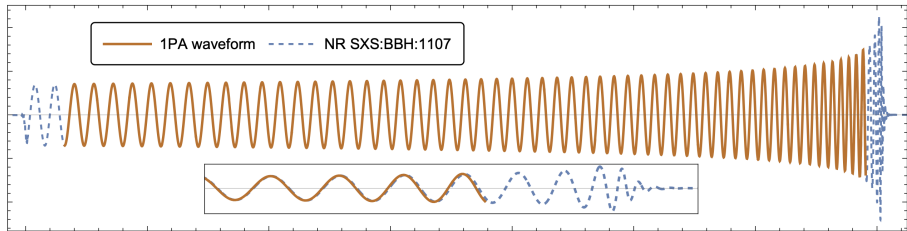


Status of post-adiabatic EMRI modelling



Wardell, Pound, Warburton, Durkan, Miller, Le Tiec [arXiv:2112.12265]

Leor Barack (Southampton)

Self-force and adiabatic expansion 3-layer structure

$$\Gamma_{\mu\nu}(g_{\alpha\beta})=0 \quad / \quad g_{\alpha\beta} = g_{\alpha\beta}^{\text{KERR}} + \varepsilon h_{\alpha\beta}^{(1)} + \varepsilon^2 h_{\alpha\beta}^{(2)} + \dots$$

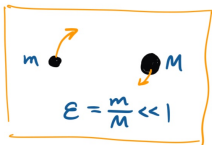
SELF-FORCE THEORY
from PDEs to
point-particle orbits

$$\ddot{X}^\alpha = \varepsilon F_{(1)}^\alpha + \varepsilon^2 F_{(2)}^\alpha + \dots$$

TWO-TIMESCALE/ADIABATIC
EXPANSION

$$\varphi = \varepsilon^{-1} \varphi_{\text{OPA}}(\varepsilon t) + \varepsilon^0 \varphi_{\text{IPA}}(\varepsilon t) + O(\varepsilon)$$

Should suffice for
parameter extraction
if ε sufficiently small



Self-force and adiabatic expansion

$$\mathcal{G}_{\mu\nu}(g_{\alpha\beta}) = 0 \quad / \quad g_{\alpha\beta} = g_{\alpha\beta}^{\text{KERR}} + \varepsilon h_{\alpha\beta}^{(1)} + \varepsilon^2 h_{\alpha\beta}^{(2)} + \dots$$

Equation
of Motion

$$\ddot{X}^\alpha = \varepsilon F_{(1)}^\alpha + \varepsilon^2 F_{(2)}^\alpha + \dots$$

$\langle F_{(1)} \rangle$

$\langle F_{(2)} \rangle$

phase evolution
(Radiation-reaction
timescale)

$$\varphi = \varepsilon^{-1} \varphi_{\text{OPA}}(\varepsilon t) + \varepsilon^0 \varphi_{\text{IPA}}(\varepsilon t) + O(\varepsilon)$$

Rapid waveforms

- Treat $h_{\alpha\beta}^{(n)}$ as functions on extended manifold:

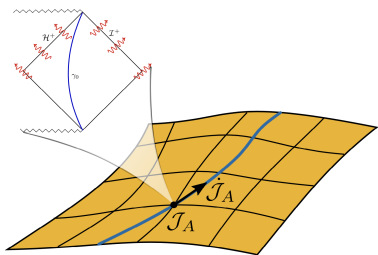
$$h_{\alpha\beta}(t, x^i) \rightarrow \sum_{n=1}^{\infty} \epsilon^n h_{\alpha\beta}^{(n)}(x^i; \mathcal{J}_A, \varphi_A)$$

where $\mathcal{J}_A = (J_A, M_{\text{BH}}, J_{\text{BH}}, \dots)$.

- With a suitable choice of (J_A, φ_A) :

$$h_{\alpha\beta}^{(n)} = \sum_{k^A} h_{\alpha\beta}^{(n)\Omega_k}(x^i; \mathcal{J}_A) e^{-ik^A \varphi_A}$$

with $\Omega_k := k^A \dot{\varphi}_A$.



