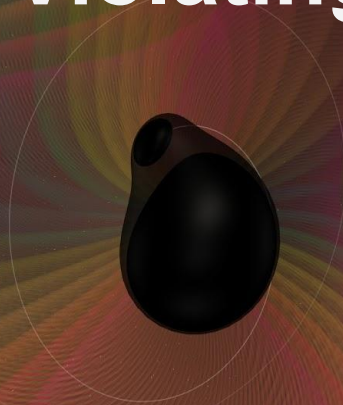


Regular Black Holes in Lorentz-Violating Gravity



Jacopo Mazza

International School for
Advanced Studies – Trieste, Italy

Trieste

Septembre 5, 2023

**XXV SIGRAV
Conference**





Based on:

JM, S. Liberati, JHEP 03 (2023) 199 [[2301.04697](#)]

“Regular Black Holes and Ultra-Compact Objects in Lorentz-Violating Gravity”



Why

Lorentz Invariance Violations (**LIV**) in gravity?

- LI well tested in matter, not so much in gravity
- Many QG scenarios point to LIV
[LIV cornerstone of QG phenomenology]
- LIV can help build QG theories

Hořava gravity

Framework: (non-projectable) **Hořava gravity**

QFT of gravity

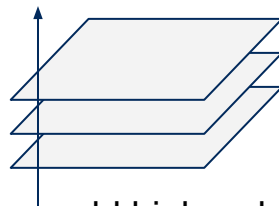
with higher
derivatives

no ghosts

power-counting
renormalisable

how?

(probably perturbatively
renormalisable)



fundamental split
between
space & time

add higher derivatives only in space
(not time)

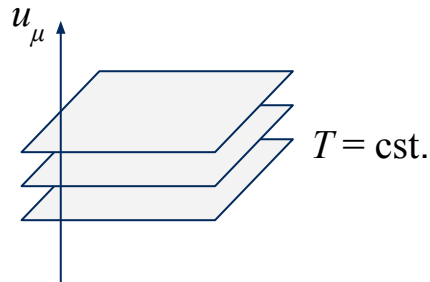
Focus on low-energy limit:
Kronometric theory

————— scalar-tensor theory, related to Einstein-æther theory

Building blocks

Basic ingredients

$$\left(\begin{array}{l} \text{metric} \\ \mathbf{g}_{\mu\nu} \\ \\ \text{khronon } T(x) \text{ or } u_\mu = N \nabla_\mu T \\ \text{æther} \end{array} \right.$$



coupling with æther
gives LIV

constraint

$$\left(\begin{array}{l} \mathbf{u}_\mu \mathbf{u}^\mu = +1 \\ \text{æther everywhere} \\ \text{timelike} \end{array} \right.$$

there can be
superluminal causal
signals

æther determines
causal structure

Black Holes

Surprisingly,
BHs exist!

Killing horizons (**KHs**)
not causal horizons

new kind of horizons:
universal horizon (UHs)

compact leaf of preferred foliation

Static and spherically symmetric IR solution (in corner of parameter space):
Schwarzschild + “painted on” æther

