

Amplitudes in AdS from Conformal Field Theory

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Conformal invariance

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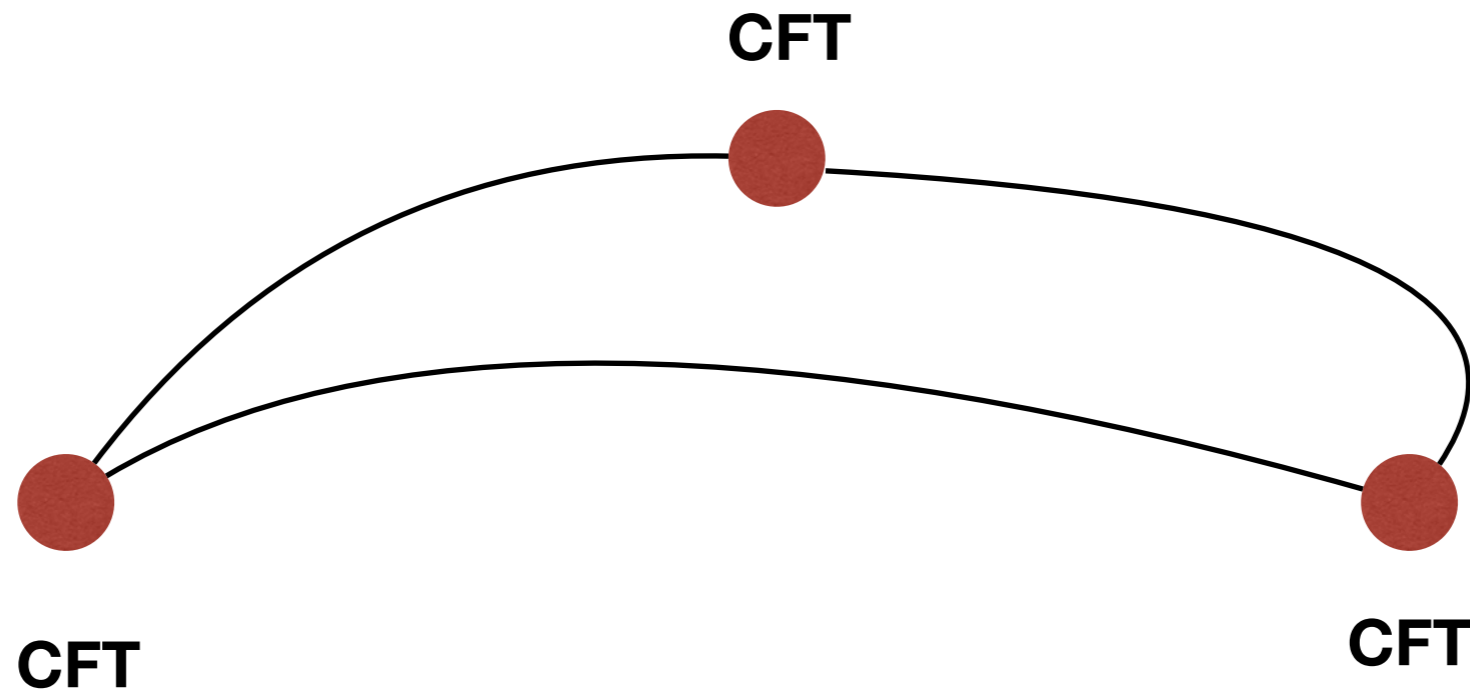
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The description of fixed points boils down to classifying conformal field theories.

Centrality of CFTs

Conformal Field Theories (CFTs) are central also in the characterisation of QFTs.

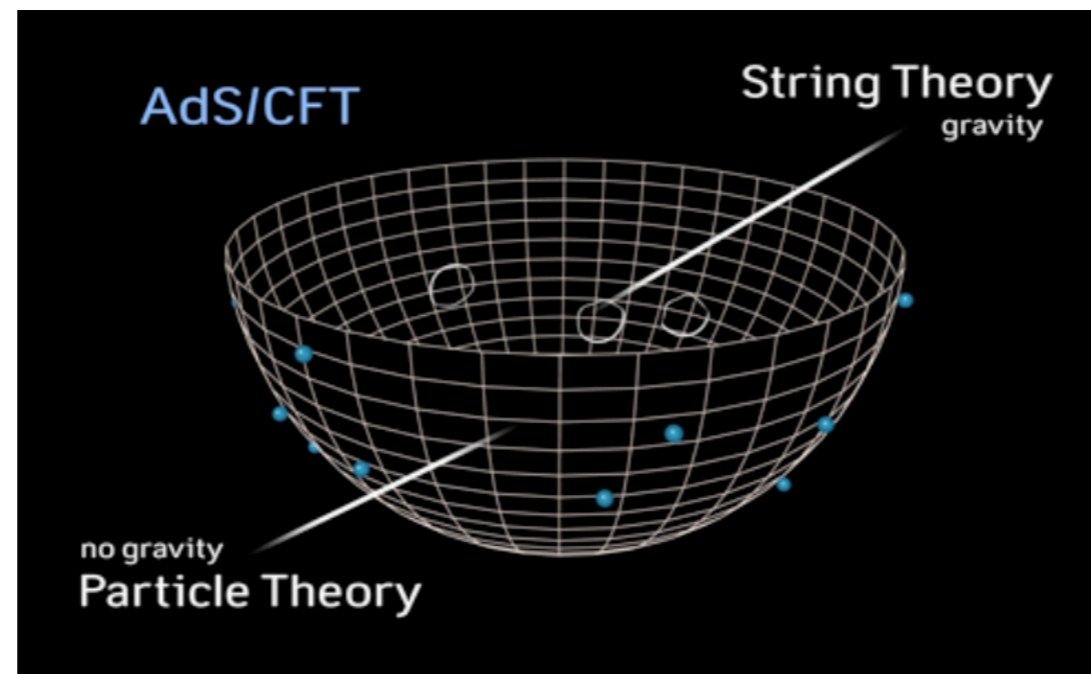


Large classes of QFTs can be seen as RG flows which emerge from a CFT (UV fixed point) and another non trivial CFT (IR fixed point)

Centrality of CFTs

They are related to theories of quantum gravity via the AdS/CFT correspondence

Operative mapping: observables (correlation functions and scattering amplitudes) in both theories are related in a very specific way.



Maldacena 1998

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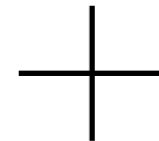
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existence of the operator product expansion (OPE)

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Here it is shown for scalars, but it is similar for spinning operators.

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What about four point correlators?

$$\langle \mathcal{O}(x_1)\mathcal{O}(x_2)\mathcal{O}(x_3)\mathcal{O}(x_4) \rangle = \frac{\mathcal{G}(u, v)}{x_{12}^{2\Delta_0} x_{34}^{2\Delta_0}}$$

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cross ratios

$$u = \frac{x_{12}^2 x_{34}^2}{x_{13}^2 x_{24}^2} \quad v = \frac{x_{14}^2 x_{23}^2}{x_{13}^2 x_{24}^2}$$

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Using the OPE inside correlators of n -points, with $n \geq 4$, it is possible to reduce them to two point functions.

Conformal blocks

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$$\downarrow$$
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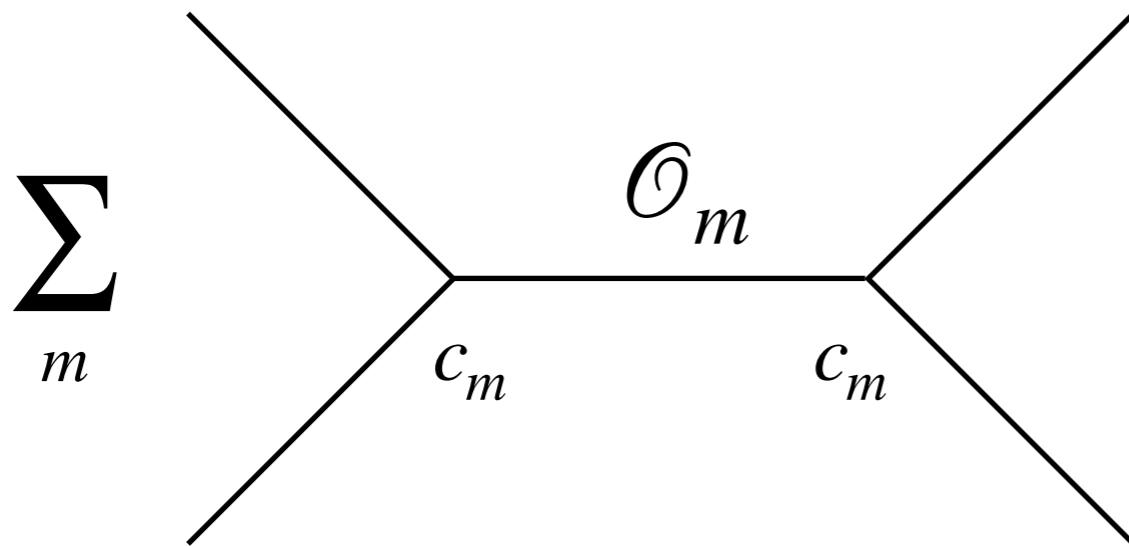
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What is m ? It denotes the quantum numbers of the exchanged operators, which in this particular case are the conformal dimension and the Lorenz spin (Δ, ℓ) .