Offline step: Solve field equations for amplitudes $h_{\alpha\beta}^{(n)\Omega_k}$ on grid of \mathcal{J}_A values.

Online step (FEW): Rapidly evolve through parameter space

$$\dot{\varphi}_A = \Omega_A(\mathcal{J}_B)$$

$$\dot{\mathcal{J}}_A = \epsilon \tilde{F}^{(0)}(\mathcal{J}_B) + \epsilon^2 \tilde{F}^{(1)}(\mathcal{J}_B) + \dots$$

Status summary

• OPA:

- Full-SF generic inspirals in Kerr are available (numerical $h_{\alpha\beta}^{(0)\Omega_k}$).
- FEW implemented in Schwarzschild with full SF
- FEW implemented in Kerr with 5.5PN fluxes and AAK waveform amplitudes
- (“Soon”) FEW for equatorial inspirals in Kerr with full SF

• 1PA:

- Quasicircular inspiral in Schwarzschild, generic/precessing CO spin
- Quasicircular inspiral with slowly-spinning primary, aligned CO spin

• Synergies

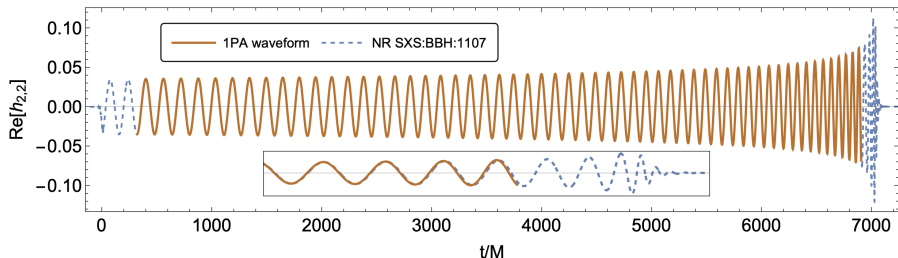
- With PN: High-order PA augmentation of 1PA SF model
- With EOB: benchmarking & calibration for EOB
- With NR: benchmarking for SF; NR+SF worldtube excision model

• Misc.

- Advanced methods for solving field equations
- Work on transition to plunge, merger & ringdown—relevant for IMRIs.

1PA waveforms

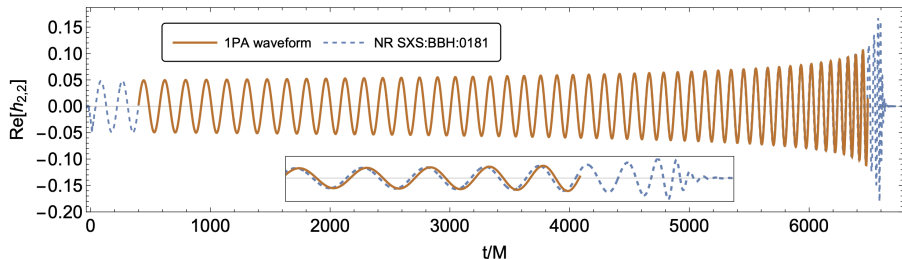
Mass ratio $\epsilon = 1/10$



(Wardell, Pound, Warburton, Durkan, Miller, Le Tiec)

1PA waveforms

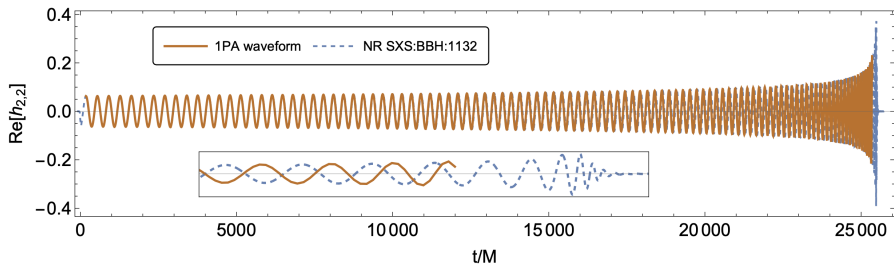
Mass ratio $\epsilon = 1/6$



(Wardell, Pound, Warburton, Durkan, Miller, Le Tiec)