$$ds^2 = F(r)dv^2 - dvdr - r^2 d\Omega^2$$

$$F(r) = 1 - \frac{2M}{r}$$

M mass

$$u^\mu \partial_\mu = A(r)\partial_v + y(r)\partial_r$$

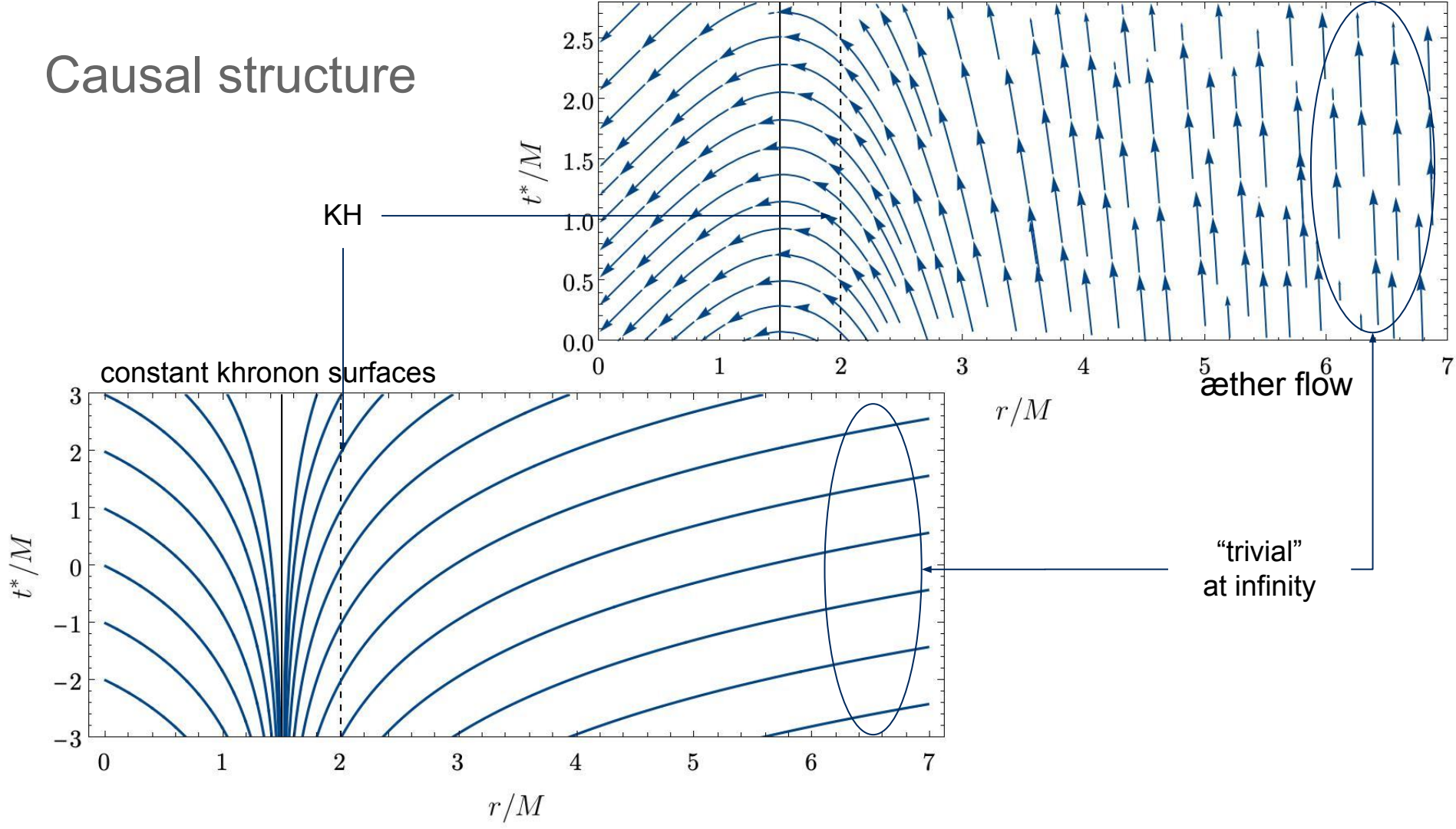
$$y(r) = -\frac{r_{\text{ae}}^2}{r^2}$$

$$r_{\text{ae}} = M \left[\frac{27}{16} \right]^{1/4}$$

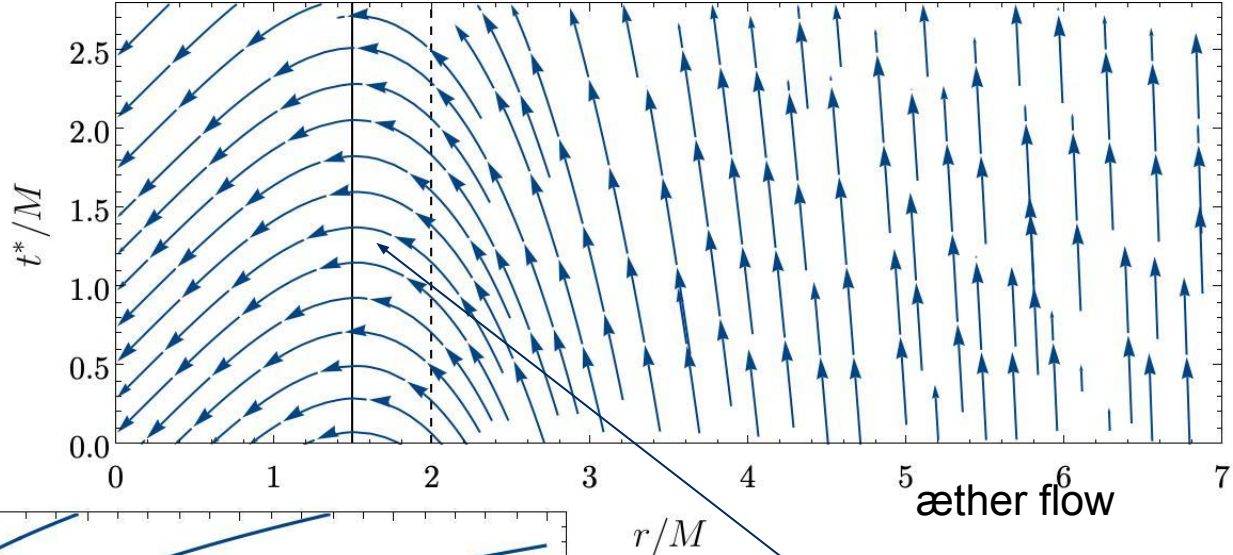
(choice)

spacetime singularity at $r = 0$