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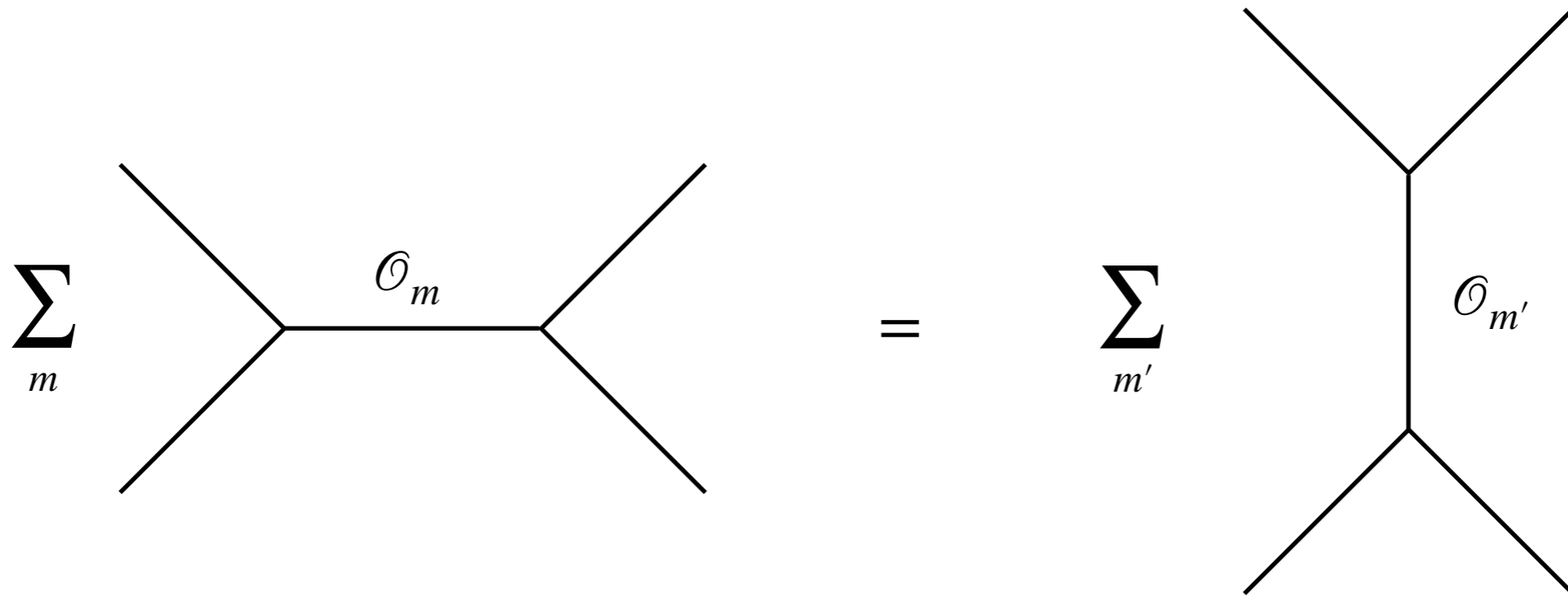
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We won't be able to solve them completely, but we will discuss some approaches to find solutions consistent with the crossing relations.

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It is related to the fact that the norm of a state is positive. In this context it means

that $\Delta \geq \frac{d-2}{2}$ for scalars and

$\Delta \geq d + \ell - 2$ for operators of spin ℓ ,
and that the $c_m \in \mathbb{R}$

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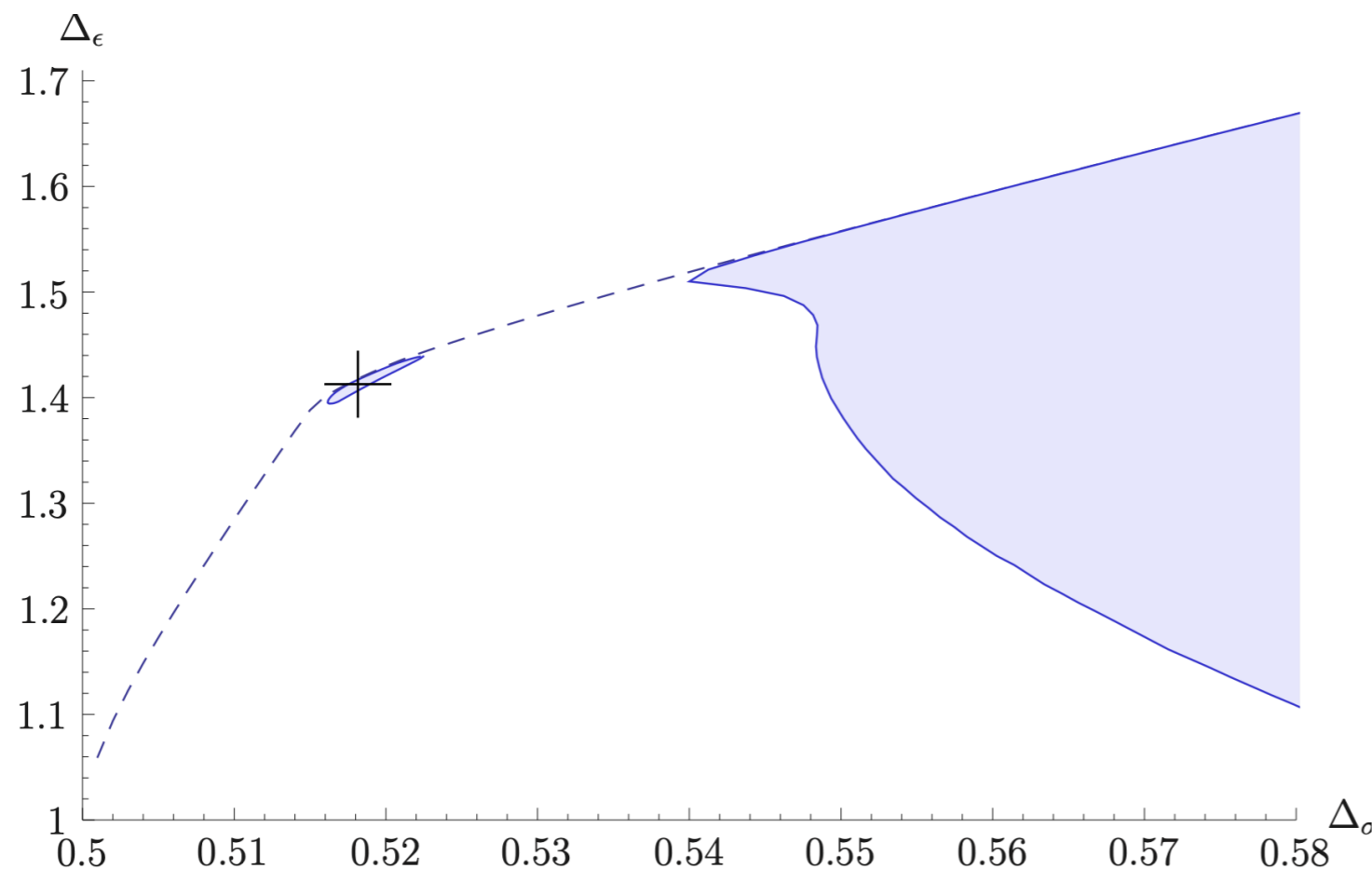
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3d Ising Model

Rattazzi Rychkov Tonni Vichi 2008
Kos Poland Simmons Duffin 2014

Numerical Bootstrap

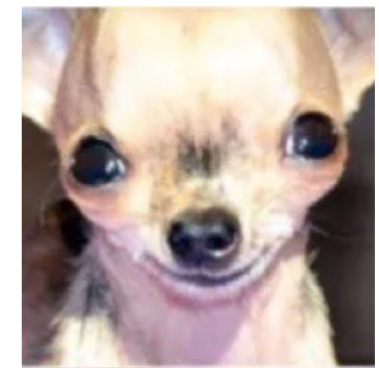
The idea is to use the crossing relations as necessary conditions for conformal dimensions and OPE coefficients to belong to a CFT.