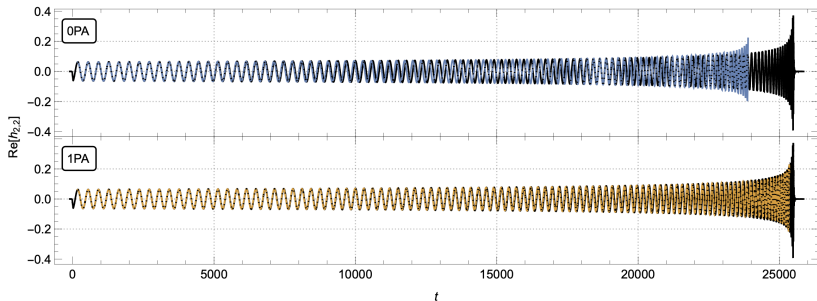
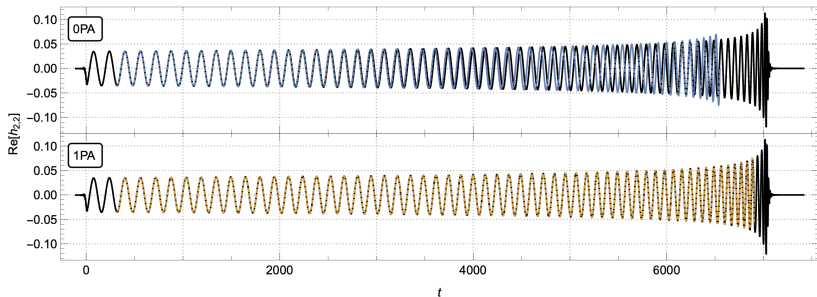
1PA waveforms

Mass ratio $\epsilon = 1$

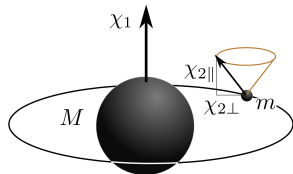
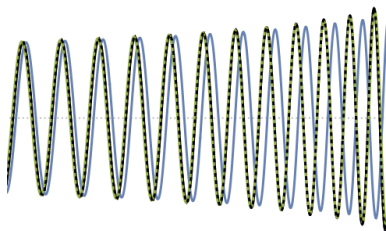
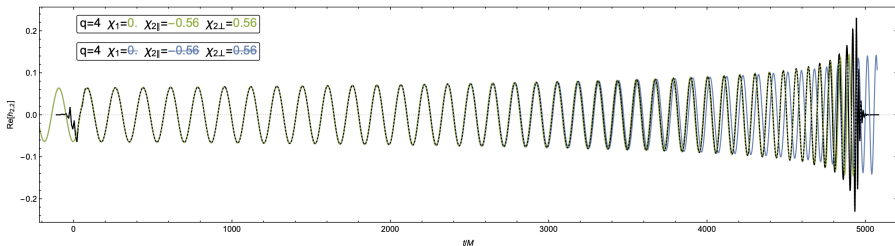


Error estimate: $\sim 7.5\epsilon$ rad from $R = 20M$ to ISCO

0PA vs. 1PA [$\epsilon = 1 : 10$ (top), $1 : 1$ (bottom)]

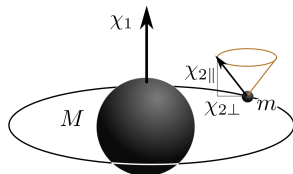
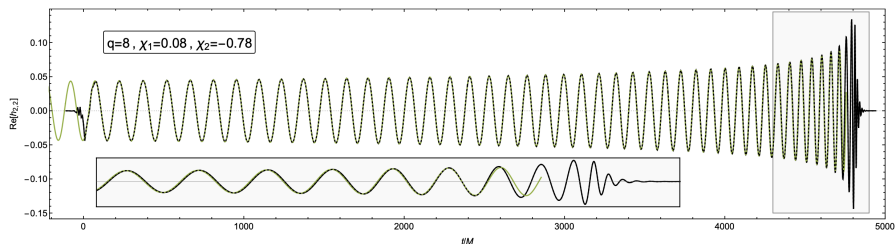