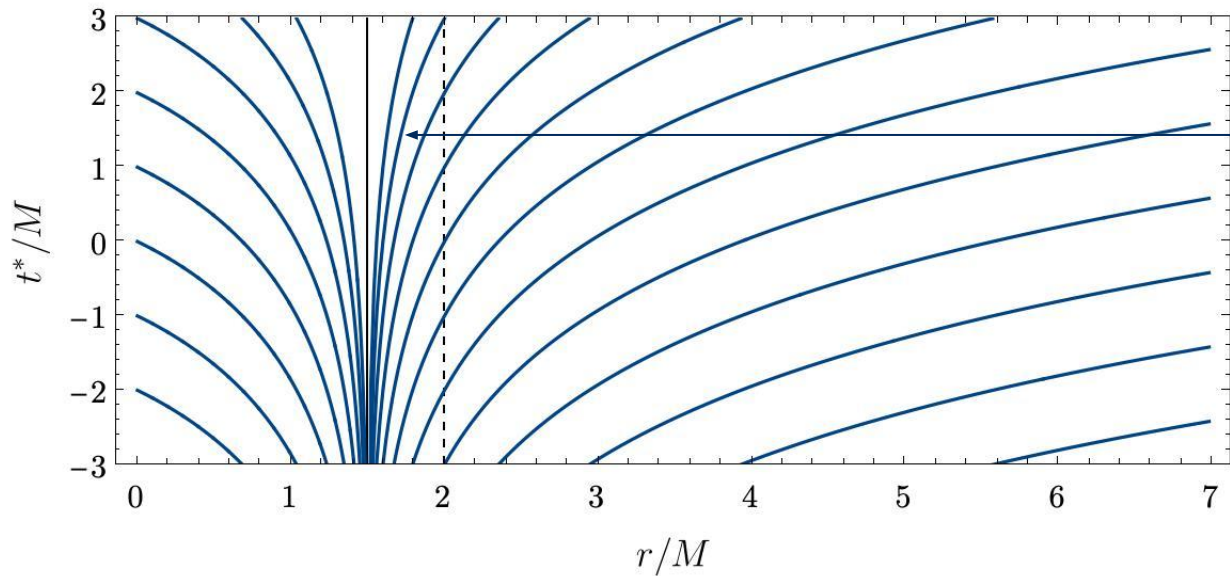
Causal structure



Causal structure



constant khronon surfaces



UH

universal horizon:
æther \perp Killing vector

- traps modes of any speed
- some BH mechanics
- seems to have thermal properties

Goal



If Hořava is UV-complete QG, its BHs might be non-singular.

What would they look like?