Tentative CFT data



crossing relations

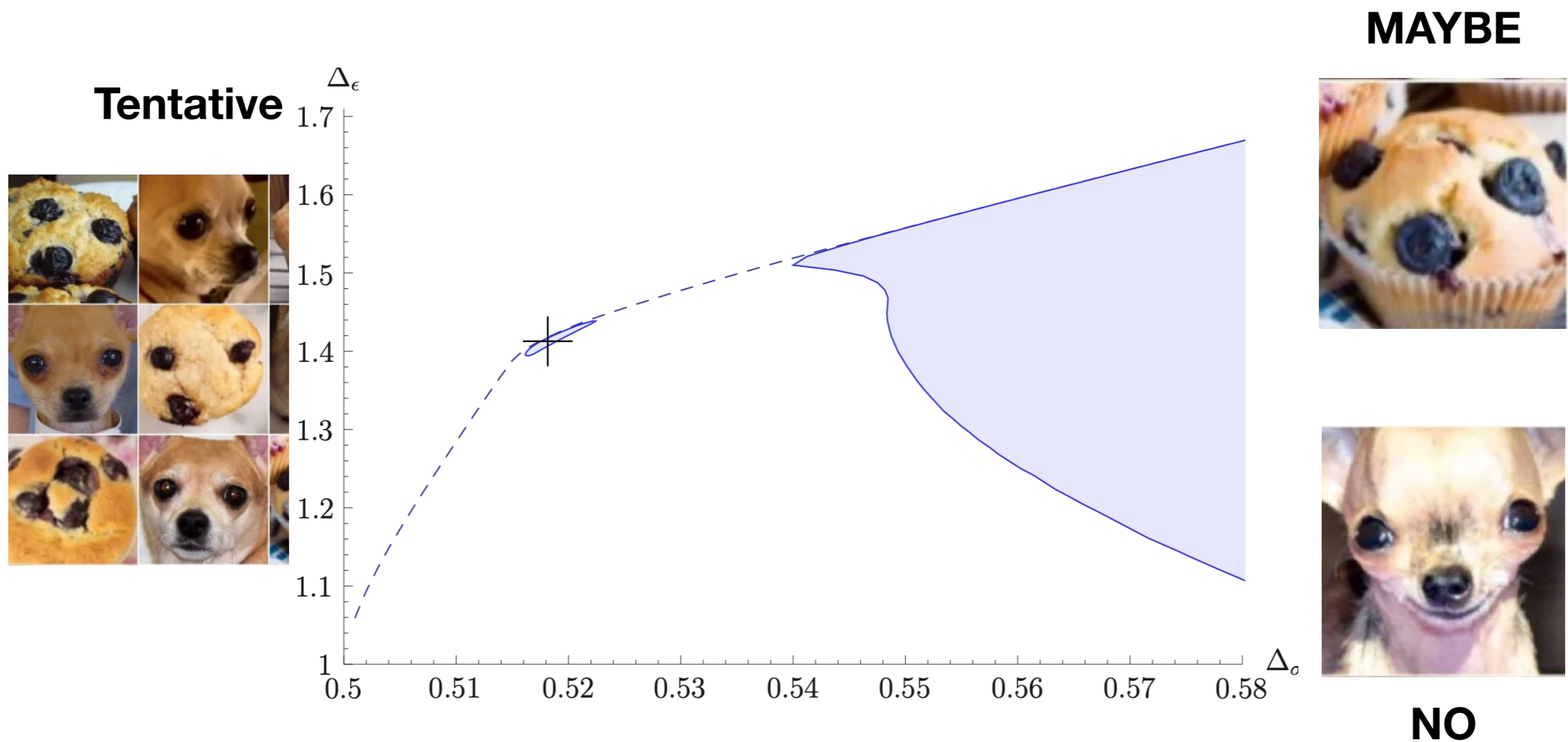
MAYBE



NO

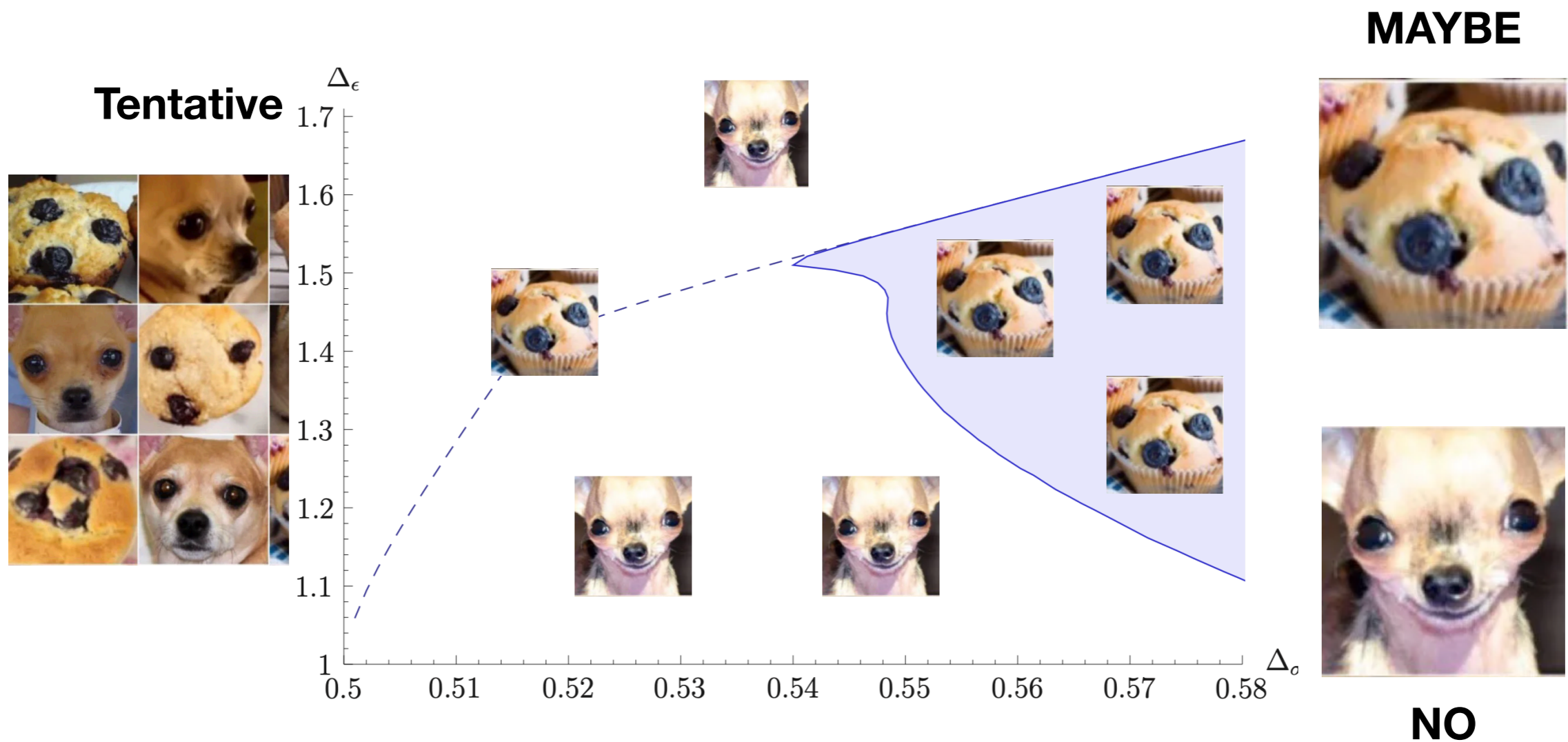
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Komargodski Zhiboedov 2013
Fitzpatrick Kaplan Poland Simmons Duffin 2012
Alday AB 2013

Analytical approach

Another viable approach is analytic.

In particular it is possible to find regimes in which one can find analytic solutions to the crossing relations.

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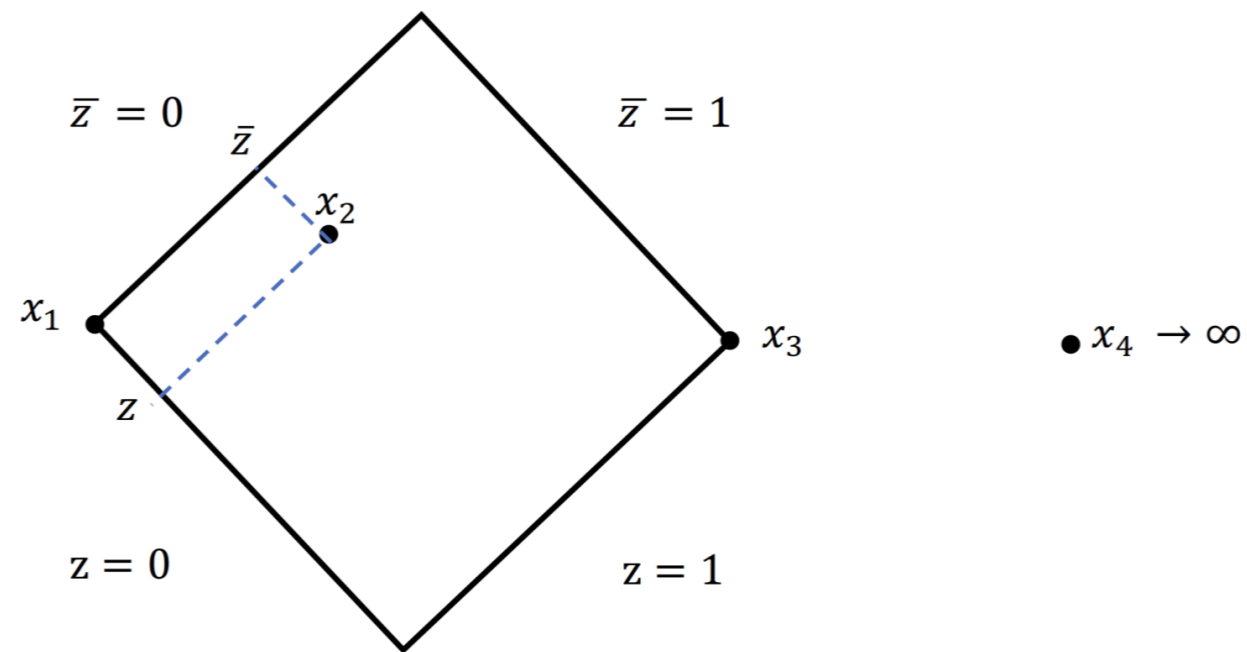
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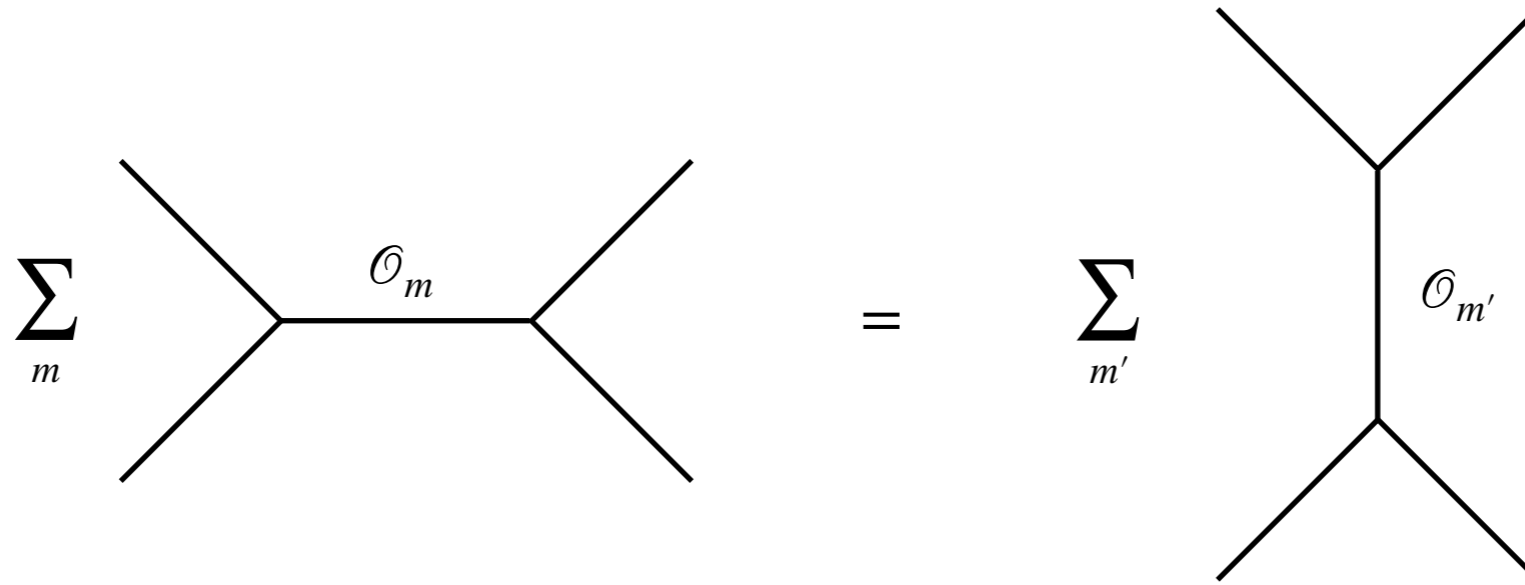
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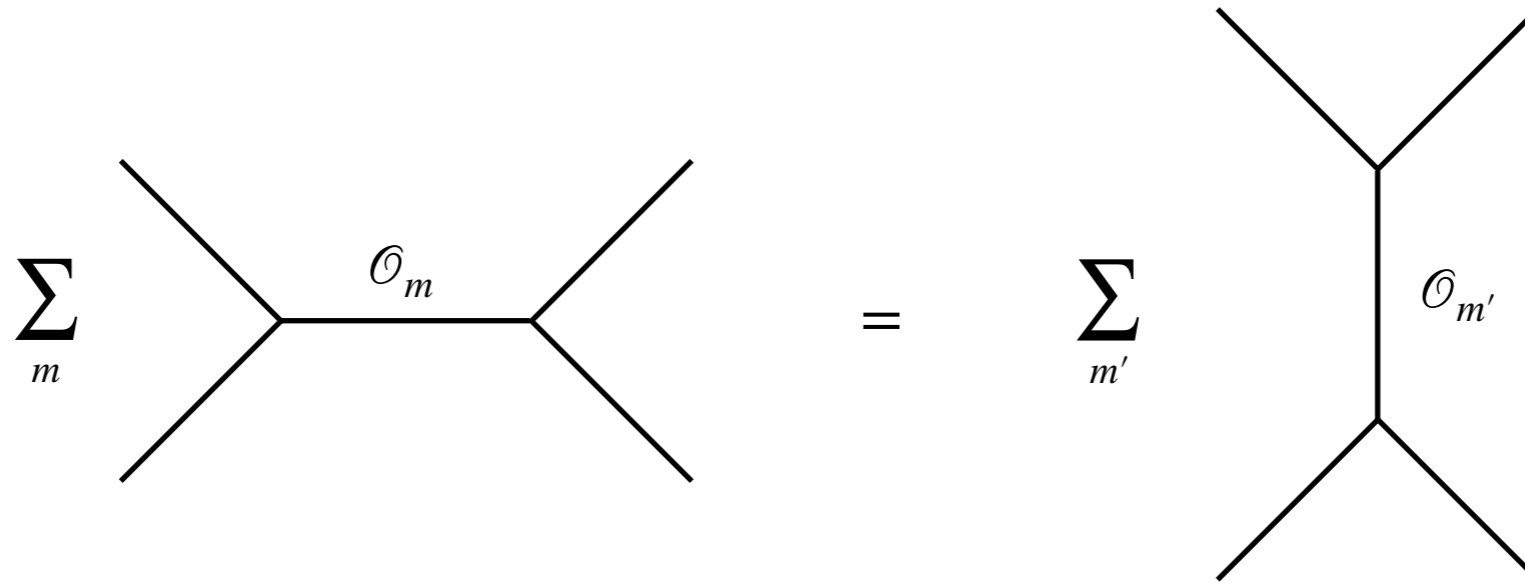
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“double” trace