1PA inspiral with CO spin



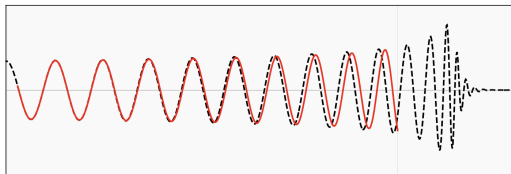
(Mathews, Pound, Wardell)

1PA inspiral with small MBH spin and (anti-)aligned CO spin

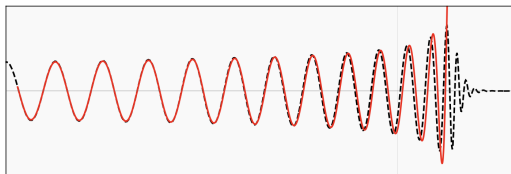


(Mathews, Pound, Wardell)

Transition to plunge



$$\frac{d\Omega}{dt} = F_{\Omega}^{(0)}$$
$$h_{lm} = \epsilon h_{lm}^{(1)}(\Omega) e^{-im\phi_p}$$

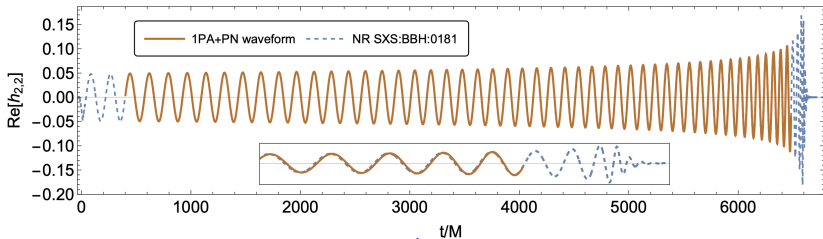
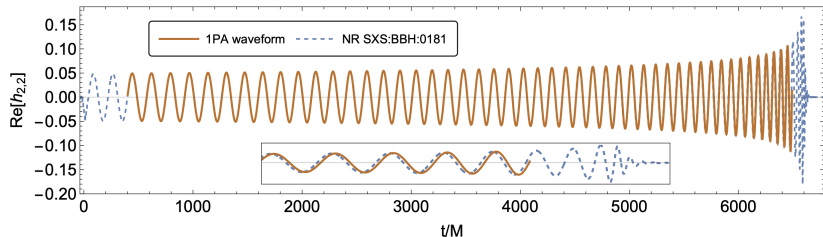


$$\frac{d\Delta\tilde{\Omega}}{dt} = \epsilon^{1/5} \left(F_{\Delta\tilde{\Omega}}^{(0)} + \epsilon^{2/5} F_{\Delta\tilde{\Omega}}^{(2)} \right)$$
$$h_{lm} = \epsilon \left(j_{lm}^{(0)} + \epsilon^{2/5} j_{lm}^{(2)} + \epsilon^{3/5} j_{lm}^{(3)} \right) e^{-im\phi_p}$$

(Compere, Durkan, Kuchler, Pound)

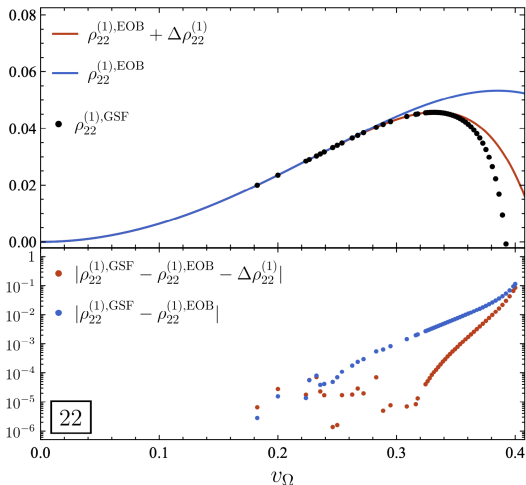
5PA PN augmentation of 1PA SF model

Mass ratio $\epsilon = 1/6$



(Pound, Warburton, Wardell, Durkan, Miller)