Build (effective) models of non-singular BHs
[“regularisations” of low-energy singular solution]



Equations are hard, so not looking for solutions

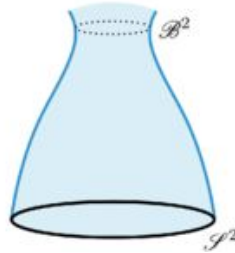
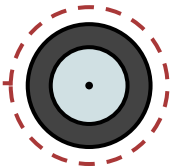
Models quantitatively wrong, but qualitatively
good

Two classes

Qualitatively, two classes of
RBHs

Simply Connected
(SC)

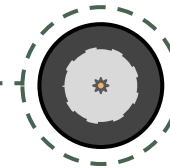
regular centre
multiple horizons
("proper" RBHs)



Carballo-Rubio, Di Filippo, Liberati, and Visser,
PRD 101, 084047 [\[1911.11200\]](#)

Multiply Connected
(MC)

wormholes
hidden by horizon
require topology change



Regularising BHs

Start from singular solution and replace



$$M \mapsto m(r)$$

$$m(r) = c \ell^{-2} r^3 + \mathcal{O}(r^4)$$

e.g. Hayward

$$m(r) = M \frac{r^3}{r^3 + 2M\ell^2}$$

in metric *and* æther

$$y_{\text{SC}}(r) = -\frac{r_{\text{ae}}^2(r)}{r^2}$$

$$r_{\text{ae}}(r) = m(r) \left[\frac{27}{16} \right]^{1/4}$$

for $r \rightarrow 0$



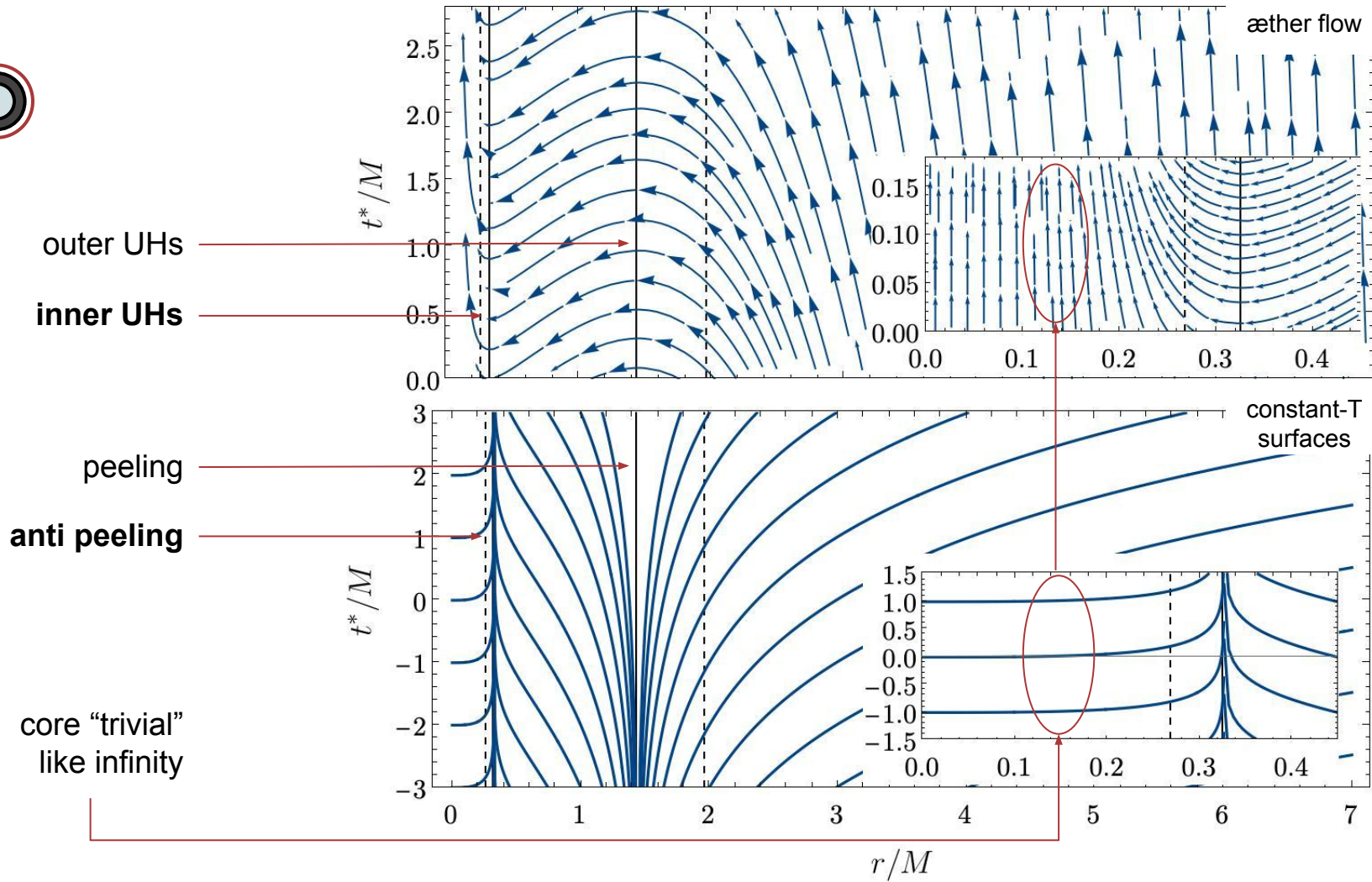
$$r \mapsto R(r)$$

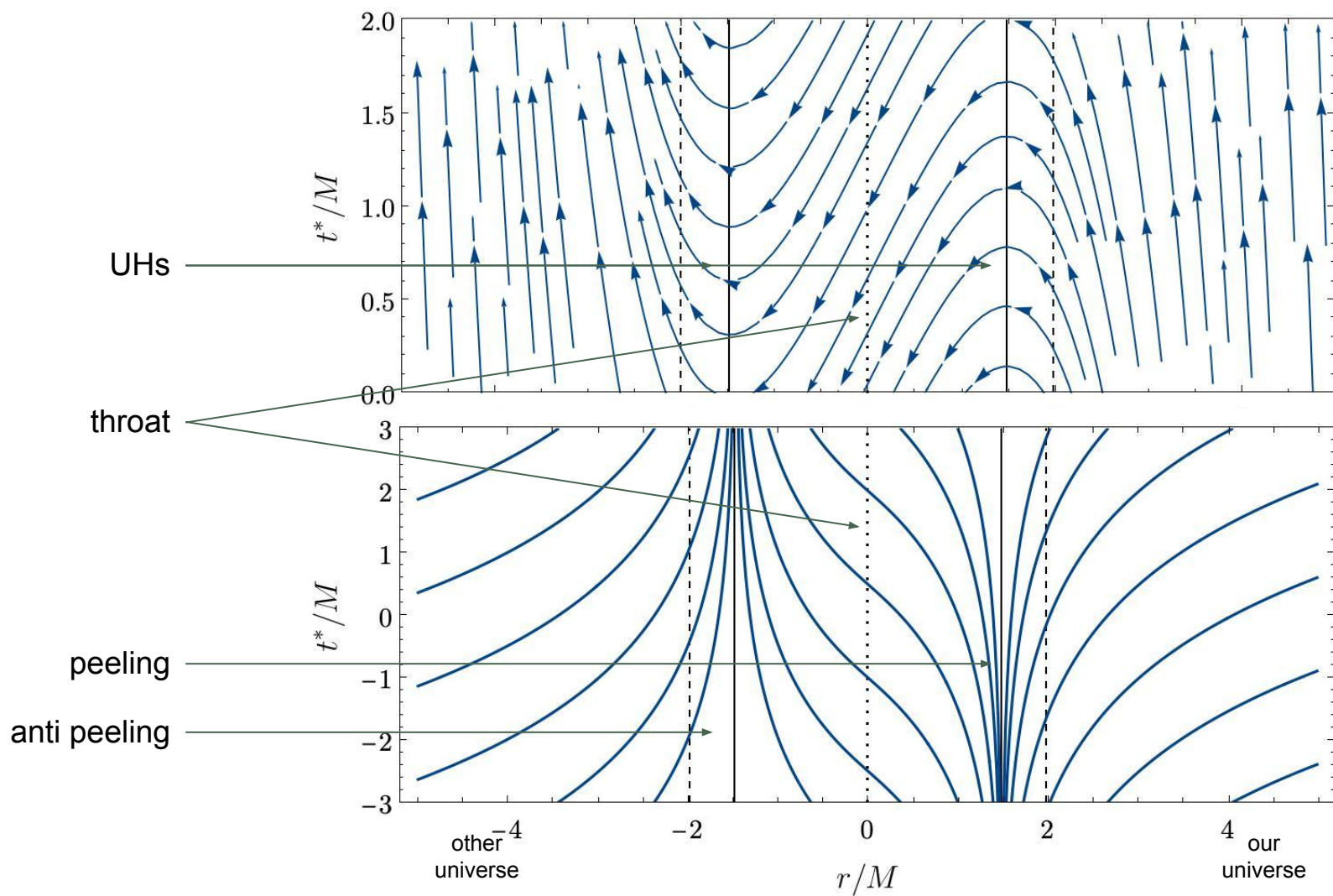
$$R(r) = \ell + \mathcal{O}(r)$$

e.g. Simpson–Visser

$$R(r) = \sqrt{r^2 + \ell^2}$$

$$y_{\text{MC}}(r) = -\frac{r_{\text{ae}}^2}{R(r)^2}$$





Recap & Outlook

Effective models of RBHs in Lorentz-violating gravity



- outer *and* inner horizons (KH/UH)



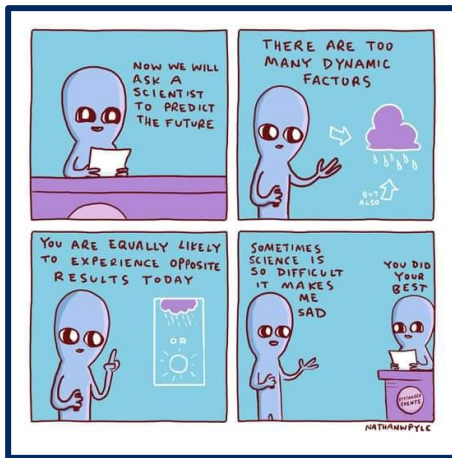
- wormhole throat hidden by horizons (KH/UH)

Common features:

- BHs *and* horizonless objects
- metric+æther simple
 - horizons' features (location, surface gravity, etc.) are tunable

Several open questions:

- Do RBHs really exist in Hořava?
- Are UHs stable?
 - mass inflation?
- How much of familiar BH physics translate to LIV gravity?
 - thermodynamics?
 - rotation, dynamics?





Thanks!

Get in touch:
jacopo.mazza@sissa.it

BckUp: T-theory

$$S = -\frac{1}{16\pi G} \int d^4x \sqrt{-g} \left[R + \lambda(\nabla_\mu u^\mu)^2 + \beta \nabla_\mu u^\nu \nabla_\nu u_\mu + \alpha a_\mu a^\mu \right]$$

$$a^\mu = u^\nu \nabla_\nu u^\mu$$

Constraints

$$|\beta| \leq 10^{-15}$$

$$\text{either } \begin{cases} |\alpha| \leq 10^{-7} & \lambda \text{ unconstrained} \\ |\alpha| \leq 0.25 \times 10^{-4} & \lambda \approx \alpha / (1 - 2\alpha) \end{cases}$$