Large spin sector

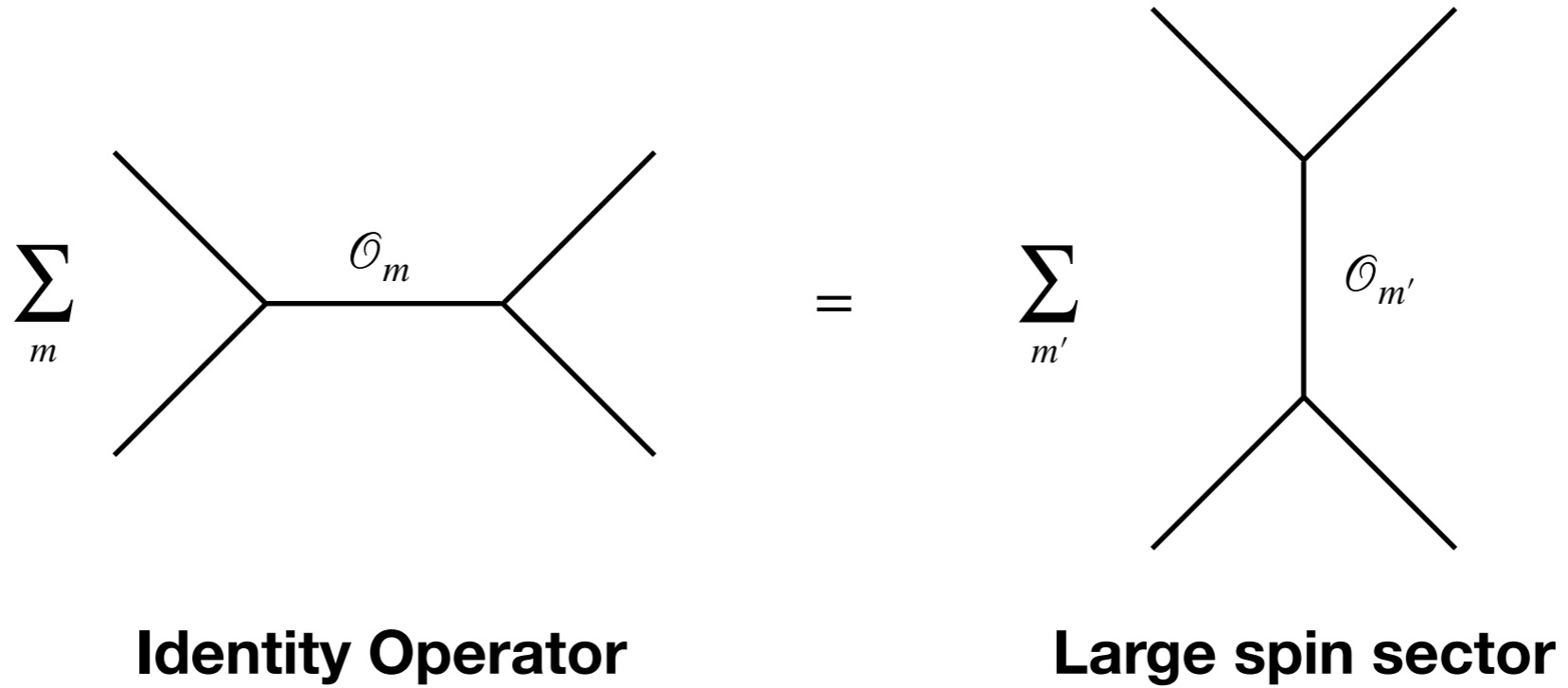


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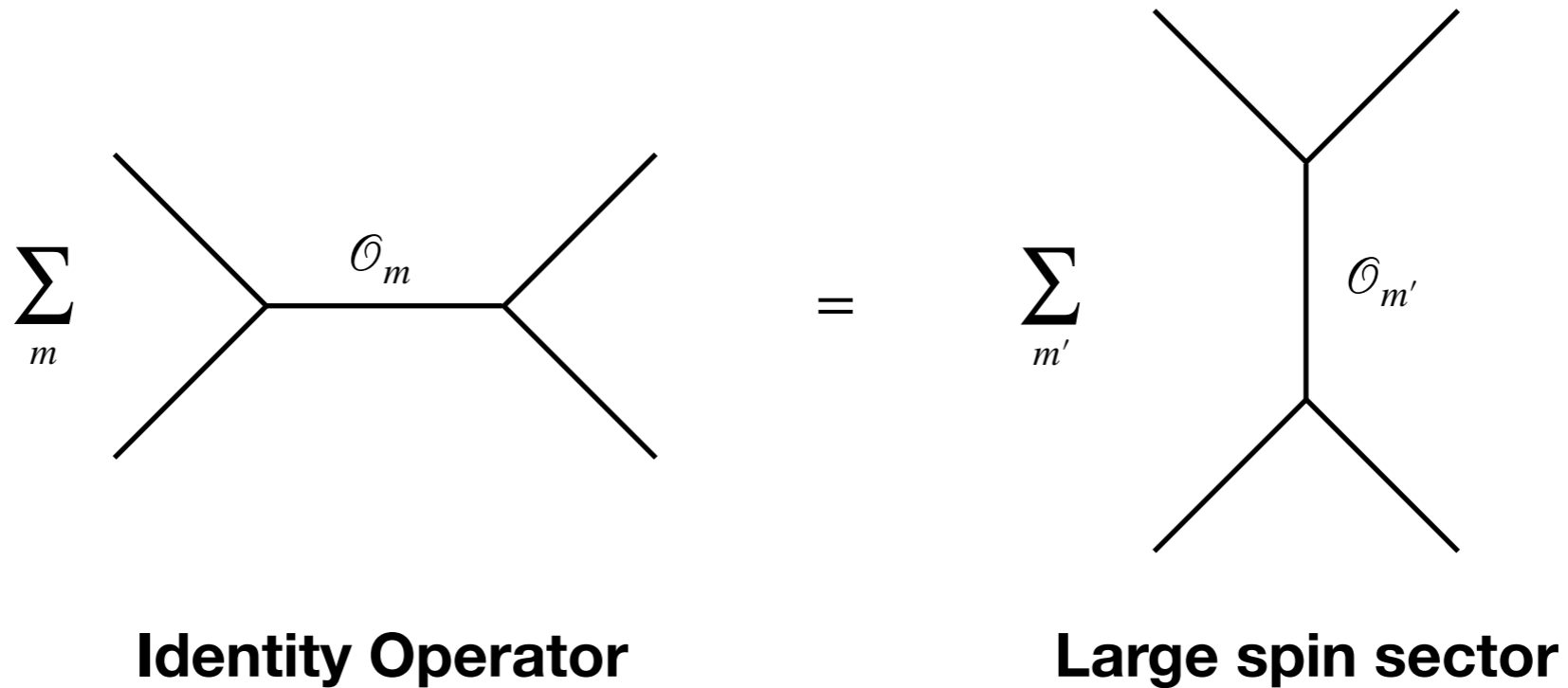