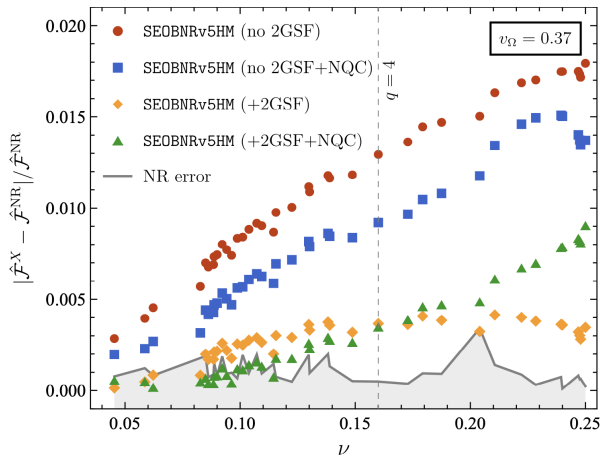
EOB Calibration modal flux, quasicircular inspiral



(Van de Meent et al 2303.18026)

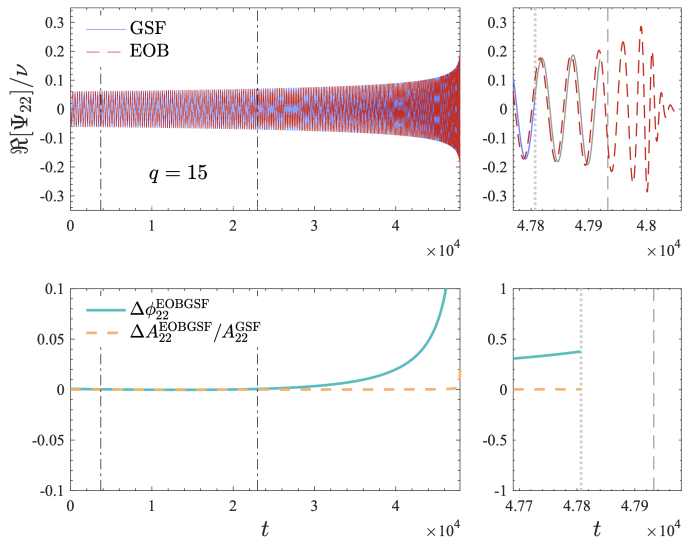
$$\Delta\rho_{22}^{(1)} = 21.2v_\Omega^8 - 411v_\Omega^{10}$$

EOB Calibration modal flux, quasicircular inspiral



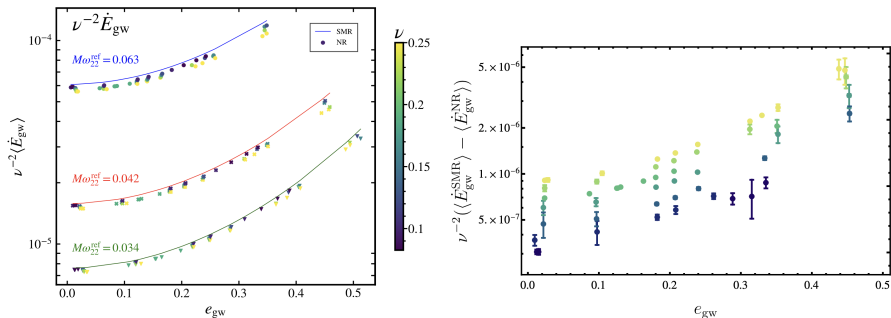
(Van de Meent et al 2303.18026)