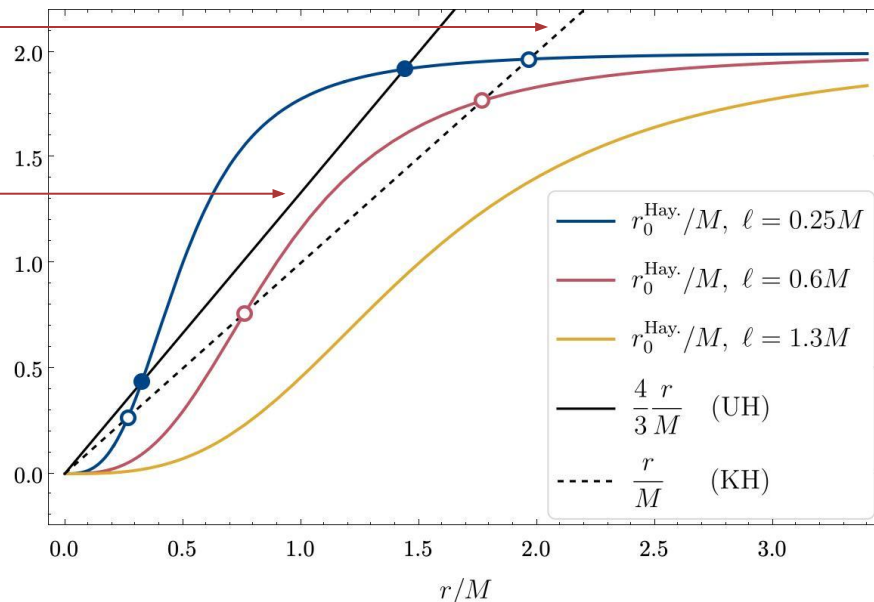
BckUp: Hayward



A common choice: Hayward $m(r) = M \frac{r^3}{r^3 + 2M\ell^2}$

intersections: KHs

intersections UHs



shape of $m(r)$
changes with ℓ
so do horizons

horizons come in pairs
(outer/inner)

UHs always in trapped
region

BckUp: SV



Simpson-Visser regularisation
a.k.a. black bounce

$$r \mapsto \sqrt{r^2 + \ell^2}$$

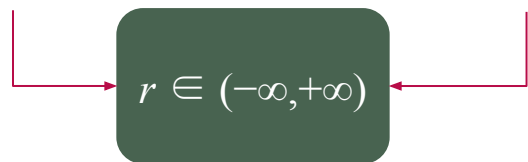
$$ds^2 = - \left(1 - \frac{2M}{\sqrt{r^2 + \ell^2}} \right) dt^2 + \frac{dr^2}{1 - \frac{2M}{\sqrt{r^2 + \ell^2}}} + (r^2 + \ell^2) d\Omega^2$$

reflection symmetry

$$r \rightarrow -r$$

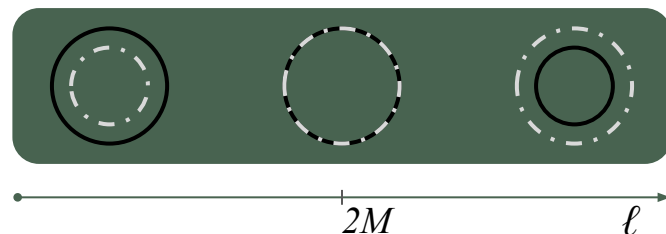
$r = 0$ sphere of area

$$4\pi\ell^2$$



two identical “universes”, glued at $r = 0$

horizon(s): Y/N depending on ℓ



BkcUp: Effective Sources

$$G_{\mu\nu} - T_{\mu\nu}^{\text{ae}} = T_{\mu\nu}^{\text{eff}} \neq 0$$

measure in preferred frame



$$(c_a = c_\sigma = 0)$$



$$\rho_G^{(u)} = -p_G^{(s)} = \frac{12M^2\ell^2}{(r^3 + 12M\ell^2)^2}, \quad \rho_\theta^{(u)} = \frac{243M^6\ell^4r^6}{6(r^3 + 2M\ell^2)^6},$$

$$p_G^\perp = -\frac{24M^2\ell^2(M\ell^2 - r^3)}{(r^3 + 2M\ell^2)^3}, \quad p_\theta^{(s)} = p_\theta^\perp = \frac{243M^5\ell^2r^6(r^3 - 2M\ell^2)}{6(r^3 + 2M\ell^2)^6},$$

$$\rho_G^{(u)} = -\frac{\ell^2}{8\varrho^8}(8\varrho^4 - 32\varrho^3M + 27M^4), \quad \rho_\theta^{(u)} = 0,$$

$$p_G^{(s)} = -\frac{\ell^2}{8\varrho^8}(8\varrho^4 + 27M^4), \quad p_\theta^{(s)} = 0,$$

$$p_G^\perp = \frac{\ell^2(\varrho - M)}{\varrho^5}, \quad p_\theta^\perp = 0,$$

$$(\varrho = \sqrt{r^2 + \ell^2})$$

$$\mathcal{O}(r^{-6})$$

$$[\mathcal{O}(r^{-4}) \text{ with } c_a]$$

at infinity

$$\mathcal{O}(r^{-4})$$