Identity Operator

Large spin sector

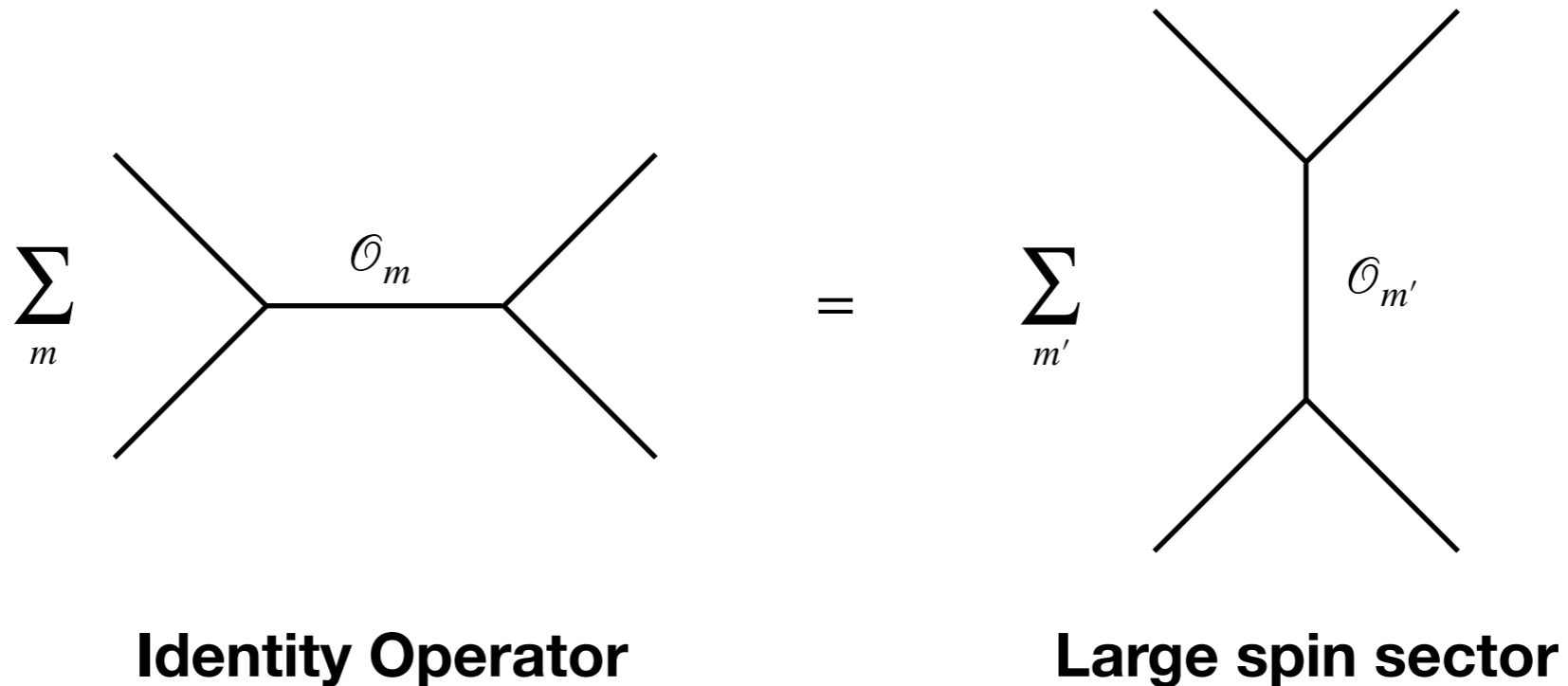


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We need to have infinitely many operators whose conformal dimension approaches $\Delta = 2\Delta_{\mathcal{O}} + 2n$ (double traces) and with very large spin.

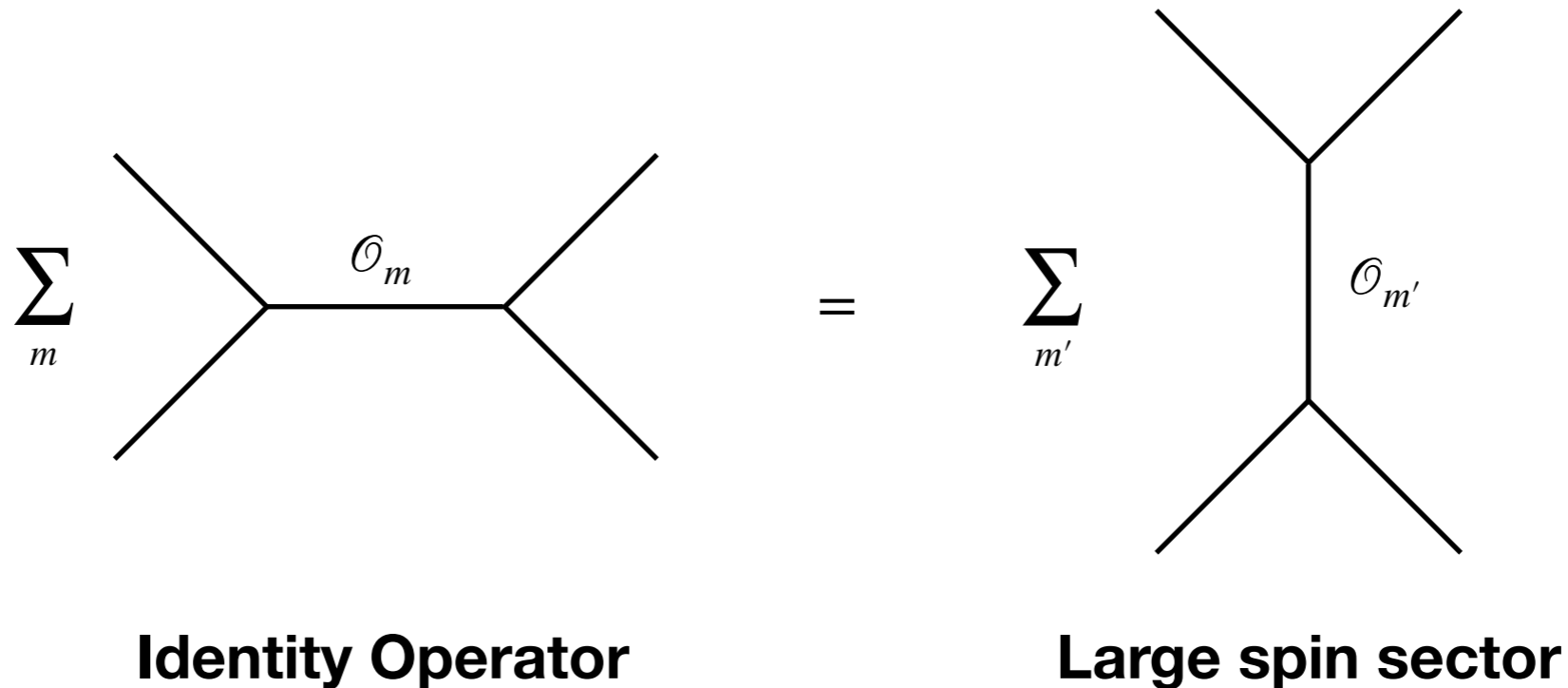
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It is possible to use the Casimir equation to iteratively find all the $1/\ell$ corrections and resum them to extrapolate for finite values of the spin.

Inversion formula

There is another approach to compute these quantities, less intuitive but computationally more powerful.

This approach allows inverting the OPE and prove why it is possible to resum the large spin series up to finite values of the spin.

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$$\mathbf{dDisc}[\mathcal{G}(z, \bar{z})] = \mathcal{G}_{Eucl}(z, \bar{z}) - \frac{1}{2} \mathcal{G}^{\cup}(z, \bar{z}) - \frac{1}{2} \mathcal{G}^{\ominus}(z, \bar{z})$$

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It has poles at the dimensions of the exchange operators with residues the square of the three point functions. The function is analytic in the spin for $\ell \geq 2$.

Applicability

The applicability of these methods is pretty vast, and it is mostly efficiently used when the theory has a small parameter (perturbation theory)

number of dimensions

$$d = 4 - \epsilon$$

rank of the gauge group

$$N$$

coupling constant

$$g \quad \text{or} \quad \lambda$$

...