TEOBResumS benchmarking



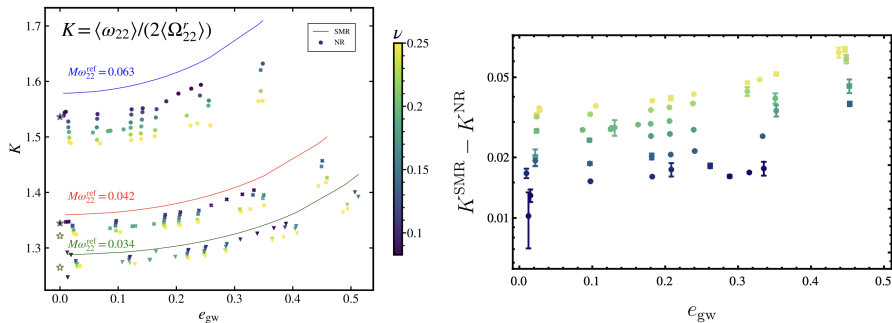
(Albertini et al 2208.01049, 2208.02055)

SF benchmarking using NR energy flux, eccentric inspiral



(A. Ramos-Buades et al. 2209.03390)

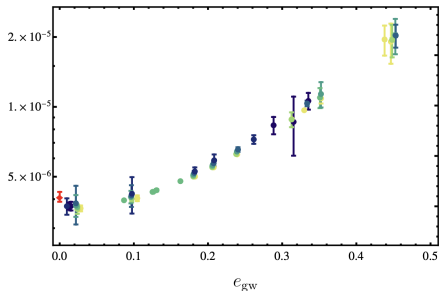
SF benchmarking using NR perisatron advance, eccentric inspiral



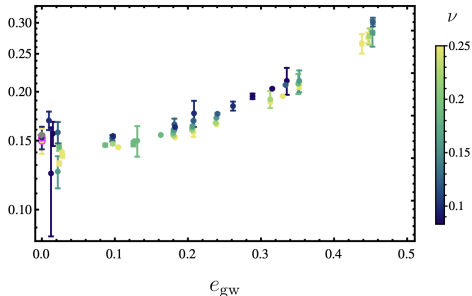
(A. Ramos-Buades et al. 2209.03390)

≥ 2 PA contributions are strangely small

$$\frac{\dot{E}^{\text{SF}} - \dot{E}^{\text{NR}}}{\nu^3}$$



$$\frac{K^{\text{SF}} - K^{\text{NR}}}{\nu}$$

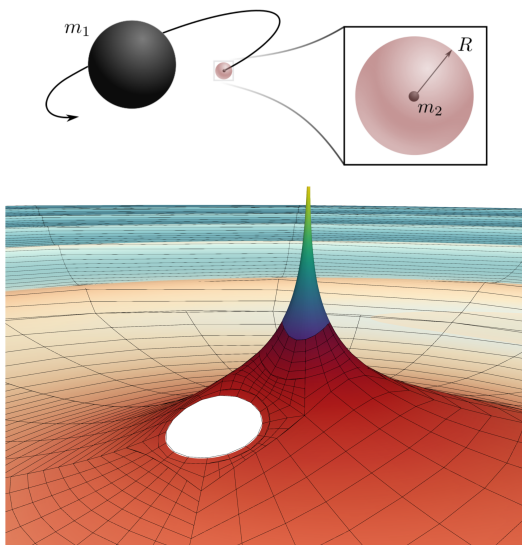


$$\frac{\dot{E}^{\text{SF}} - \dot{E}^{\text{NR}}}{\nu^3} = f_{1\text{PA}}(e) + f_{2\text{PA}}(e)\nu + f_{3\text{PA}}(e)\nu^2 + \dots$$

Weak ν dependence $\Rightarrow f_{\geq 2\text{PA}}$ very small

Putting SF and NR together: worldtube excision for IMRIs

(Dhesi, Wittek, LB, Pfeiffer, Pound, Rüter)



PA modelling and Environmental Effects

- Variety of EEs potentially discernible at 1PA
- Some EEs can be incorporated as sourcing terms within existing PA scheme
- But EEs can also come with new time/lengthscales, requiring new types of multiple-scale expansions
- Degeneracy with/between EEs resolved with multiple EMRI observations
- Whether EEs are signal to be extracted or noise to be removed, 1PA-accuracy models of clean EMRIs will be crucial
- (Also good value for money, given ≥ 2 PA terms appear to be negligible)