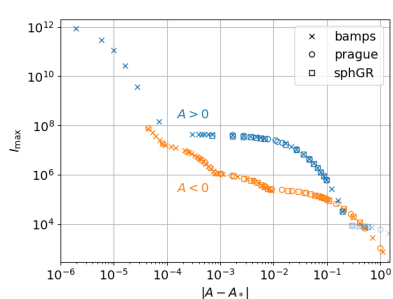


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# Conclusions

We seek a complete picture of the threshold of blow-up in vacuum!

- ▶ Mathematical and numerical relativity offer complimentary approaches to our understanding of GR. Extreme spacetimes require *ever more sophisticated* methods. We attack both.
- ▶ In spherical symmetry the threshold of gravitational collapse is well understood. In the more general setting *a new picture* is emerging. Limited aspects of the phenomenology survive.
- ▶ The deliverables of our research program have broad utility for extreme spacetimes, compact binaries and gravitational wave astronomy.