AdS/CFT correspondence

CFT

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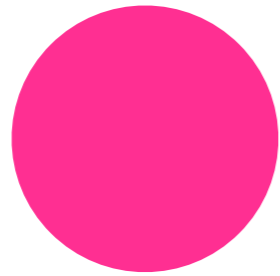
$$N \sim g_s^{-1}$$

$$\lambda = g_{YM}^2 N = (\alpha')^{-2}$$

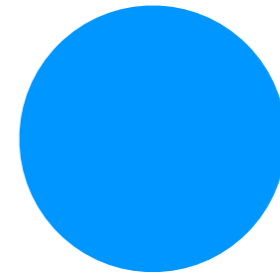
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Correlators in CFTs



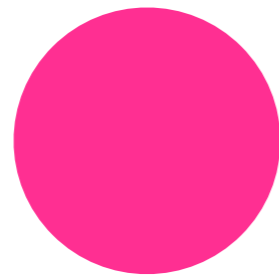
Amplitudes in AdS



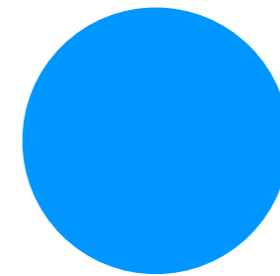
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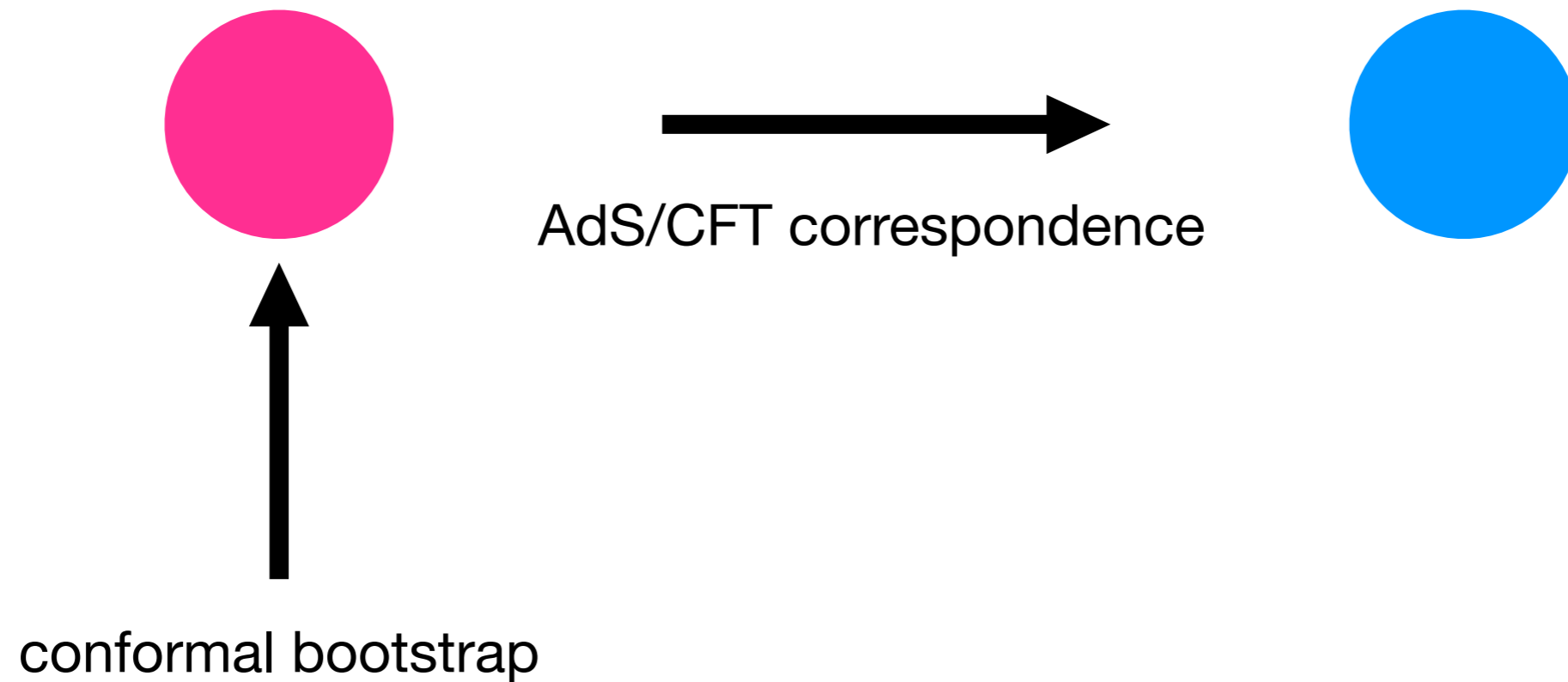
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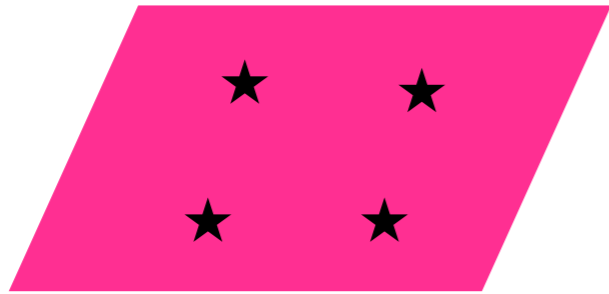
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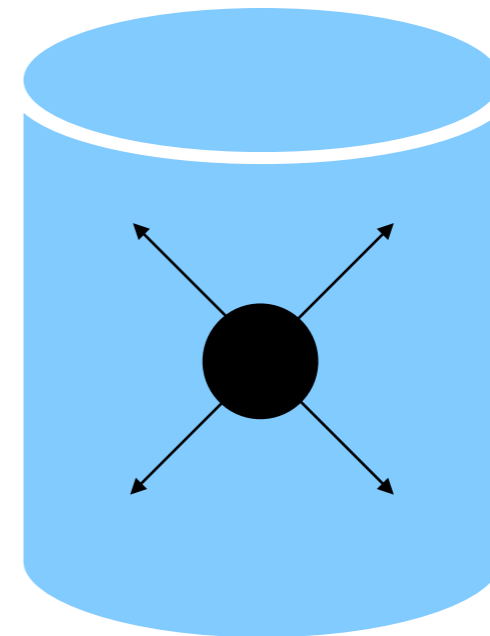
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$$\mathbb{R}^{d-1,1} = \partial\text{AdS}_{d+1}$$

Amplitudes in AdS



$$\text{AdS}_{d+1} \times S^q$$

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Double trace operators
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ignore the stress tensor

Large N

We expand all the quantities up to order N^{-4} :

$$\mathcal{G}(u, v) = \mathcal{G}^{(0)}(u, v) + \frac{1}{N^2} \mathcal{G}^{(1)}(u, v) + \frac{1}{N^4} \mathcal{G}^{(2)}(u, v) + \dots$$

$$\Delta = \Delta^{(0)} + \frac{1}{N^2} \gamma^{(1)} + \frac{1}{N^4} \gamma^{(2)} + \dots$$

$$c_{\Delta, \ell}^2 = k_{\Delta, \ell}^{(0)} + \frac{1}{N^2} k_{\Delta, \ell}^{(1)} + \frac{1}{N^4} k_{\Delta, \ell}^{(2)} + \dots$$

The idea is to compute order by order in N , they CFT data. The main aim is to understand if we can predict the order N^{2k} using the $N^{2(k-1)}$ one.

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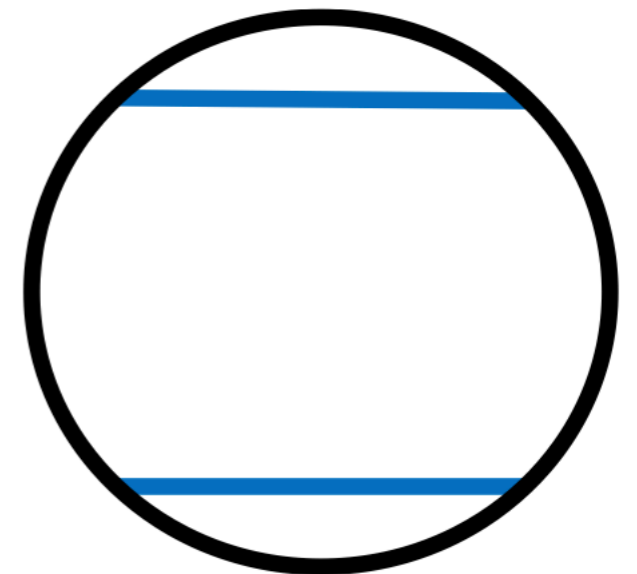
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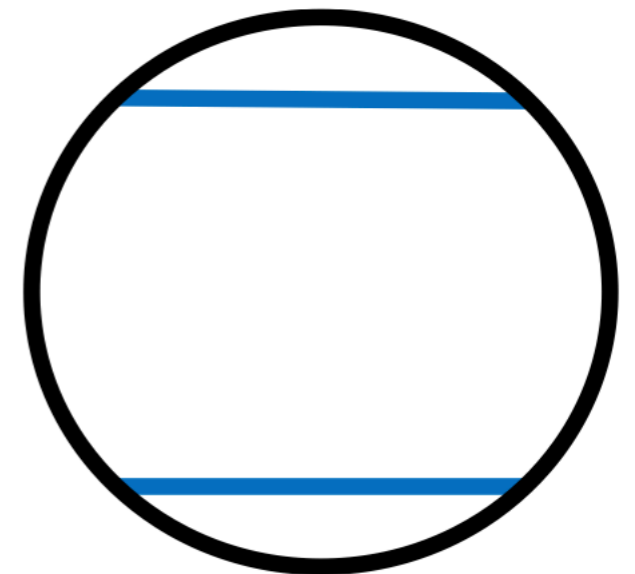
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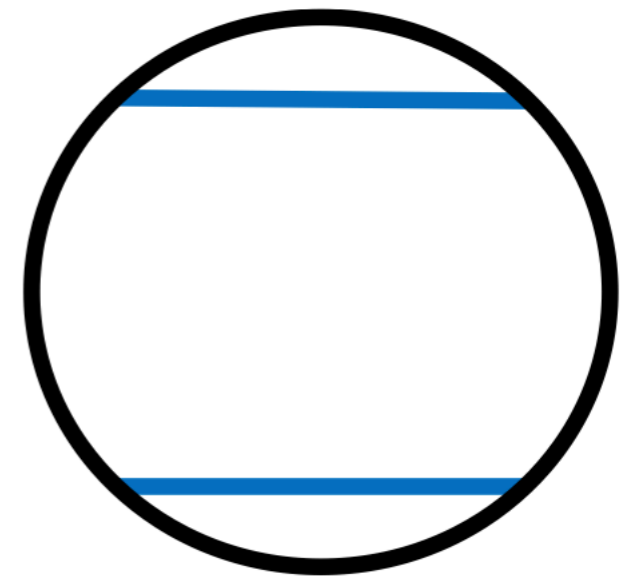
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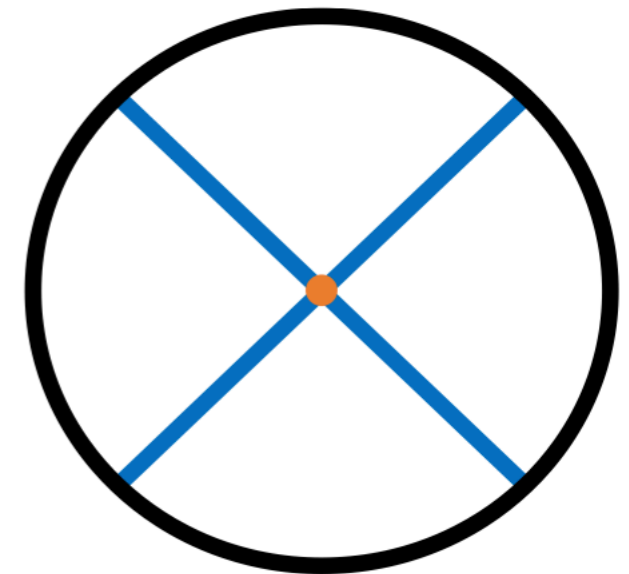
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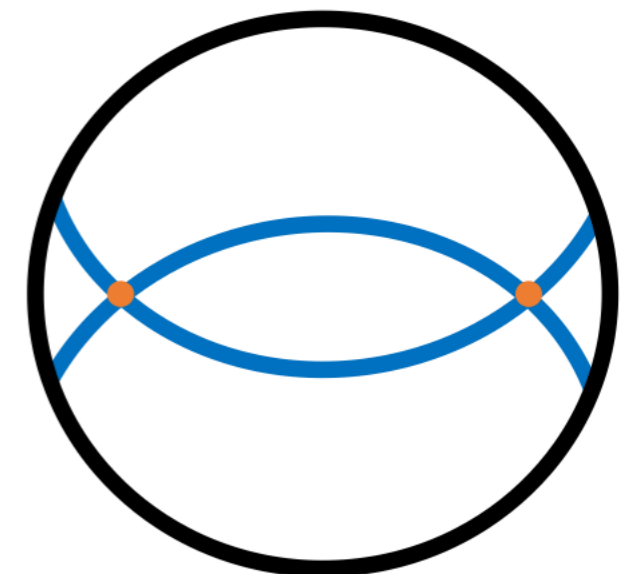
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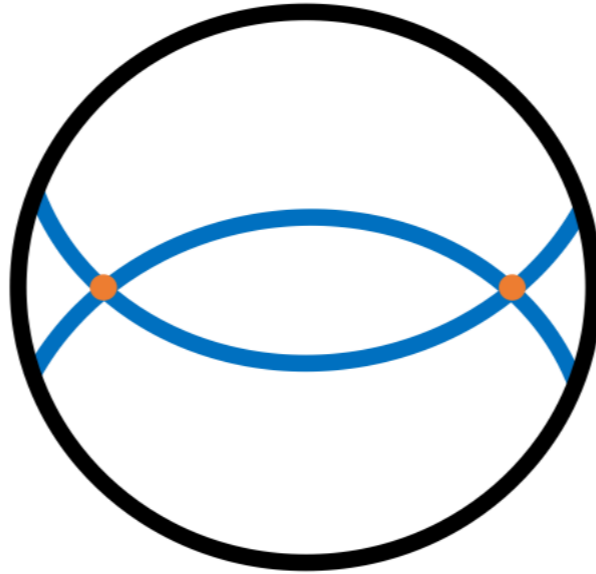
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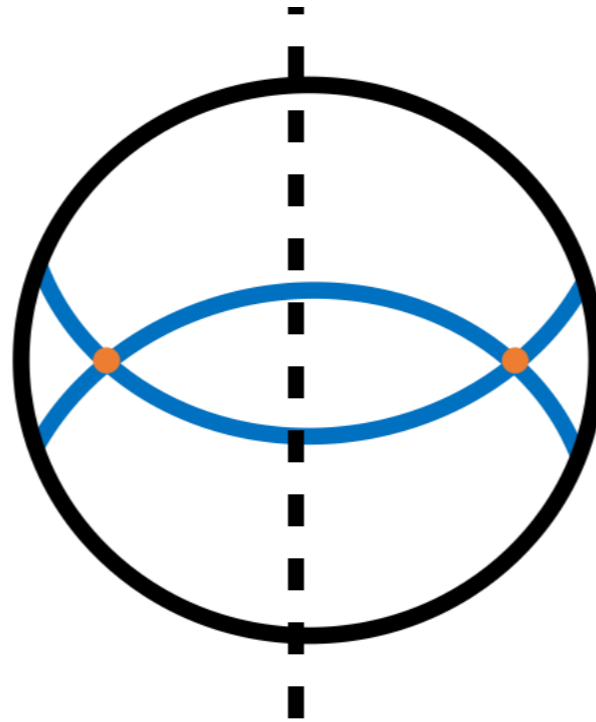


In pictures

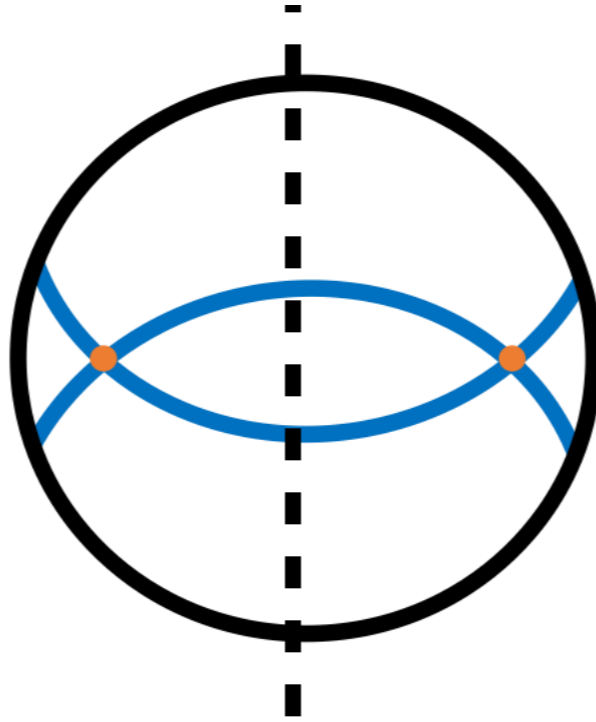
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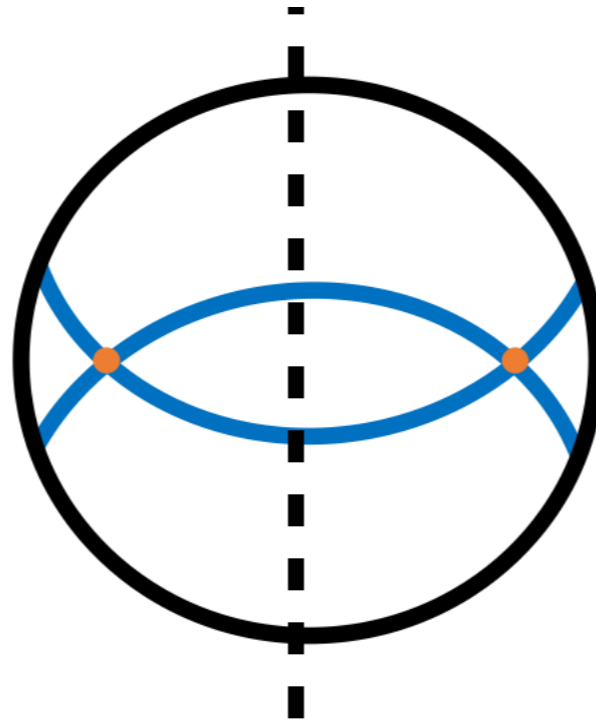


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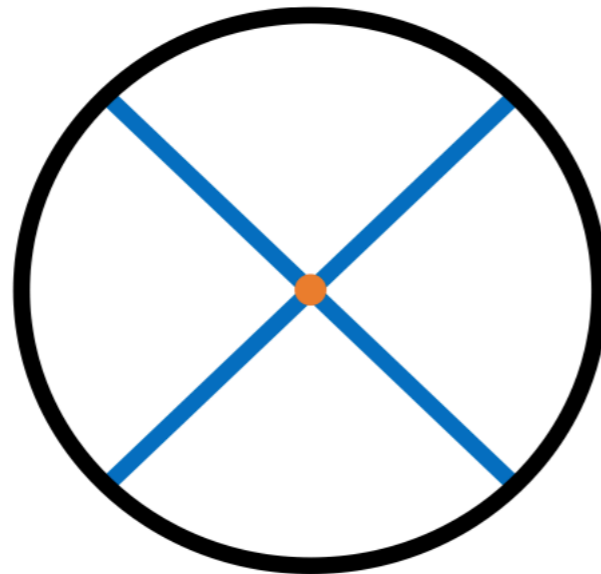


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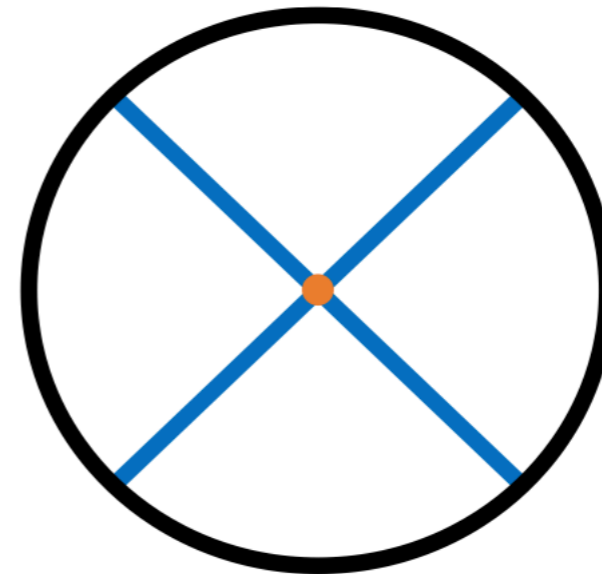
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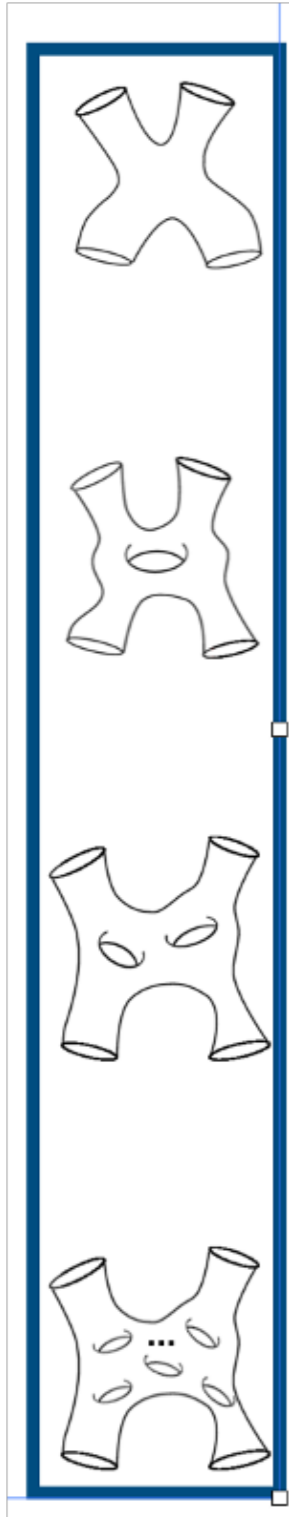
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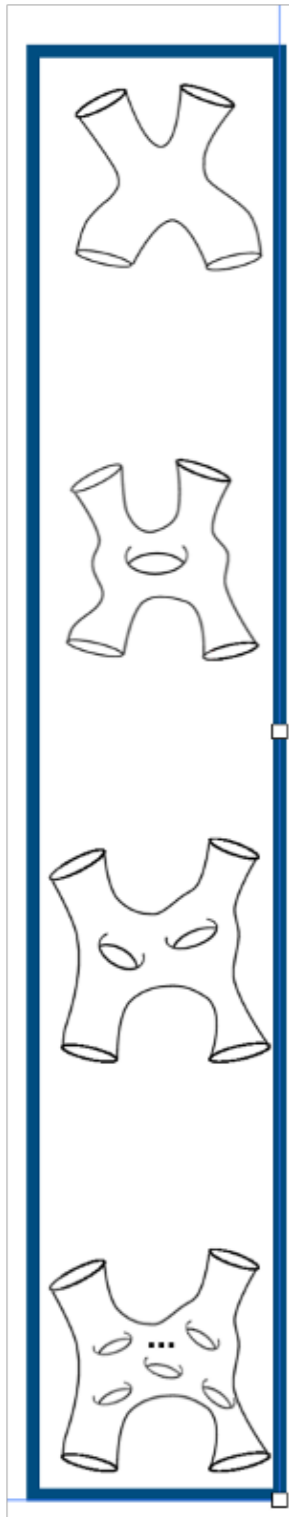


Other applications



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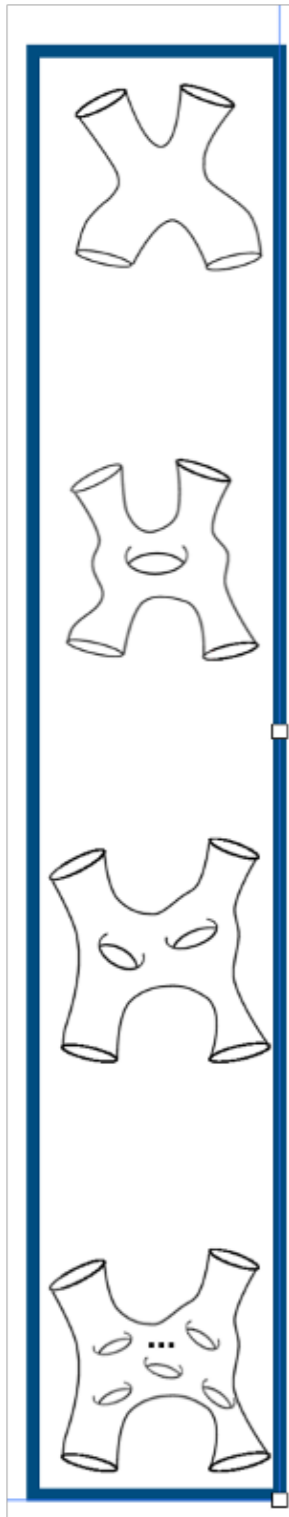
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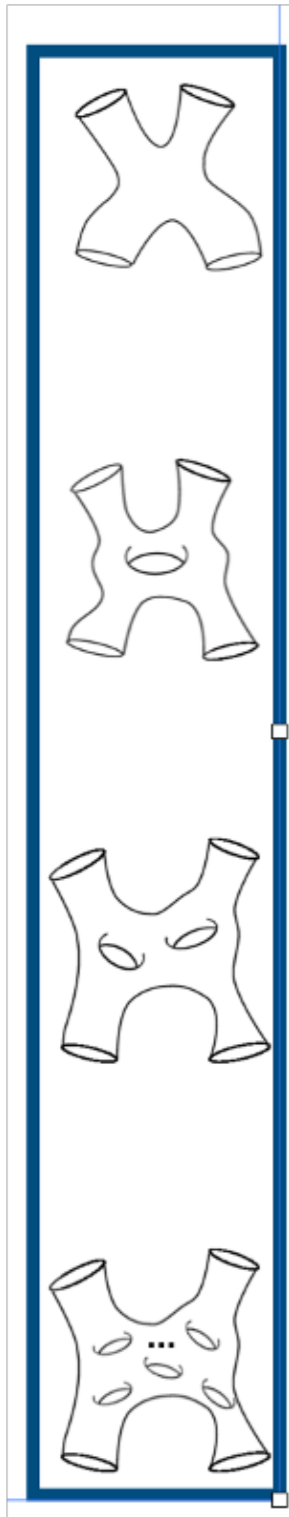


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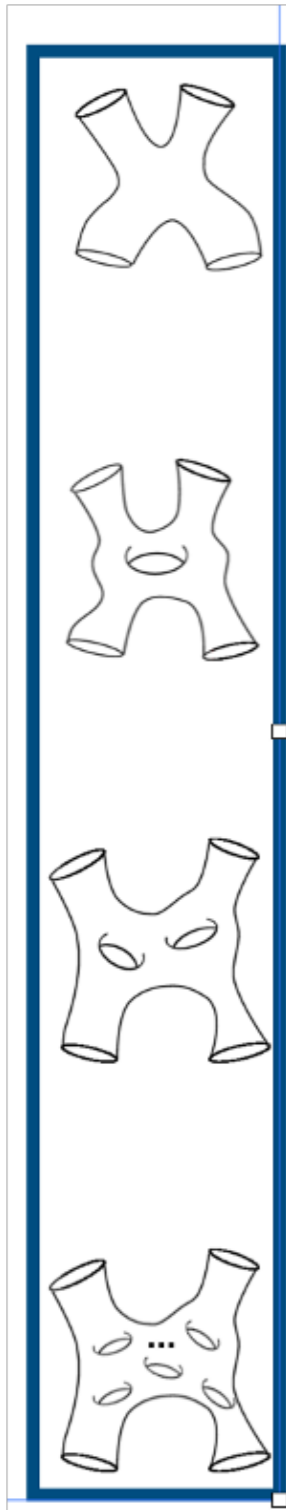
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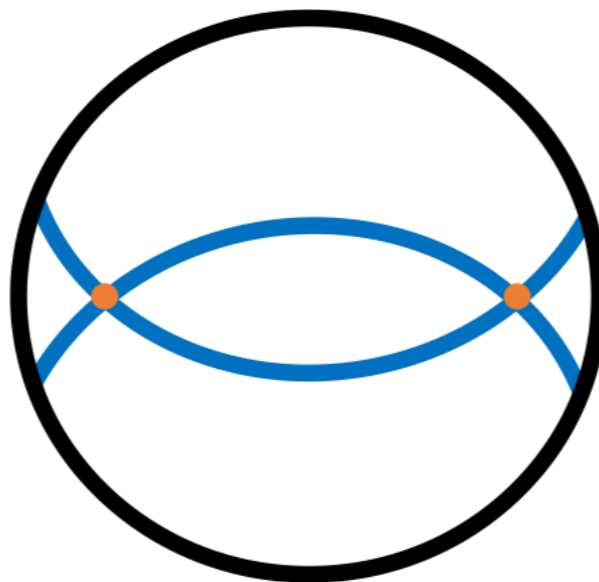
Provides a unique framework to access scattering amplitudes in curved space-times, which are generically very hard/impossible to compute with other methods.

Conclusions

I presented a framework to analytically study CFT, using only the symmetries and the presence of an OPE expansion.

Mapping between singularities and OPE data.

Access amplitudes in AdS using the AdS/CFT correspondence



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Extend the long distance behaviour to any distance: **continuous limit**